

Fundamental Physics with High-Energy Cosmic Neutrinos

Mauricio Bustamante

Niels Bohr Institute, University of Copenhagen

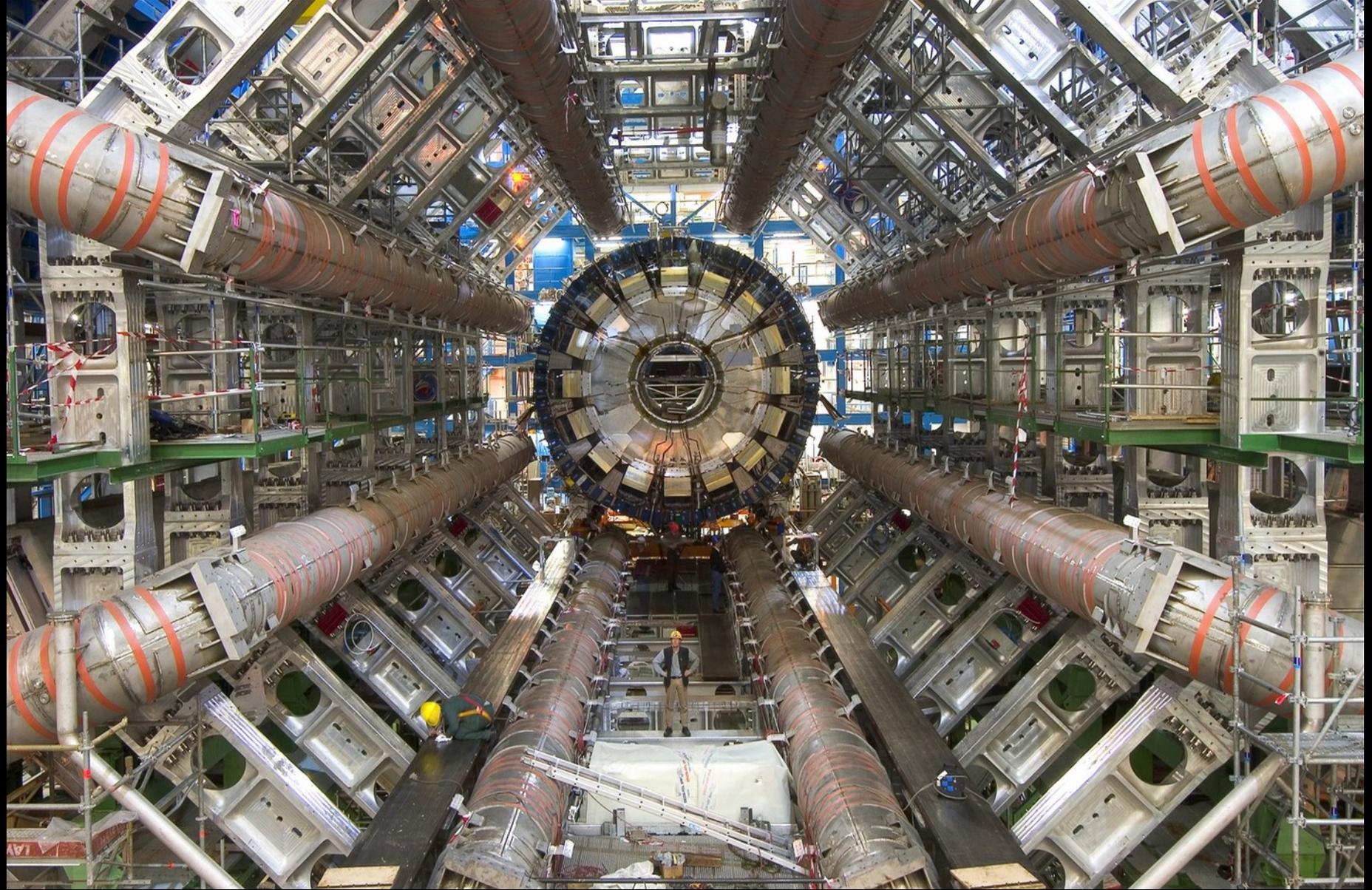
Multimessengers @ Prague
December 05, 2019

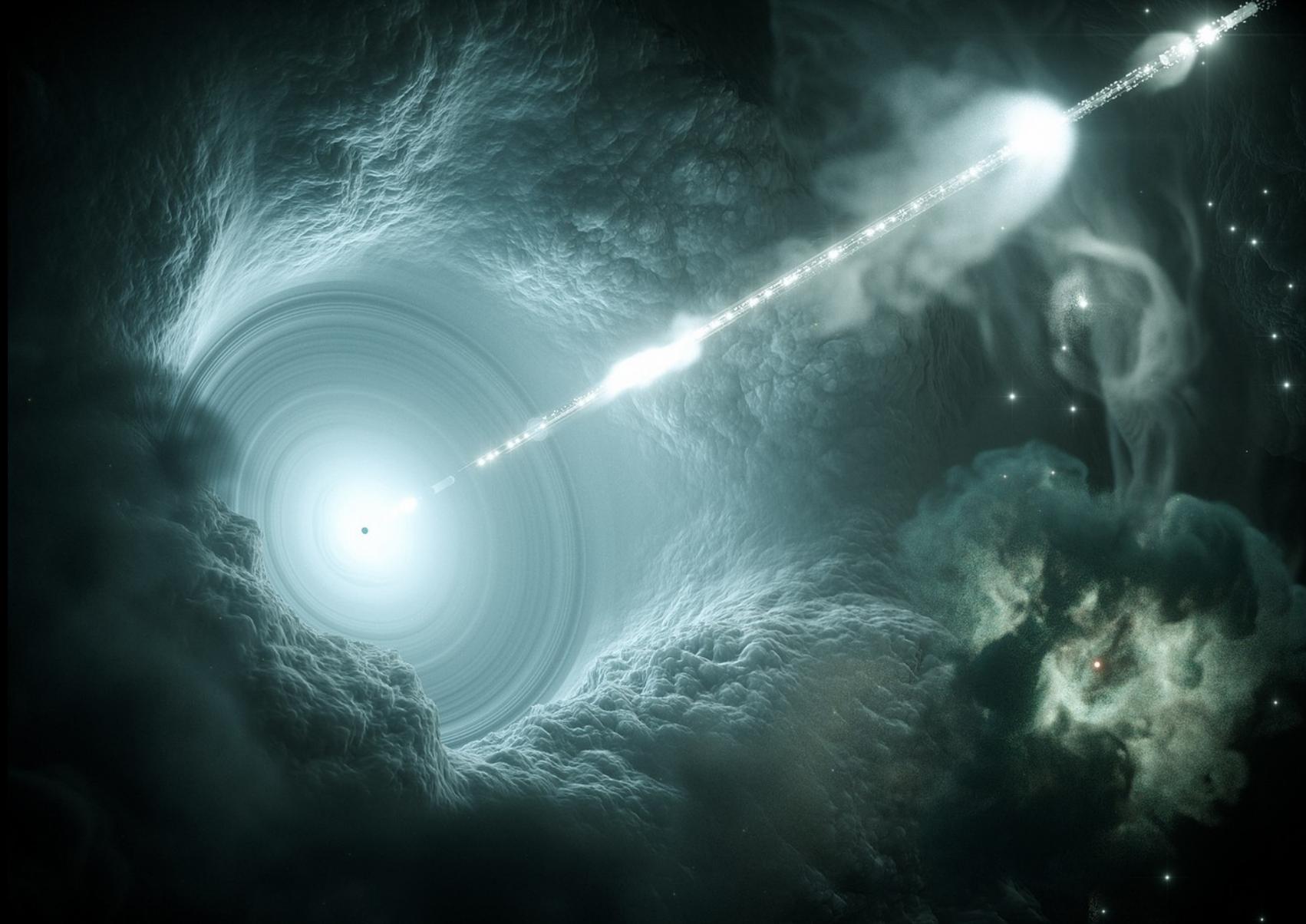
UNIVERSITY OF
COPENHAGEN

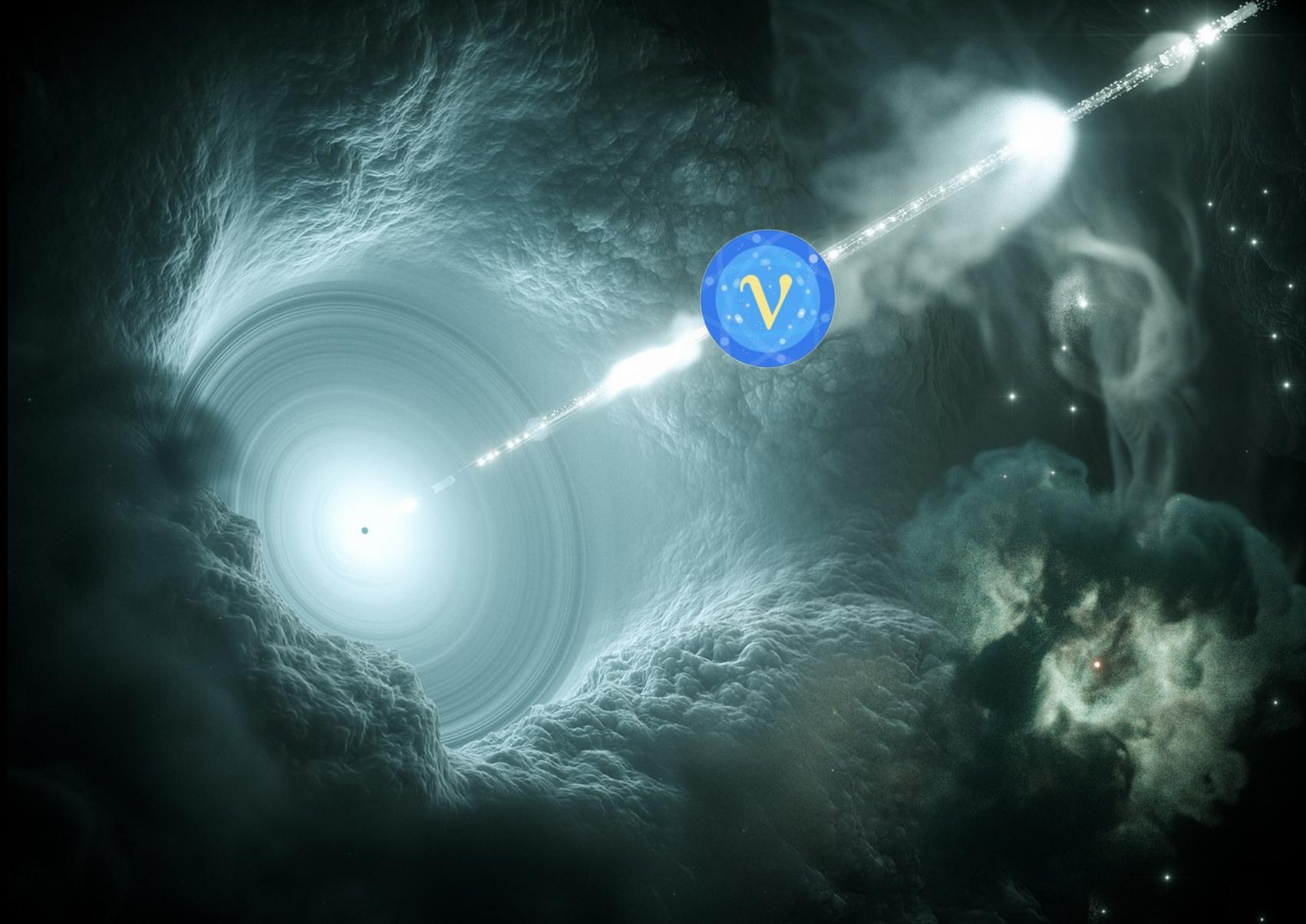


VILLUM FONDEN









Why study fundamental physics with HE cosmic ν ?

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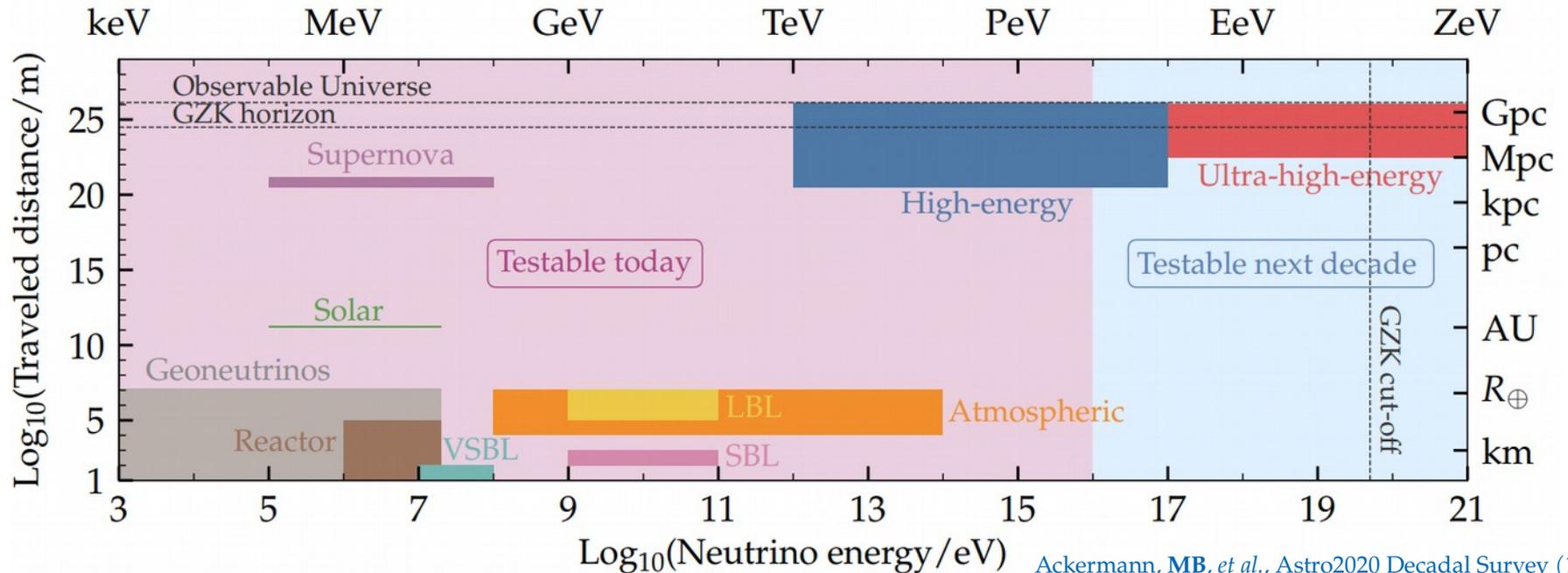
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↳ Probe physics at new energy scales

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- 2 They have the **longest baselines** (\sim Gpc)
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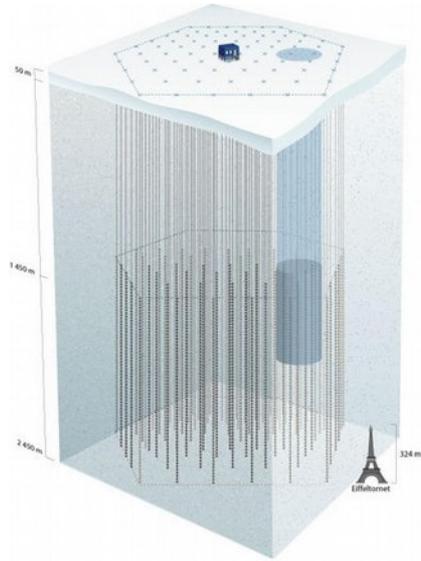
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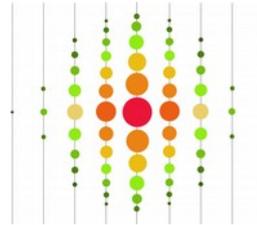
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- 5 It comes *for free*

IceCube (8 years)

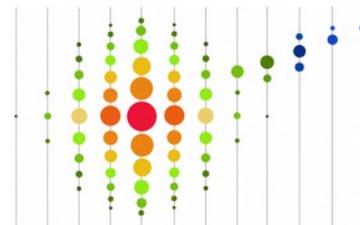
km³ in-ice
Cherenkov detector



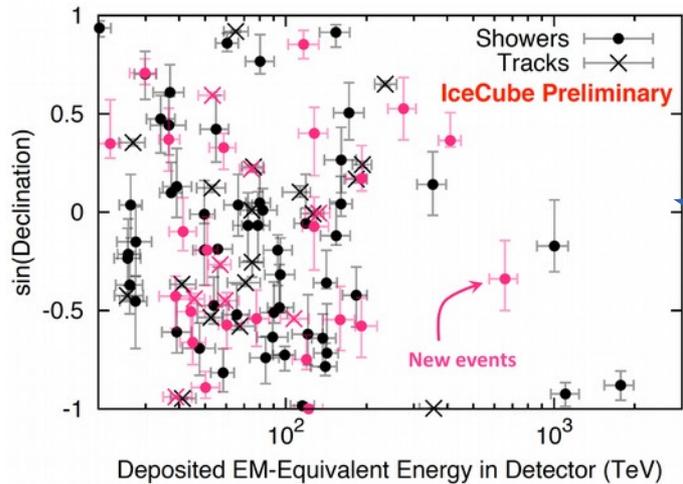
Showers
(mostly from ν_e, ν_τ)



Tracks
(from ν_μ)

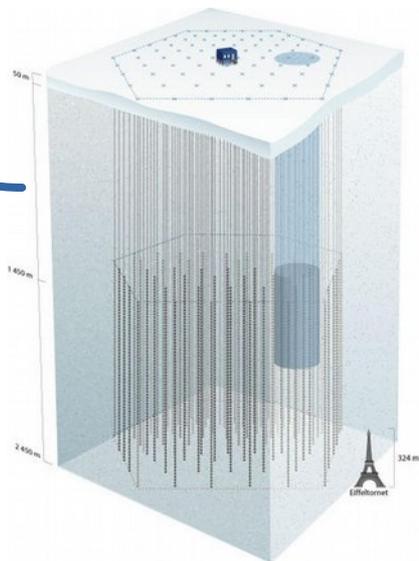


103 contained events, 15 TeV–2 PeV



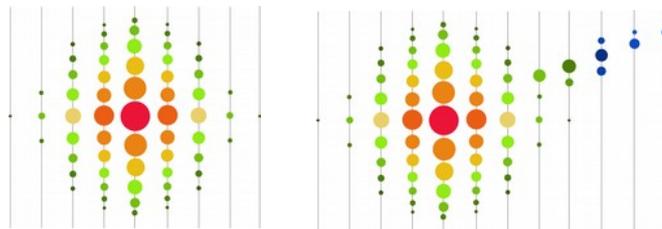
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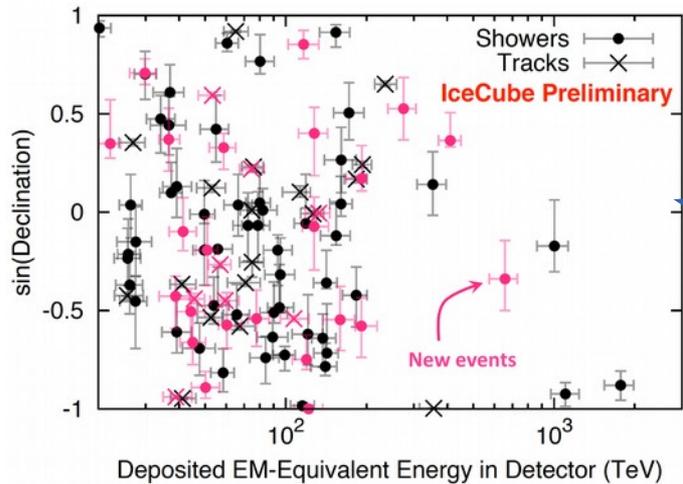


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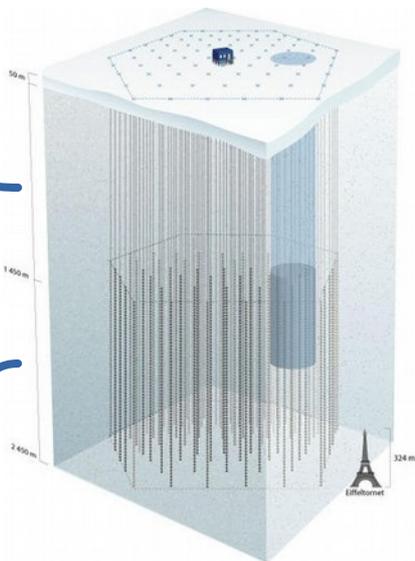


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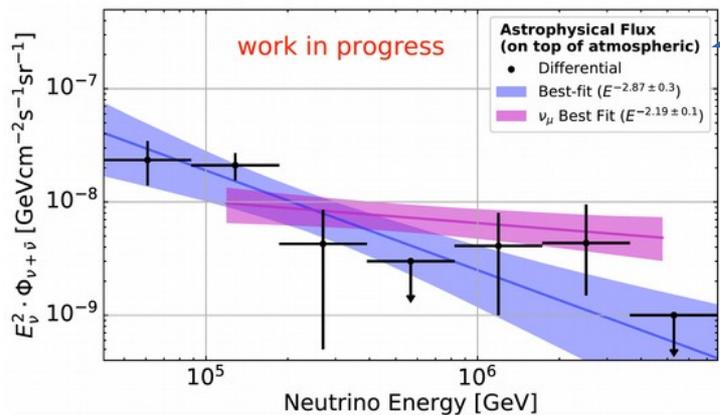


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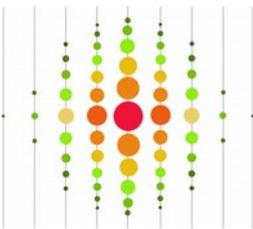
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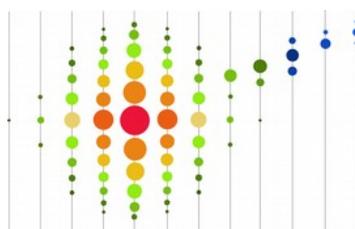
Astrophysical ν flux detected at $> 7\sigma$



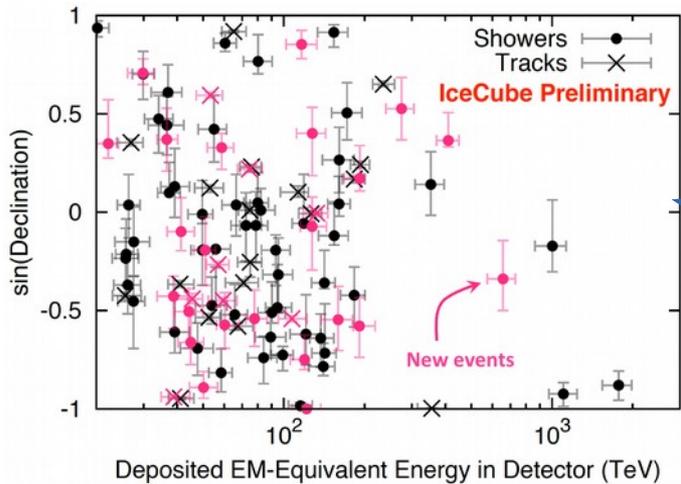
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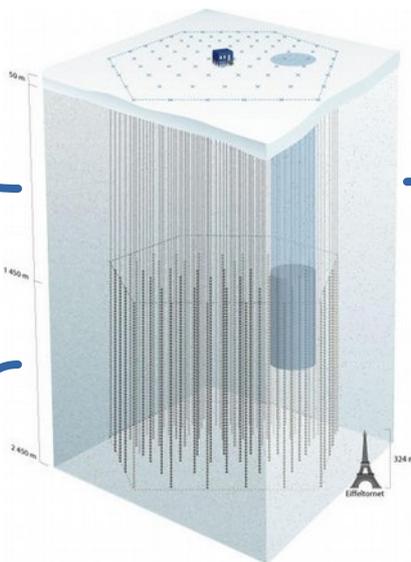


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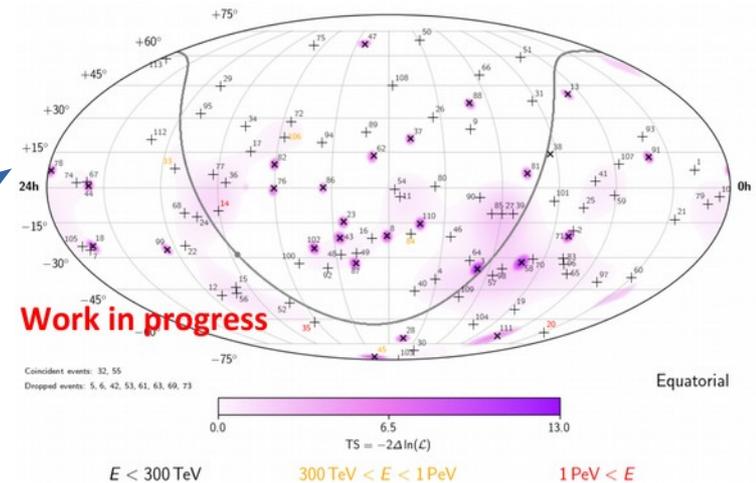


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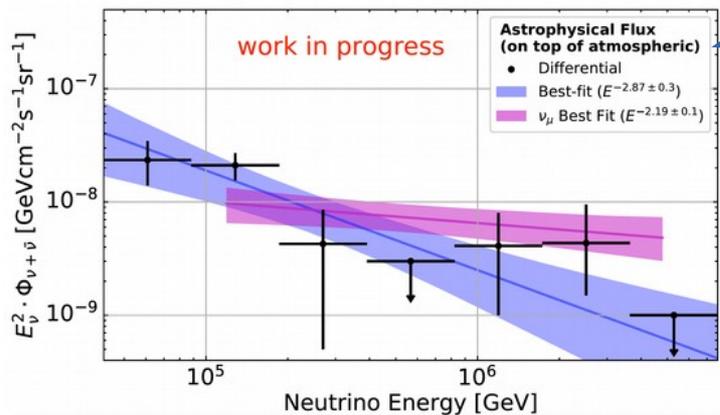
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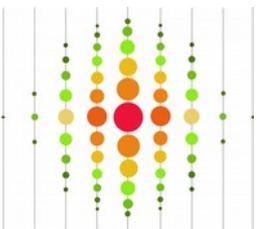
Arrival directions compatible with isotropy



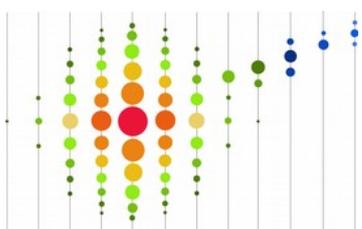
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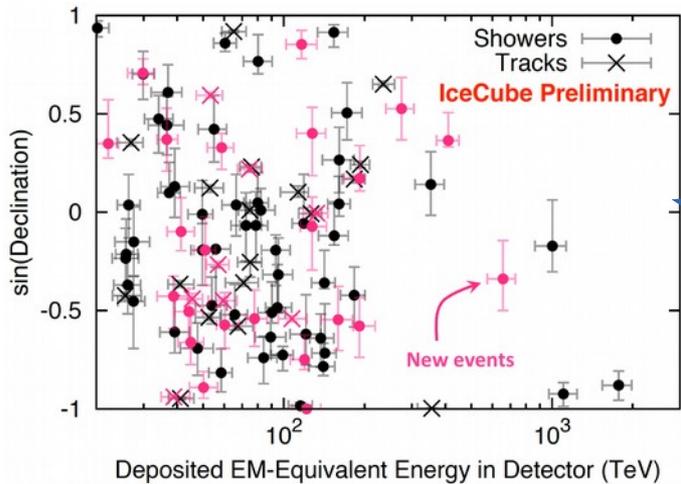
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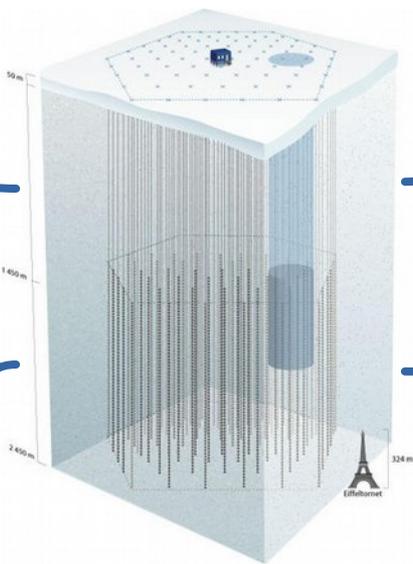


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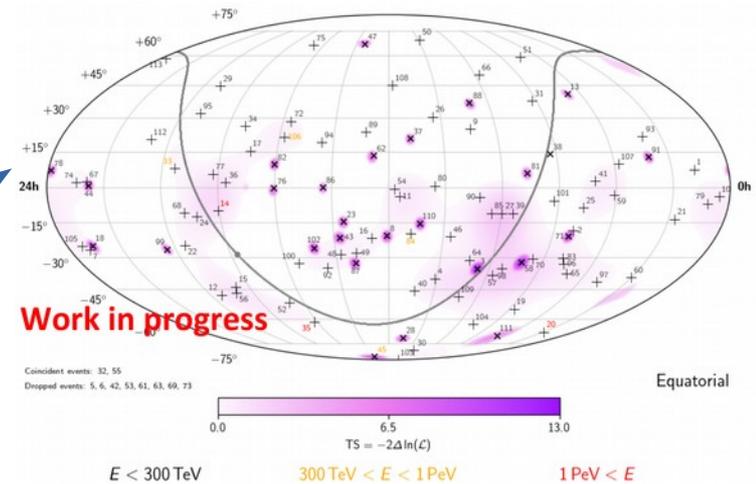


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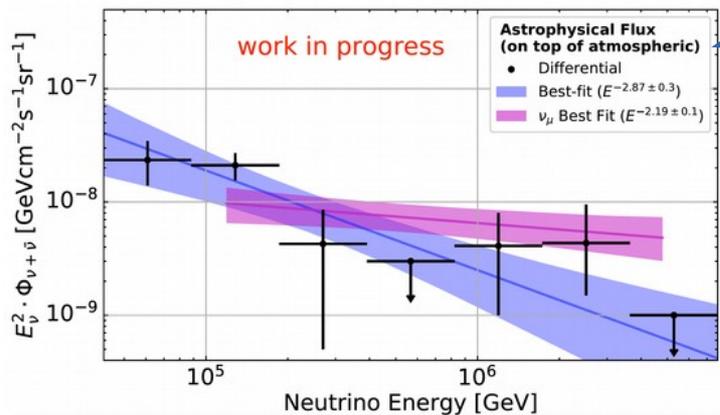
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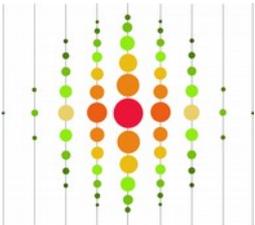
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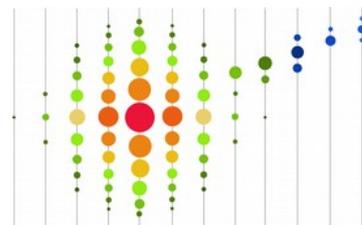
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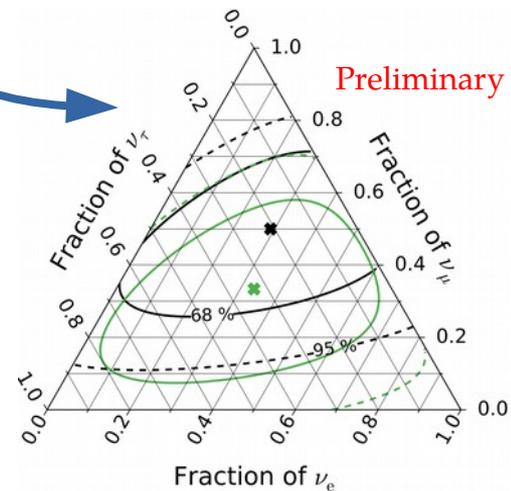
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Flavor composition



Status quo of high-energy cosmic neutrinos

What we know

- ▶ Isotropic distribution of sources
- ▶ Spectrum is a power law $\propto E^{-p}$
- ▶ At least some sources are gamma-ray transients
- ▶ No correlation between directions of cosmic rays and neutrinos
- ▶ Flavor composition: compatible with equal number of ν_e, ν_μ, ν_τ
- ▶ No evident new physics

What we don't know

- ▶ The sources of the diffuse ν flux
- ▶ The ν production mechanism
- ▶ The spectral index of the spectrum
- ▶ A spectral cut-off at a few PeV?
- ▶ Are there Galactic ν sources?
- ▶ The precise flavor composition
- ▶ Is there new physics?

Status quo of high-energy cosmic neutrinos

But we have solid theory expectations
+ fast experimental progress

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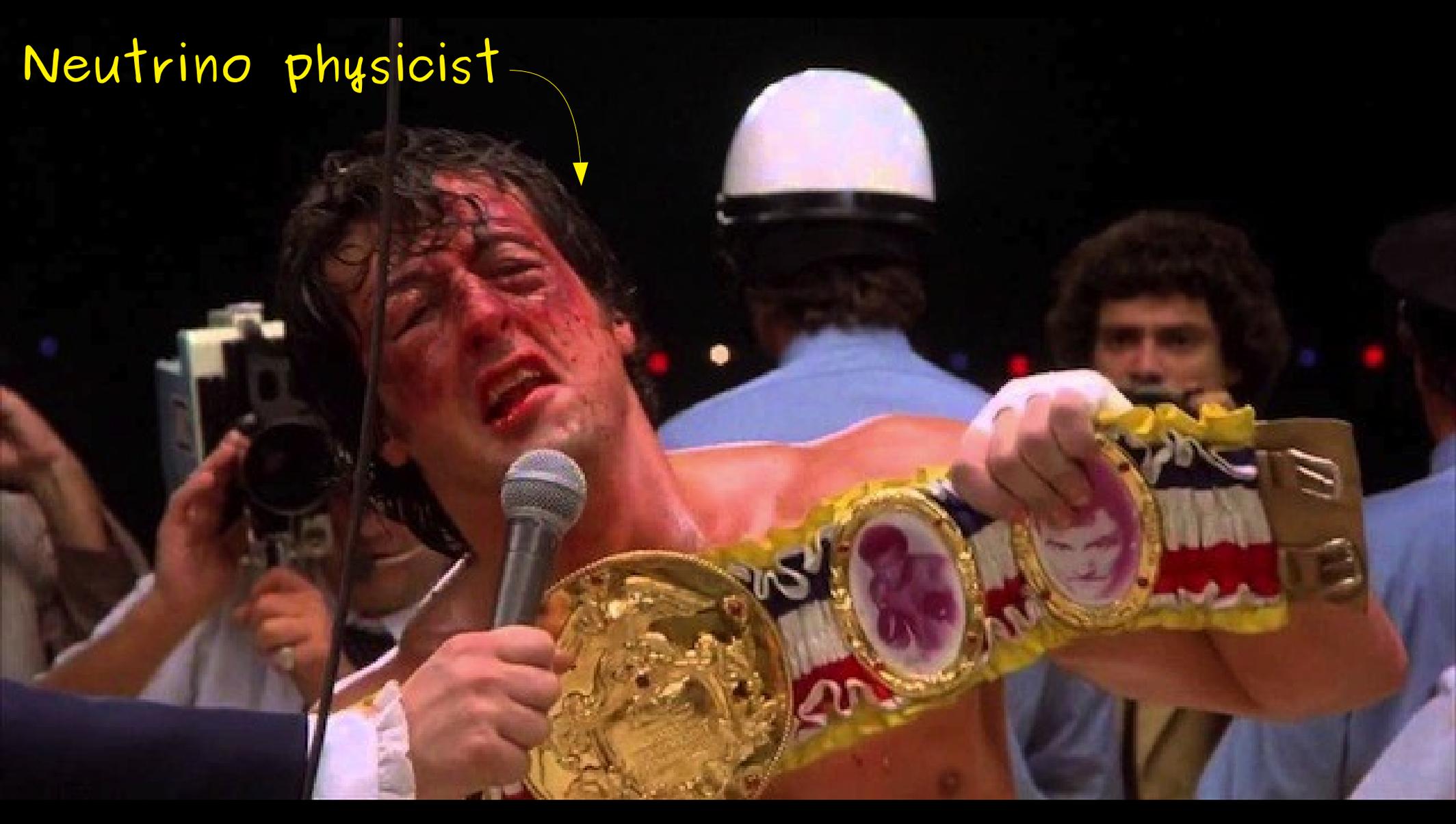
In the face of astrophysical unknowns,
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Yes.

Already today.



Neutrino physicist



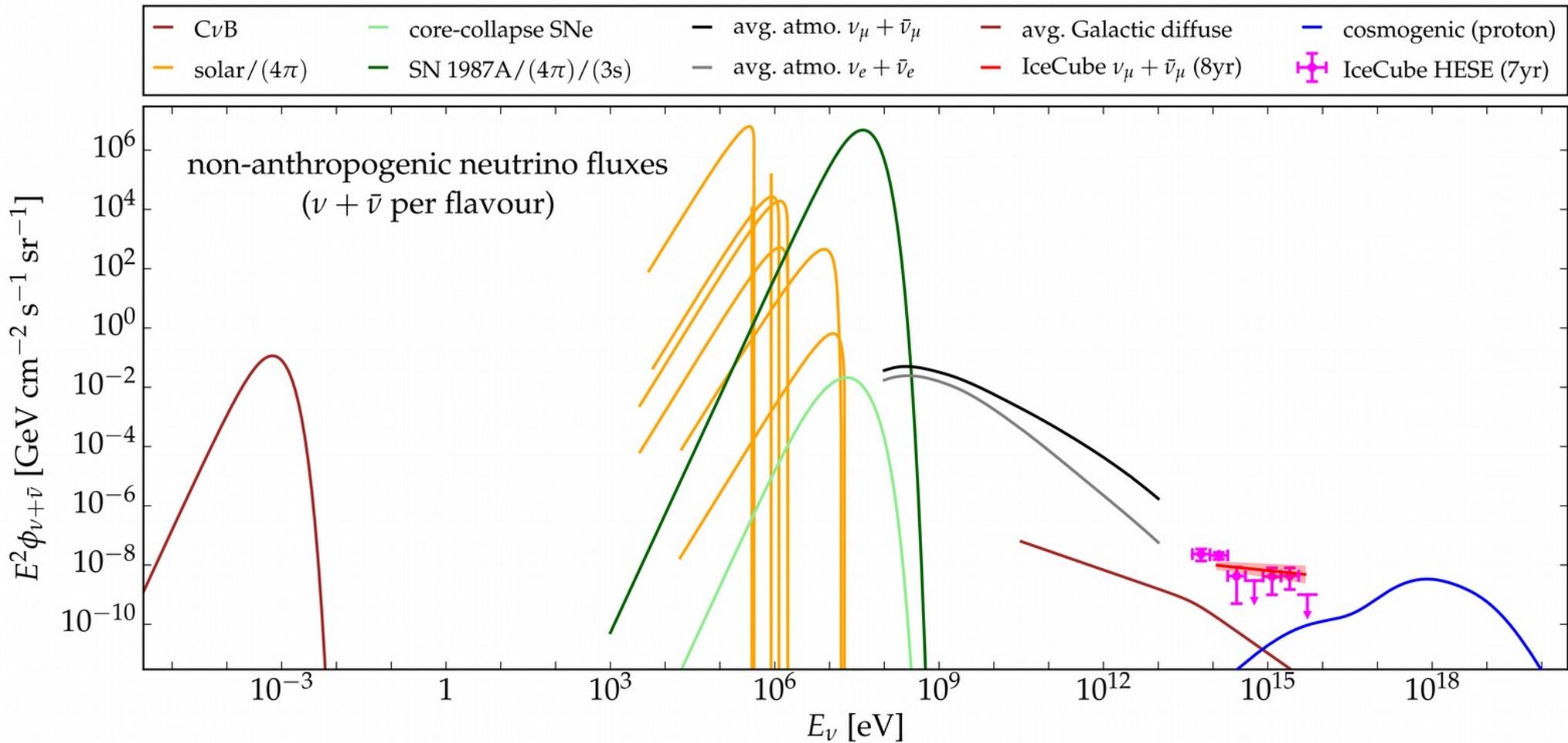


Figure courtesy of Markus Ahlers
 Also in: [Van Elewyck *et al.*, PoS\(ICRC2019\), 1023](#)

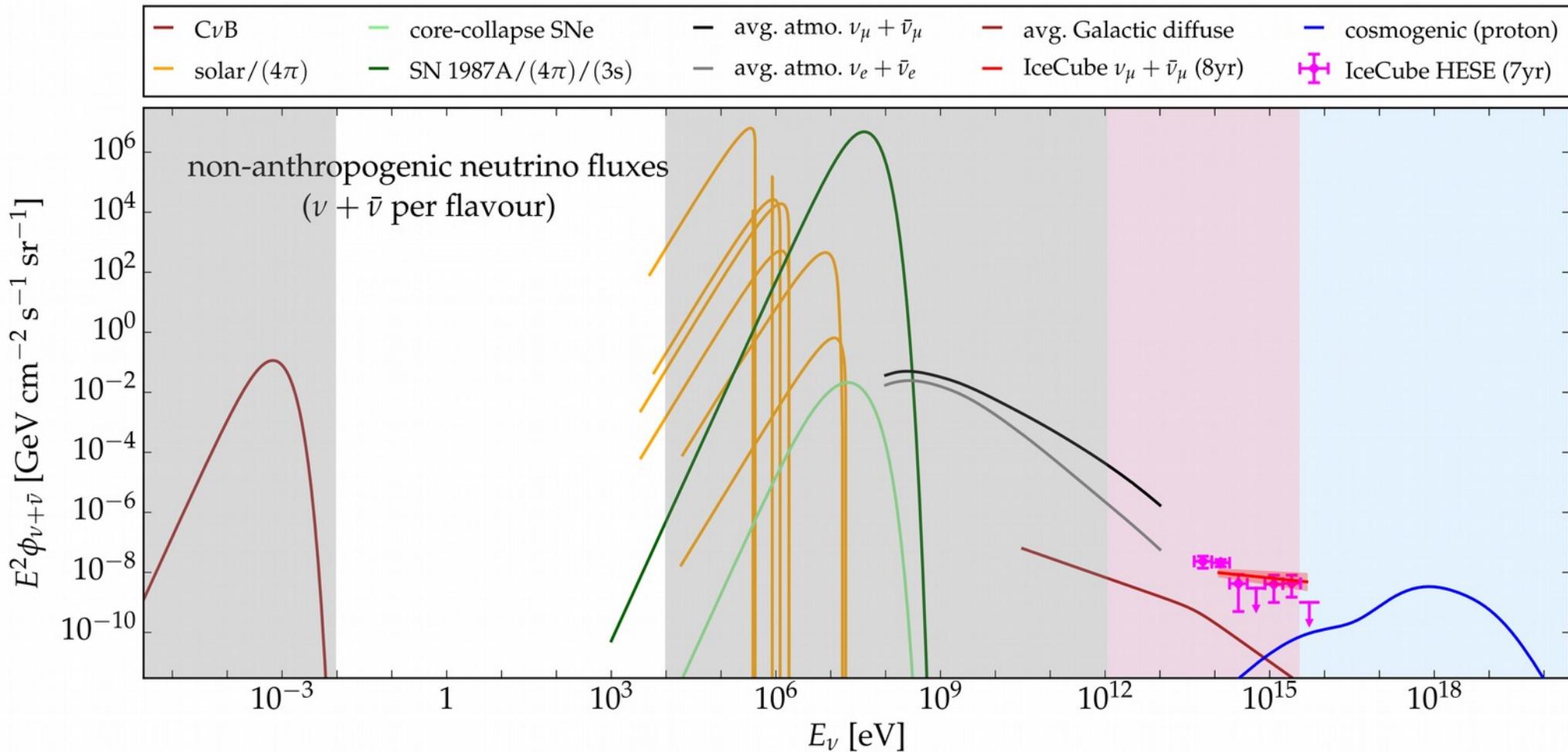


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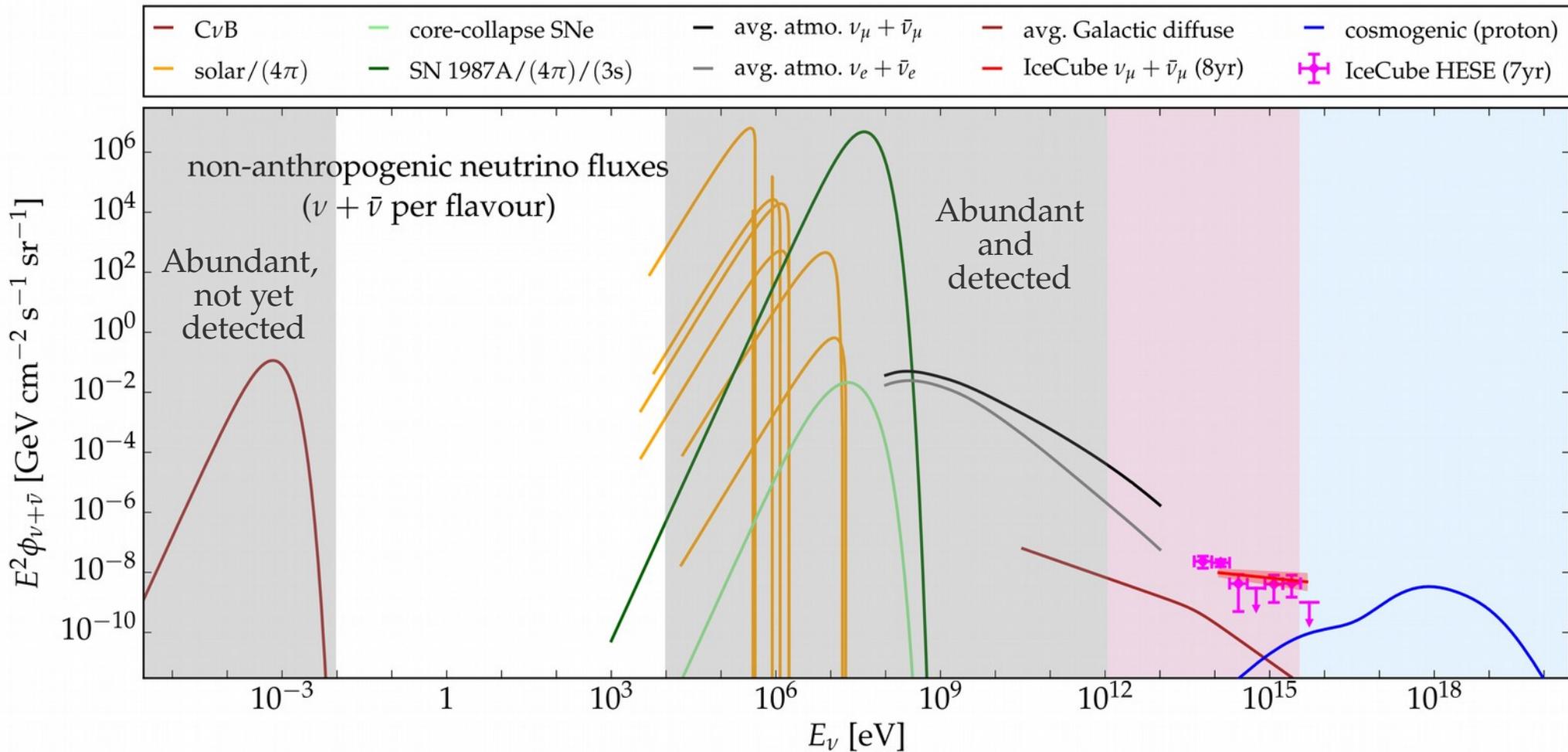


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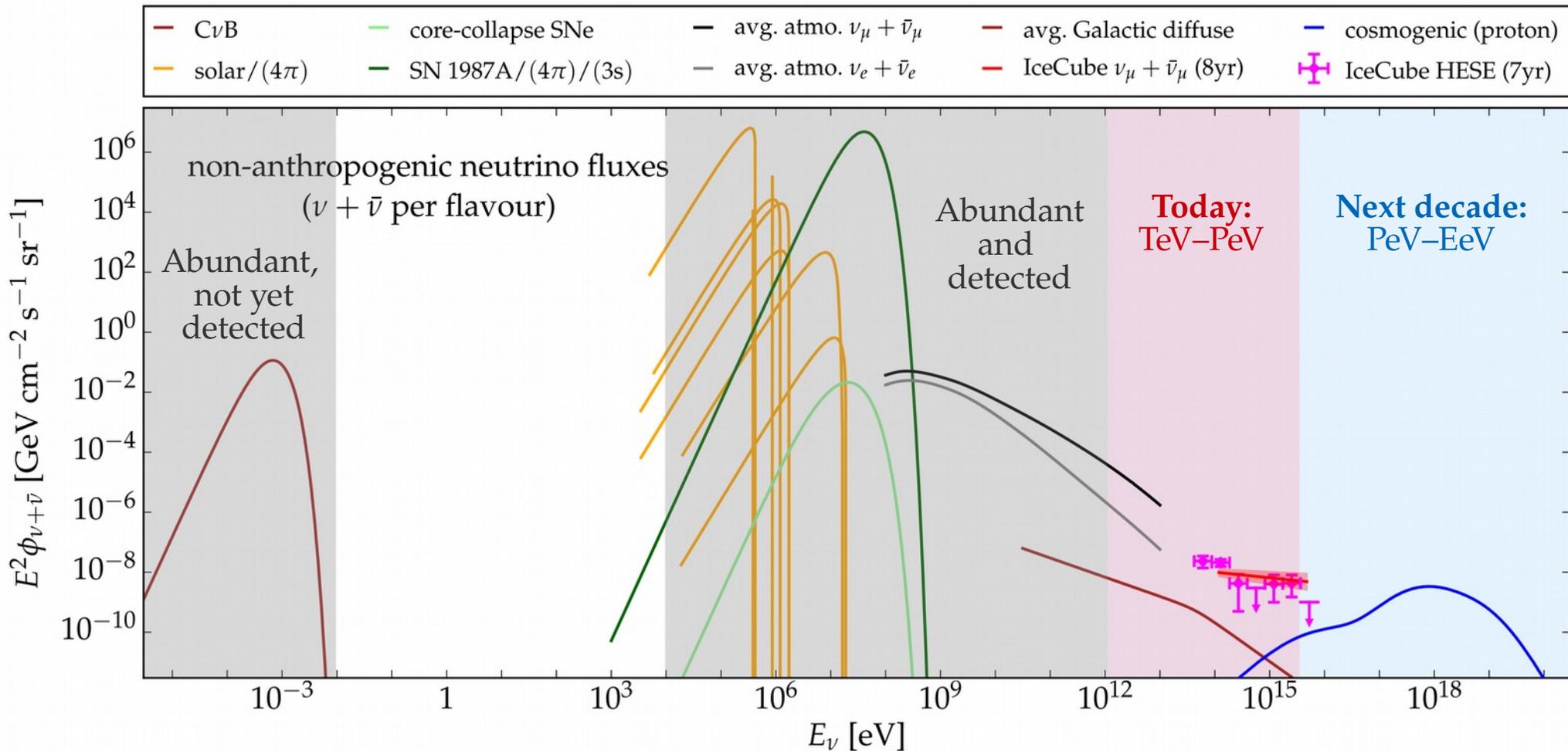
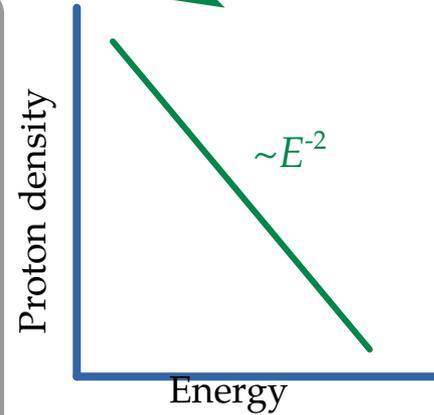
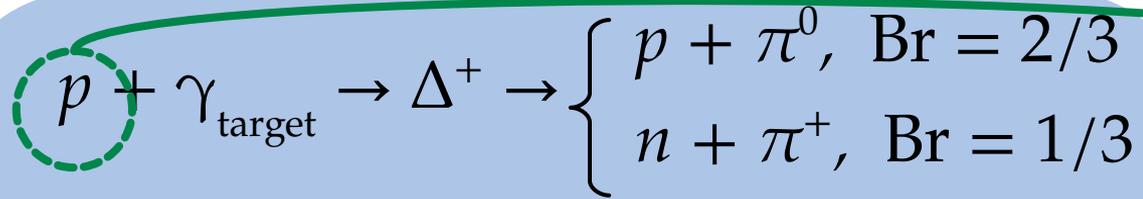


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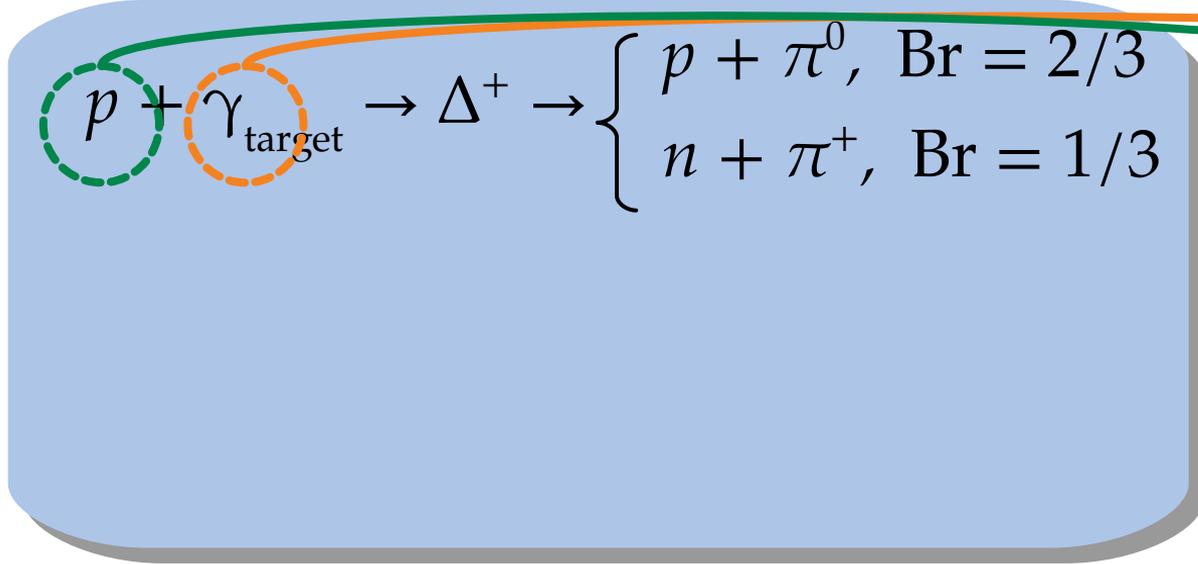
The multi-messenger connection: a simple picture

$$p + \gamma_{\text{target}} \rightarrow \Delta^+ \rightarrow \begin{cases} p + \pi^0, & \text{Br} = 2/3 \\ n + \pi^+, & \text{Br} = 1/3 \end{cases}$$

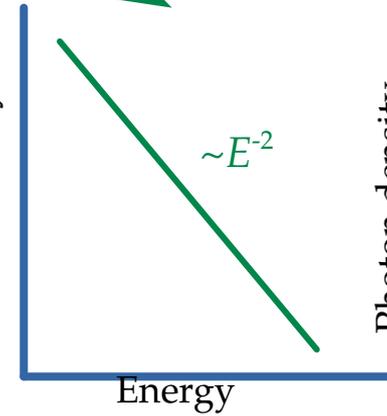
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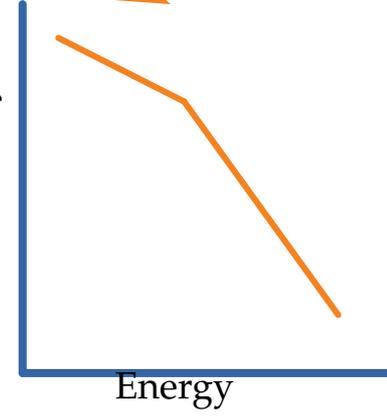
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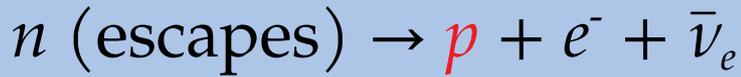
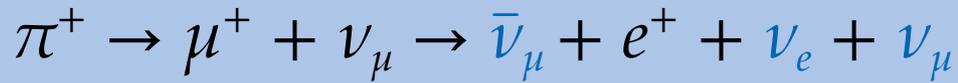
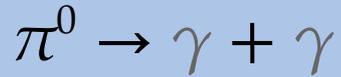
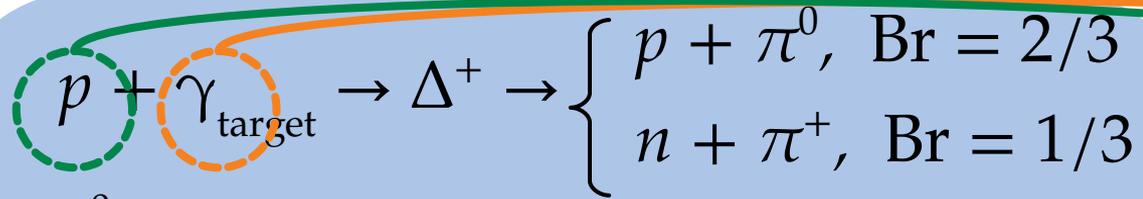
Proton density



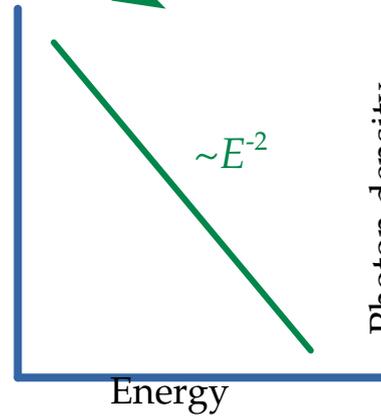
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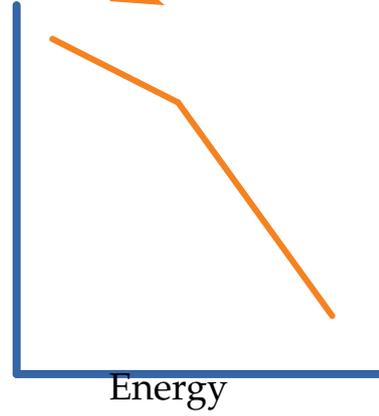
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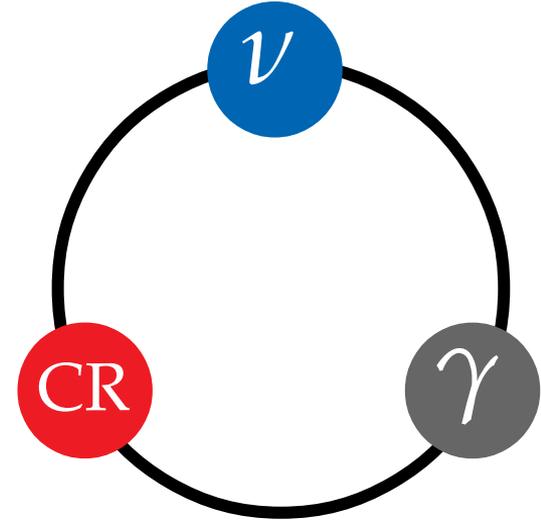
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$$\pi^0 \rightarrow \gamma + \gamma$$

$$\pi^+ \rightarrow \mu^+ + \nu_\mu \rightarrow \bar{\nu}_\mu + e^+ + \nu_e + \nu_\mu$$

$$n \text{ (escapes)} \rightarrow p + e^- + \bar{\nu}_e$$



Neutrino energy = Proton energy / 20

Gamma-ray energy = Proton energy / 10

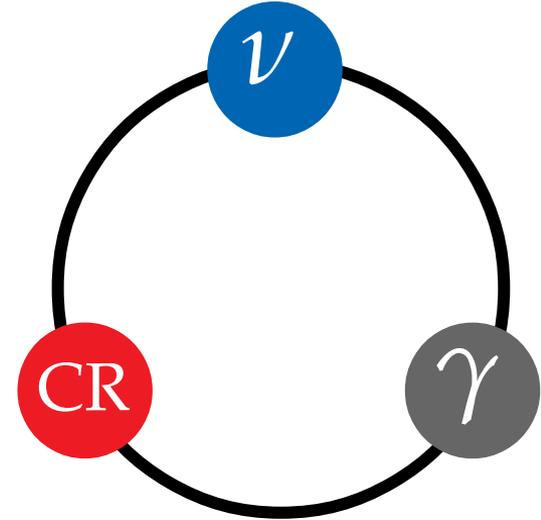
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1 PeV

20 PeV

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Neutrinos – The ultimate smoking gun of cosmic accelerators

Gamma rays

Neutrinos

UHE Cosmic rays

Point back at sources

Size of horizon

Energy degradation

Relative ease to detect

Note: This is a simplified view

Neutrinos – The ultimate smoking gun of cosmic accelerators

	Gamma rays	Neutrinos	UHE Cosmic rays
Point back at sources	Yes	Yes	No
Size of horizon			
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Neutrinos – The ultimate smoking gun of cosmic accelerators

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How does IceCube see TeV–PeV neutrinos?

Deep inelastic neutrino-nucleon scattering

Neutral current (NC)

$$\nu_x + N \rightarrow \nu_x + X$$

Charged current (CC)

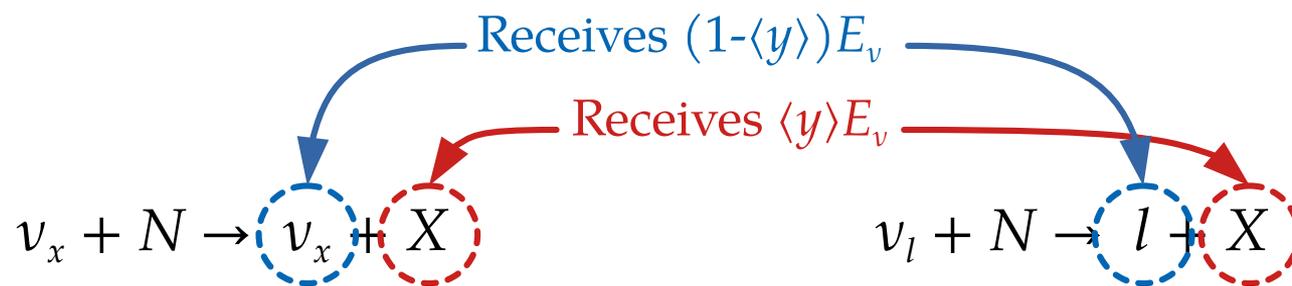
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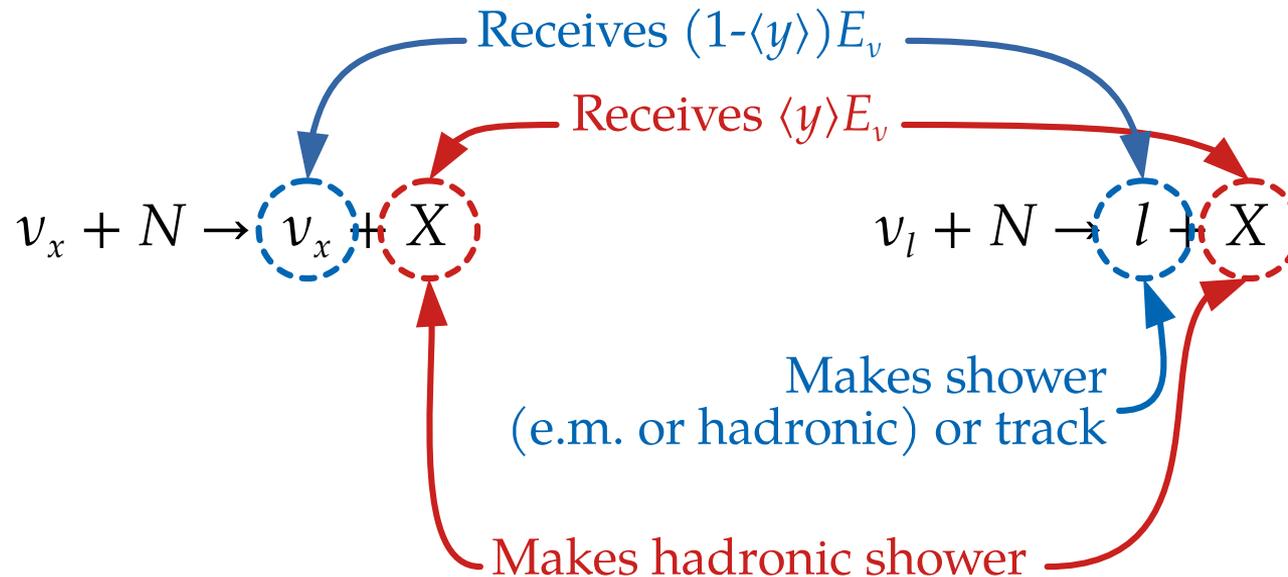
At TeV–PeV, the average inelasticity $\langle y \rangle = 0.25\text{--}0.30$

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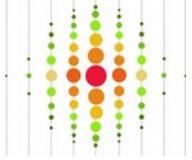
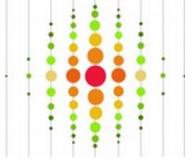
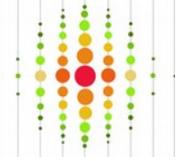
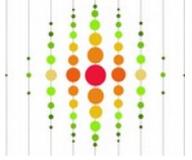
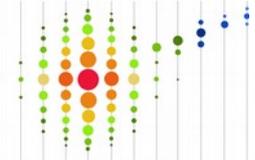
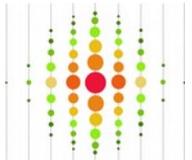
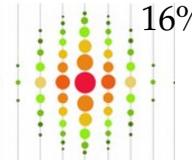
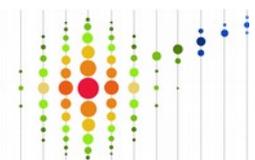
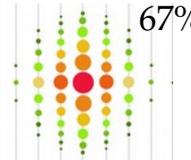
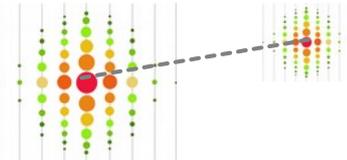
Charged current (CC)



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Detected

To be confirmed

$\nu_x + \bar{\nu}_x$ NC	 <p>Hadronic X shower</p>				
$\nu_e + \bar{\nu}_e$ CC	 <p>Hadronic X shower</p>	+  <p>E.m. shower</p>			
$\nu_\mu + \bar{\nu}_\mu$ CC	 <p>Hadronic X shower</p>	+  <p>Track</p>			
$\nu_\tau + \bar{\nu}_\tau$ CC	 <p>Hadronic X shower</p>	+  <p>E.m. shower</p>	or  <p>Track</p>	or  <p>Hadronic shower</p>	 <p>Double pulse/bang</p>

Fundamental physics with HE cosmic neutrinos

- ▶ Numerous new-physics effects grow as $\sim \kappa_n \cdot E^n \cdot L$
- ▶ So we can probe $\kappa_n \sim 4 \cdot 10^{-47} (E/\text{PeV})^{-n} (L/\text{Gpc})^{-1} \text{PeV}^{1-n}$
- ▶ Improvement over current limits: $\kappa_0 < 10^{-29} \text{PeV}$, $\kappa_1 < 10^{-33}$
- ▶ Fundamental physics can be extracted from four neutrino observables:
 - ▶ Spectral shape
 - ▶ Angular distribution
 - ▶ Flavor composition
 - ▶ Timing

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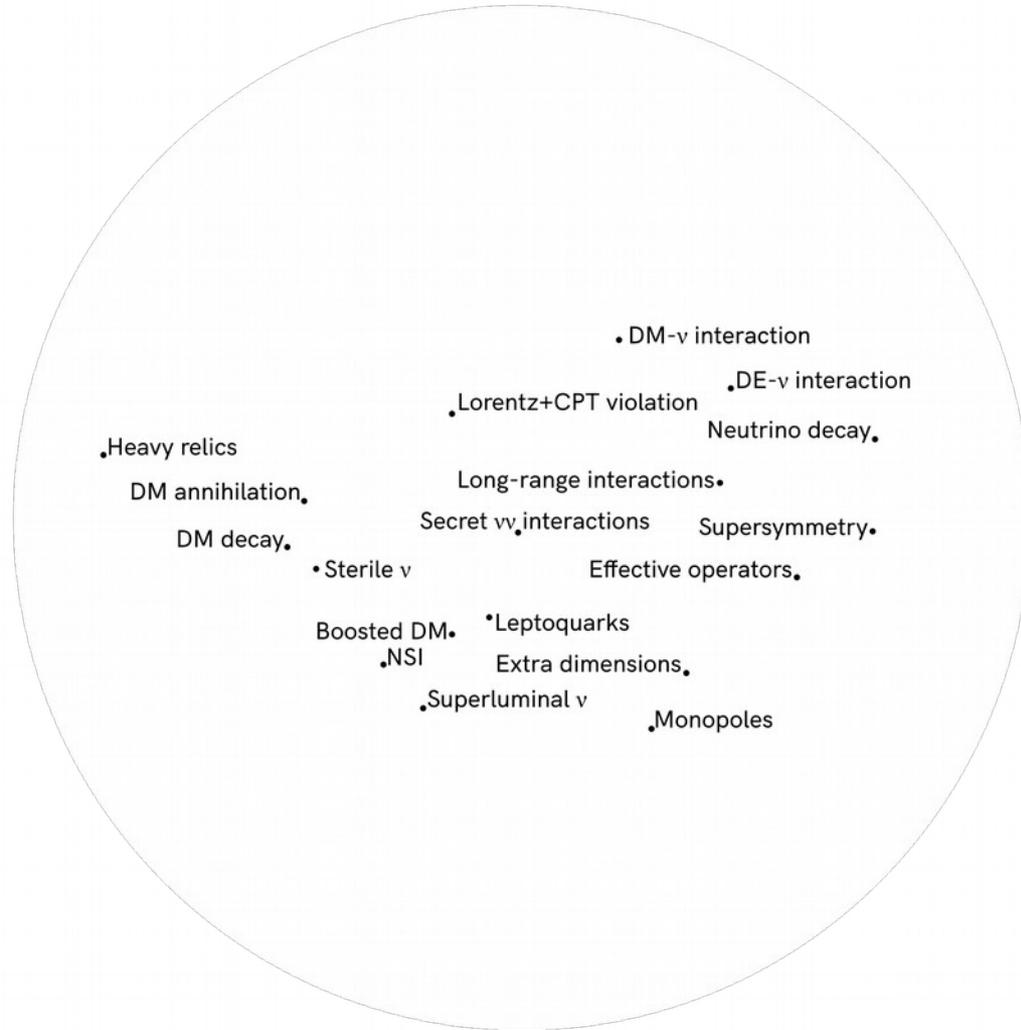
▶ So we can probe $\kappa_n \sim 4 \cdot 10^{-47} (E/\text{PeV})^{-n} (L/\text{Gpc})^{-1} \text{PeV}^{1-n}$

▶ Improvement over current limits: $\kappa_0 < 10^{-29} \text{PeV}$, $\kappa_1 < 10^{-33}$

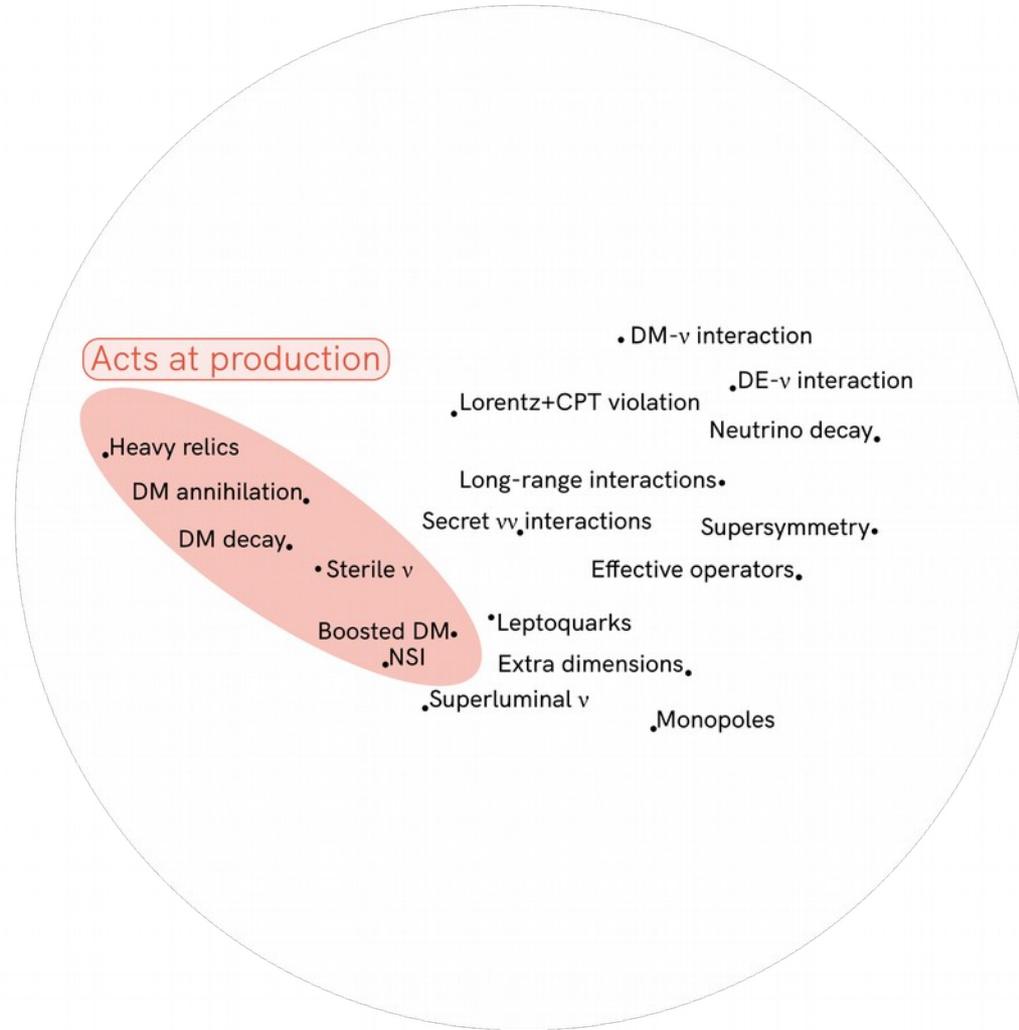
▶ Fundamental physics can be extracted from four neutrino observables:

- ▶ Spectral shape
- ▶ Angular distribution
- ▶ Flavor composition
- ▶ Timing

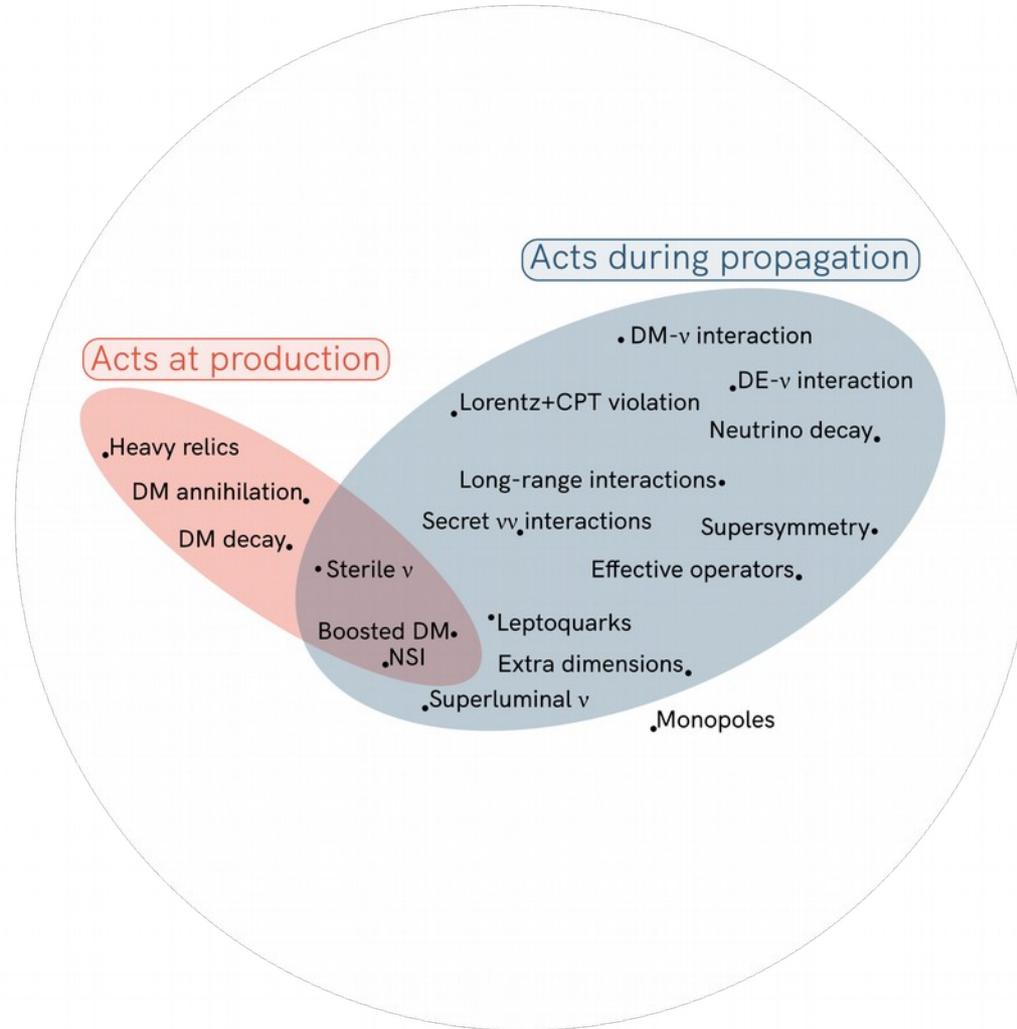
*In spite of
poor energy, angular, flavor reconstruction
& astrophysical unknowns*



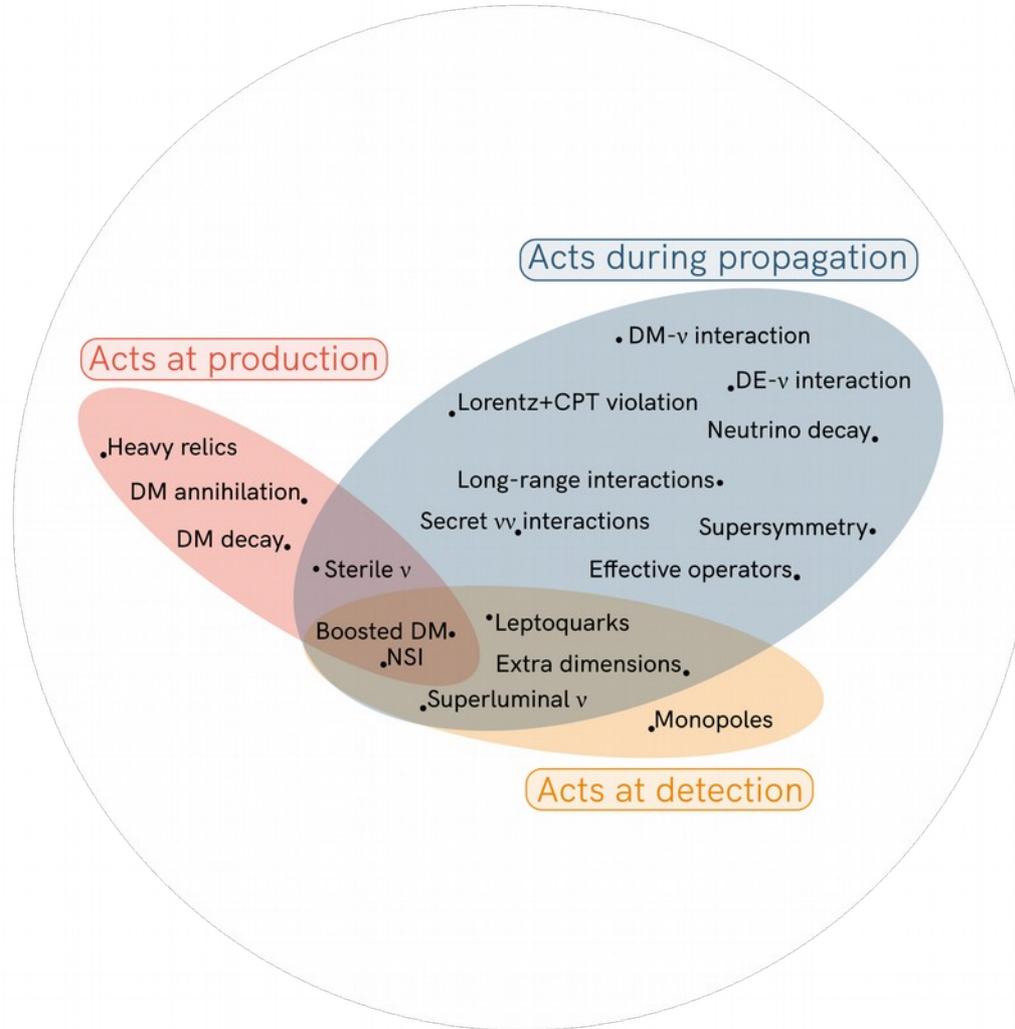
Note: Not an exhaustive list



Note: Not an exhaustive list



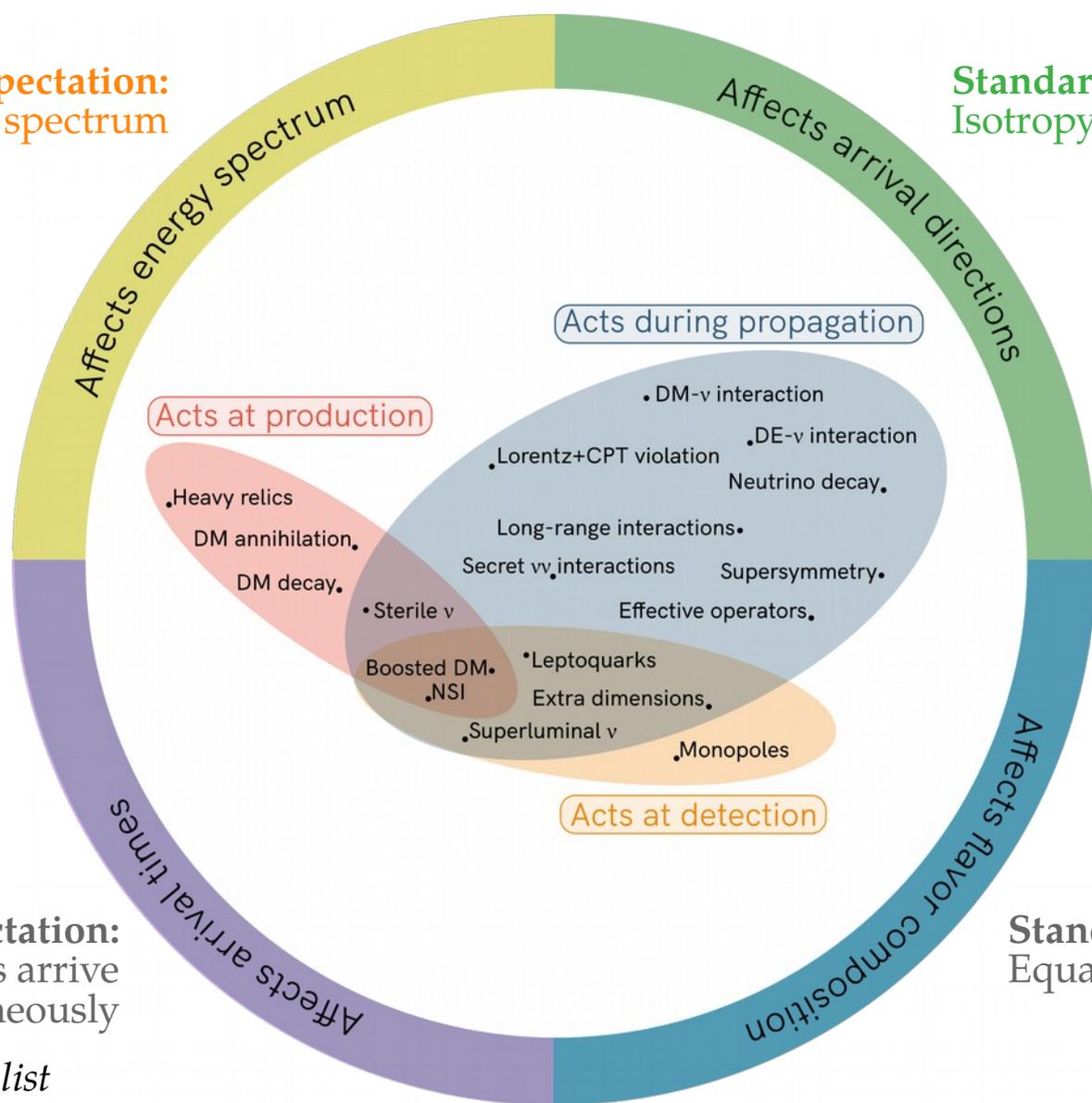
Note: Not an exhaustive list



Note: Not an exhaustive list

Standard expectation:
Power-law energy spectrum

Standard expectation:
Isotropy (for diffuse flux)



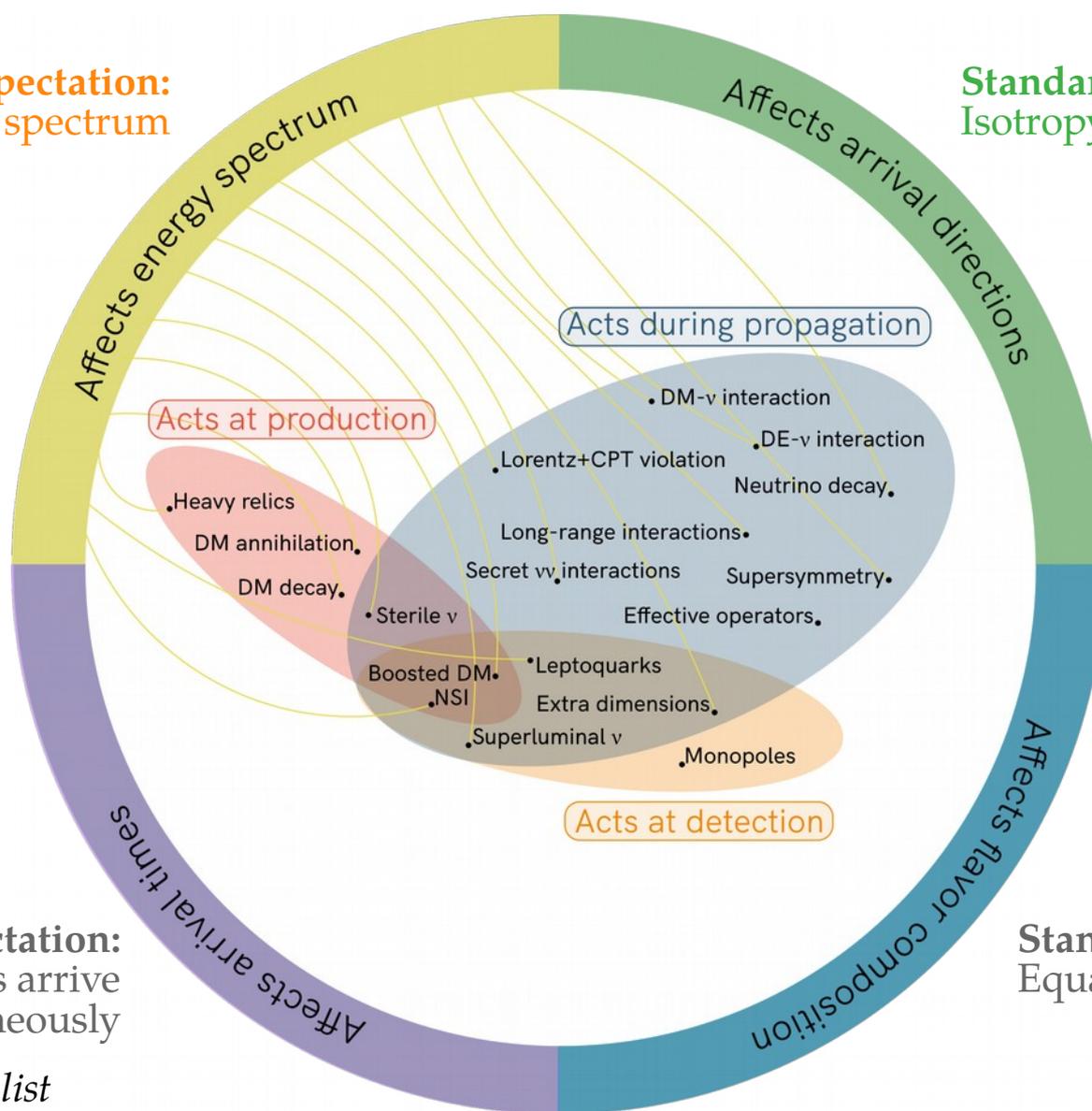
Standard expectation:
 ν and γ from transients arrive simultaneously

Standard expectation:
Equal number of ν_e, ν_μ, ν_τ

Note: Not an exhaustive list

Standard expectation:
Power-law energy spectrum

Standard expectation:
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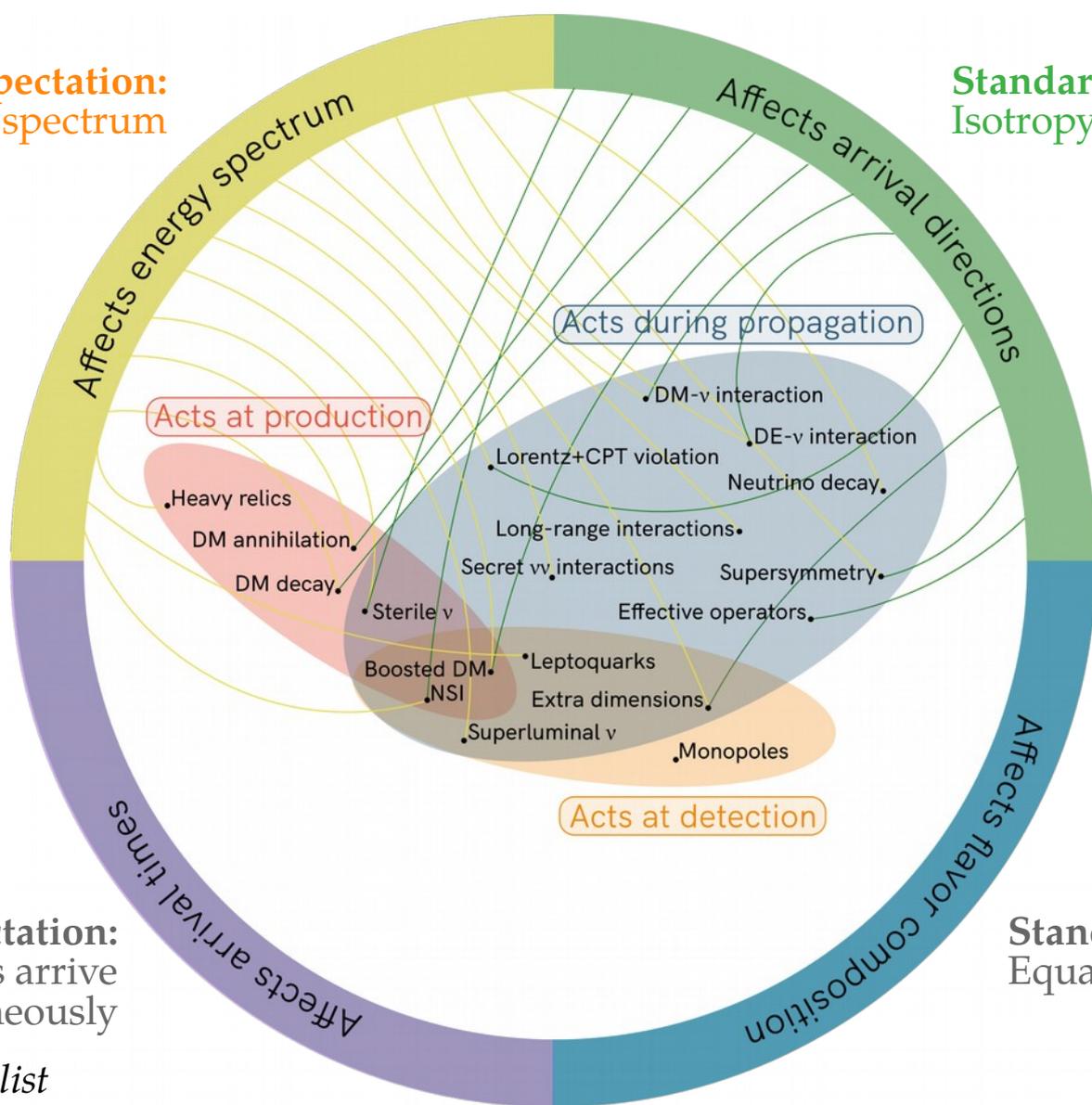
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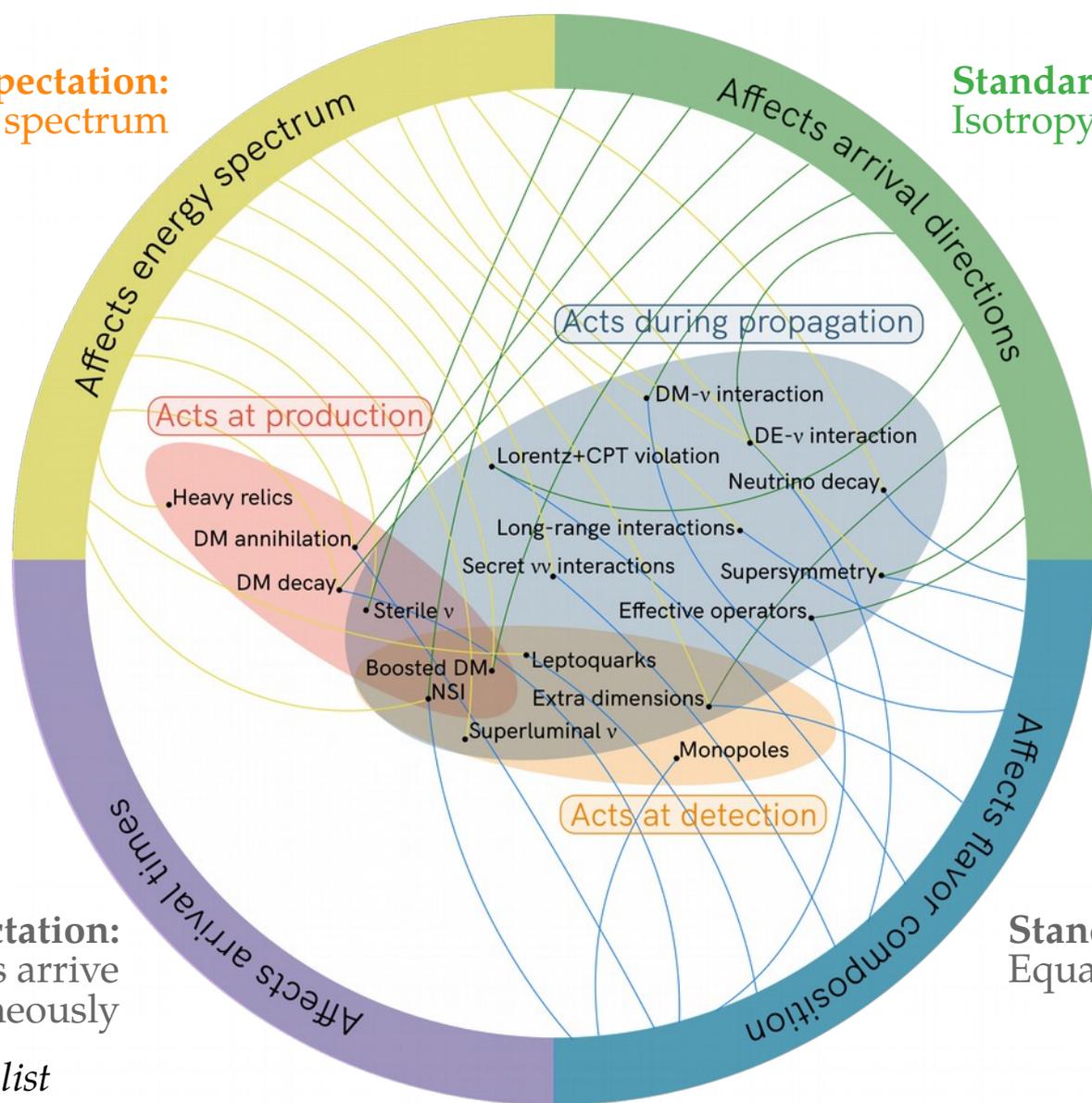
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Note: Not an exhaustive list

Standard expectation:
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Isotropy (for diffuse flux)



Acts at production

Acts during propagation

Acts at detection

Affects energy spectrum

Affects arrival directions

Affects arrival times

Affects flavor composition

Heavy relics
DM annihilation
DM decay
Sterile ν
Boosted DM
NSI
Superluminal ν

Lorentz+CPT violation
Long-range interactions
Secret $\nu\nu$ interactions
Leptoquarks
Extra dimensions
Monopoles

DM- ν interaction
DE- ν interaction
Neutrino decay
Supersymmetry
Effective operators

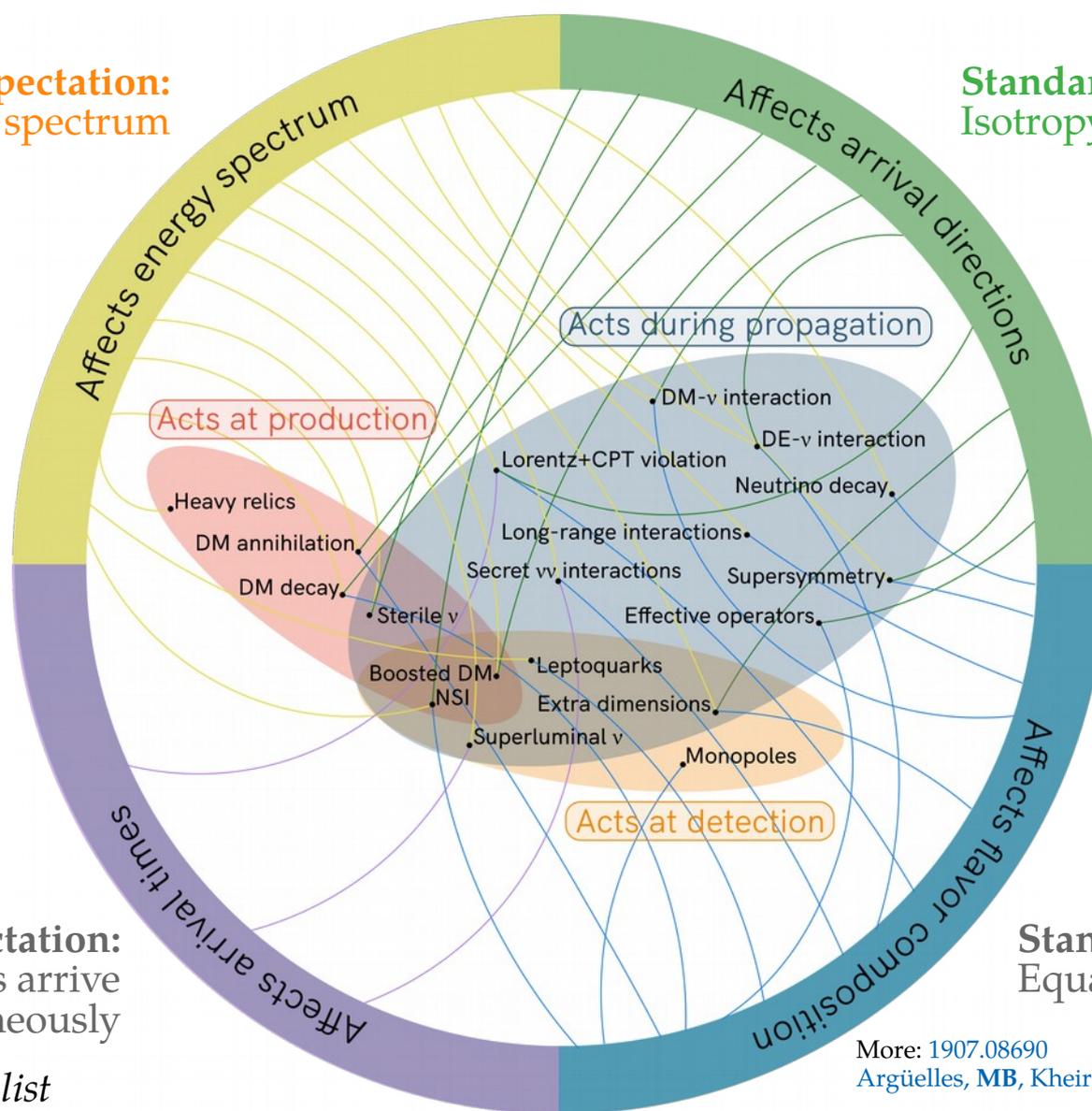
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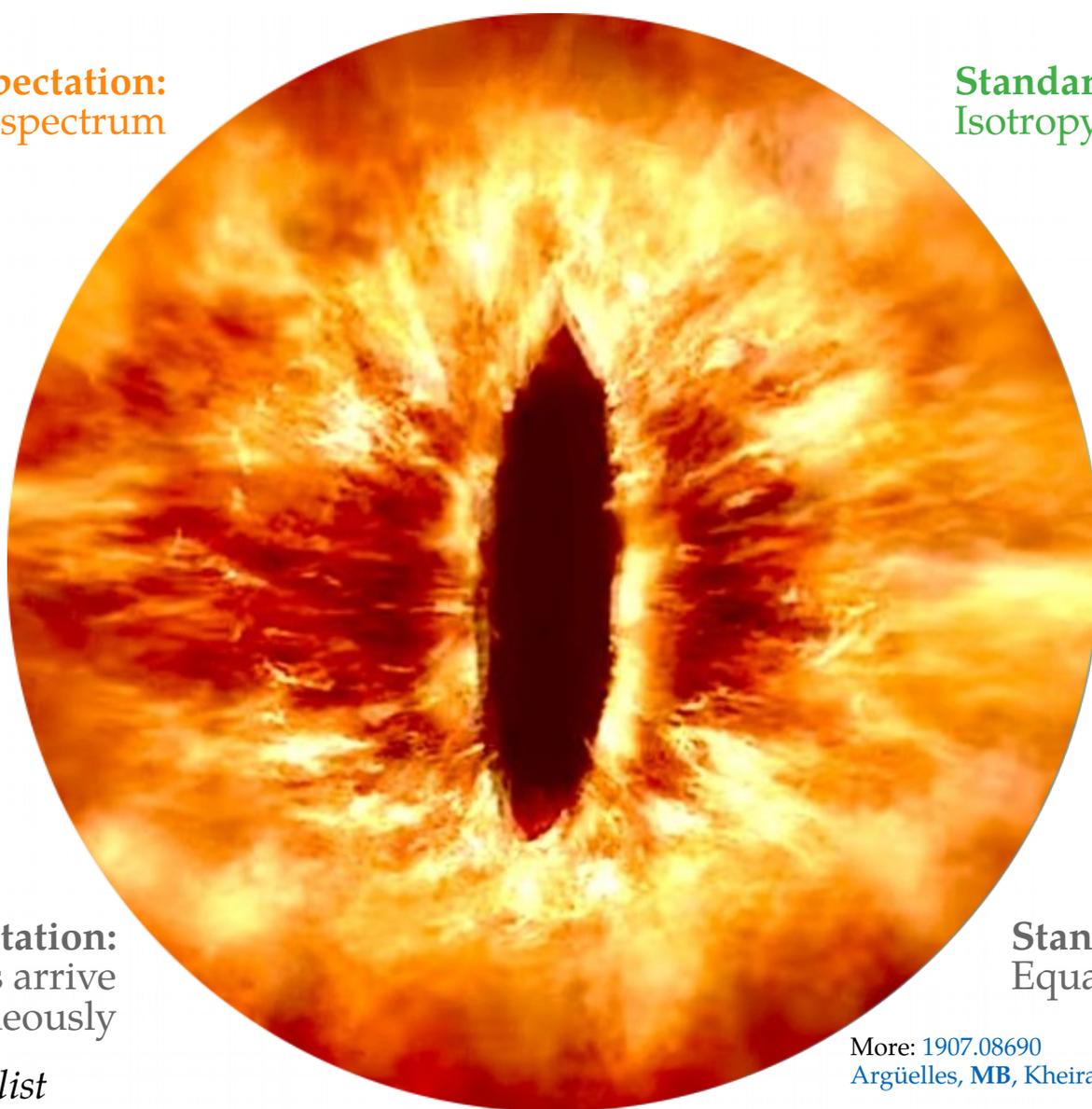
Note: Not an exhaustive list

More: 1907.08690

Argüelles, MB, Kheirandish, Palomares-Ruiz, Salvadó, Vincent

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Power-law energy spectrum

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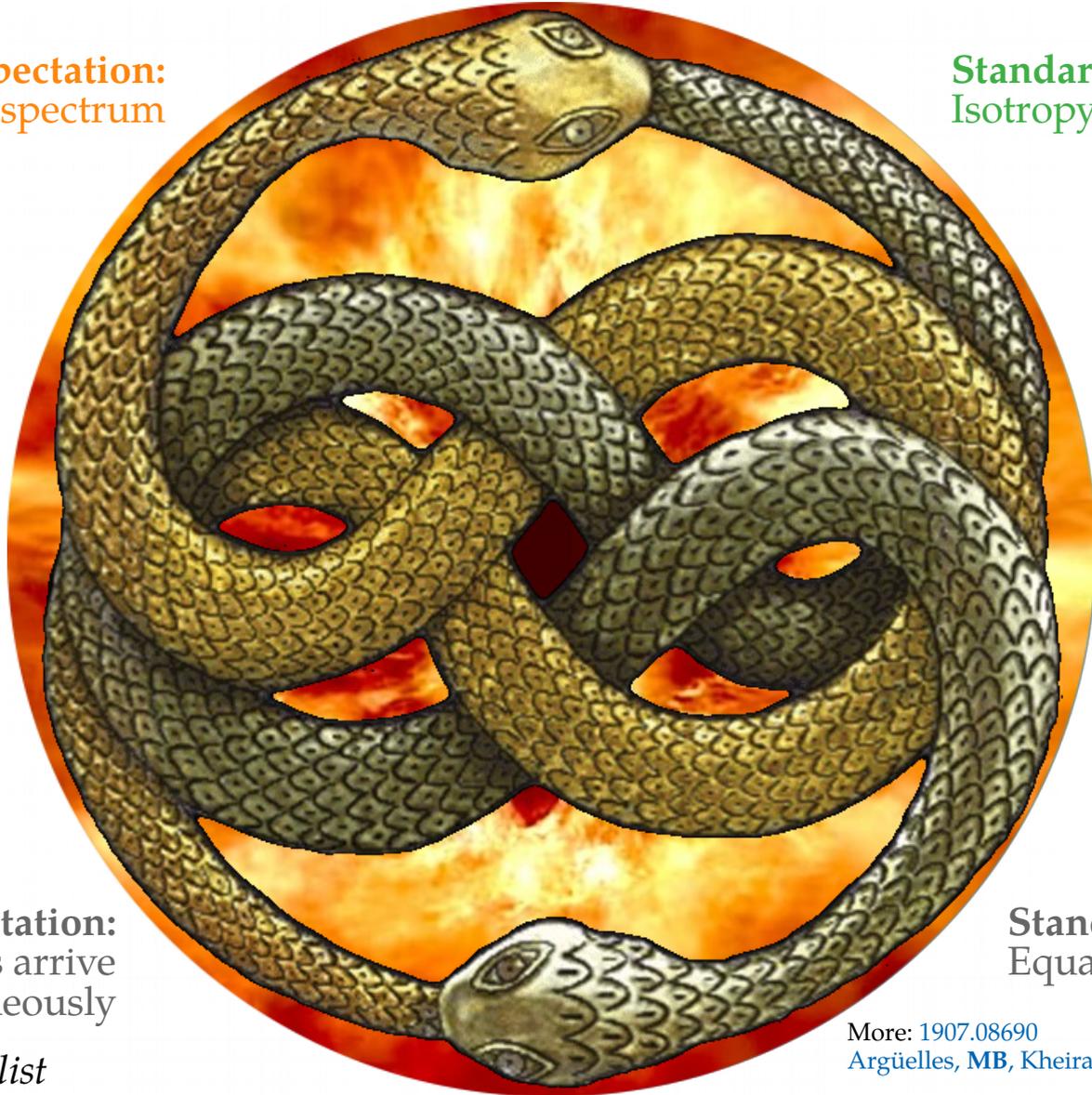
Standard expectation:
 ν and γ from transients arrive
simultaneously

Standard expectation:
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Note: Not an exhaustive list

More: 1907.08690
Argüelles, MB, Kheirandish, Palomares-Ruiz, Salvadó, Vincent

Standard expectation:
Power-law energy spectrum



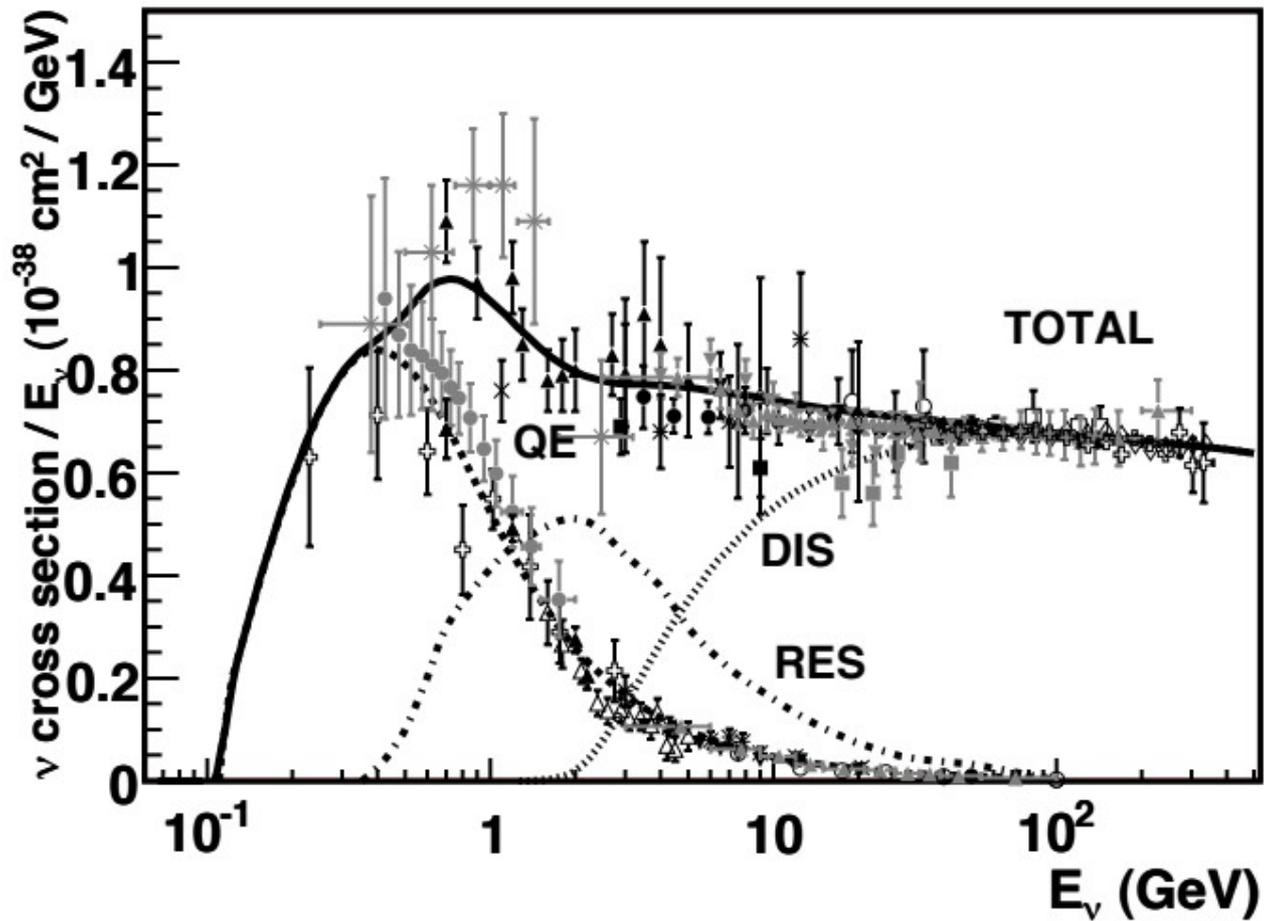
Standard expectation:
Isotropy (for diffuse flux)

Standard expectation:
 ν and γ from transients arrive
simultaneously

Standard expectation:
Equal number of ν_e, ν_μ, ν_τ

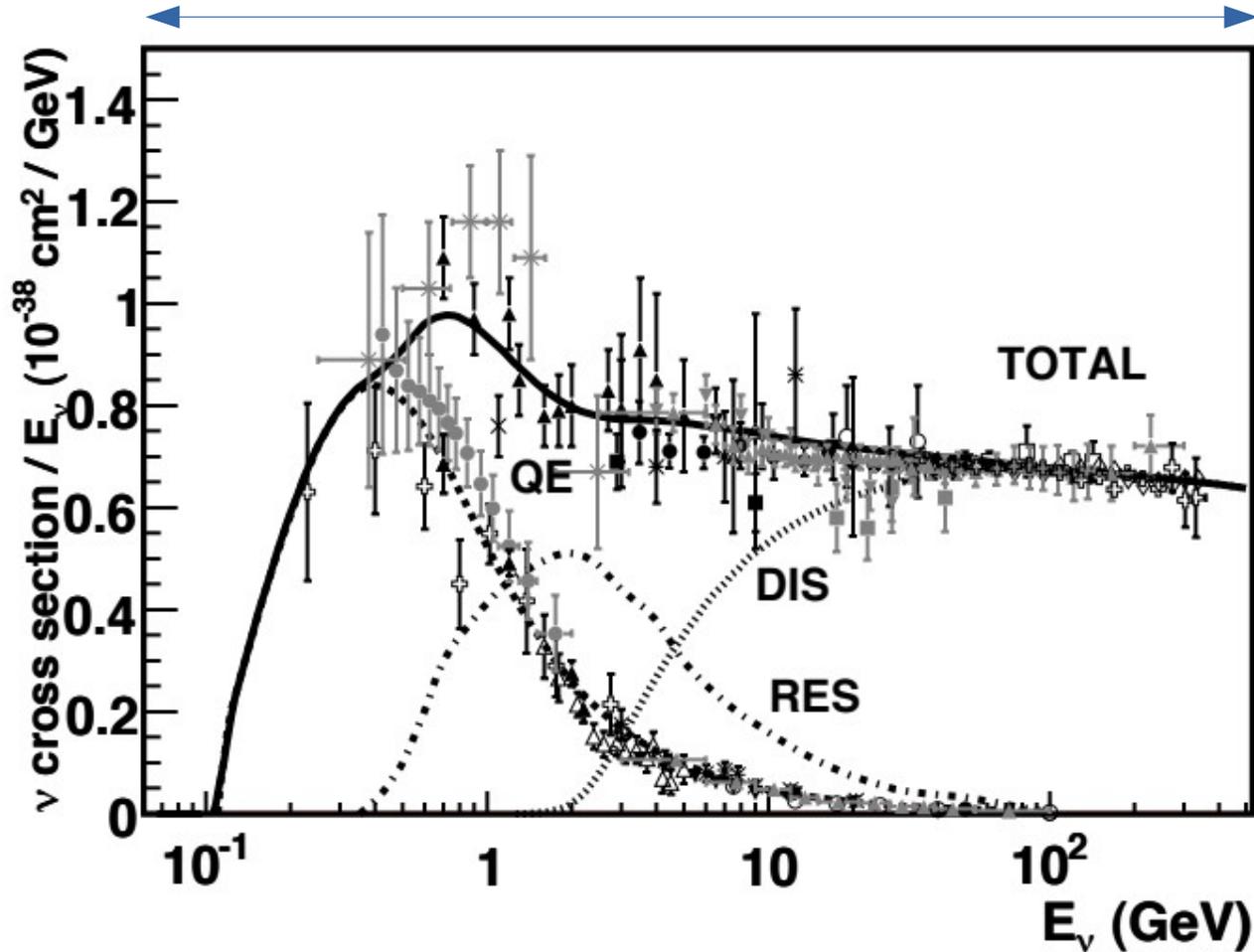
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Particle Data Group

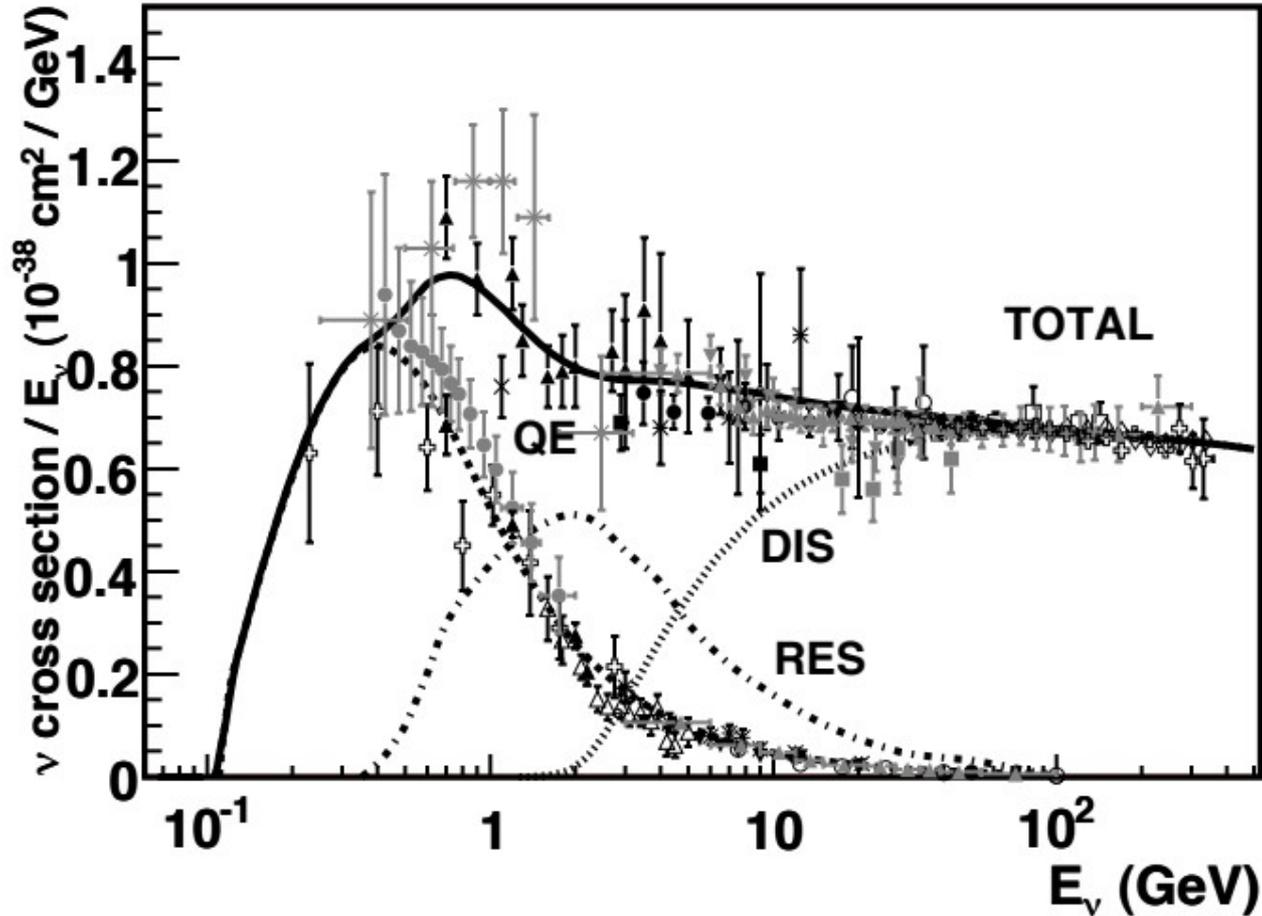
Accelerator experiments



Particle Data Group

Accelerator experiments

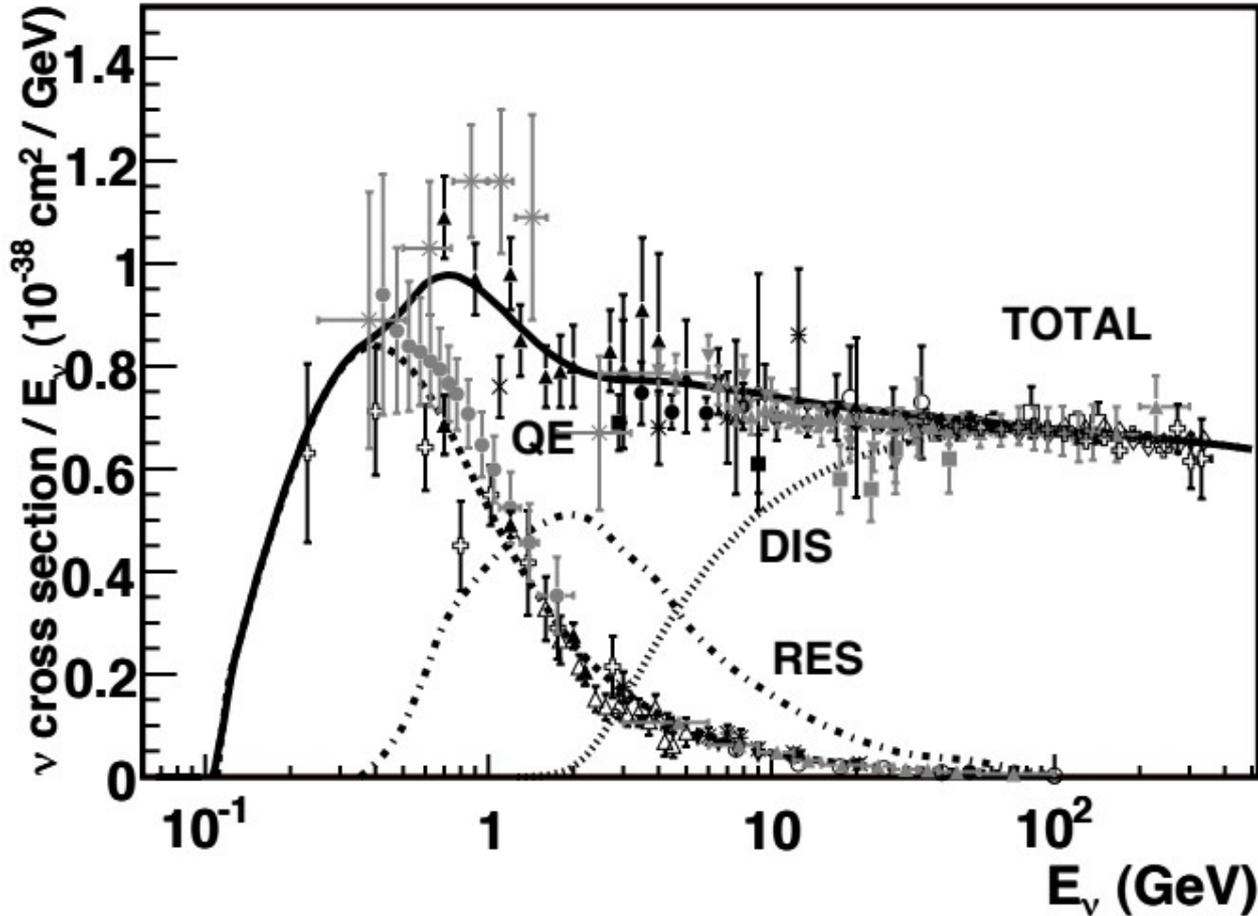
←
One recent
measurement
(COHERENT)



Particle Data Group

Accelerator experiments

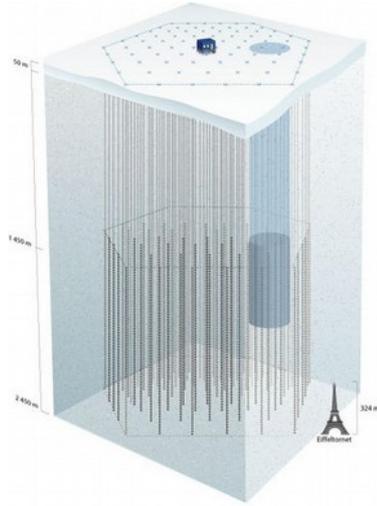
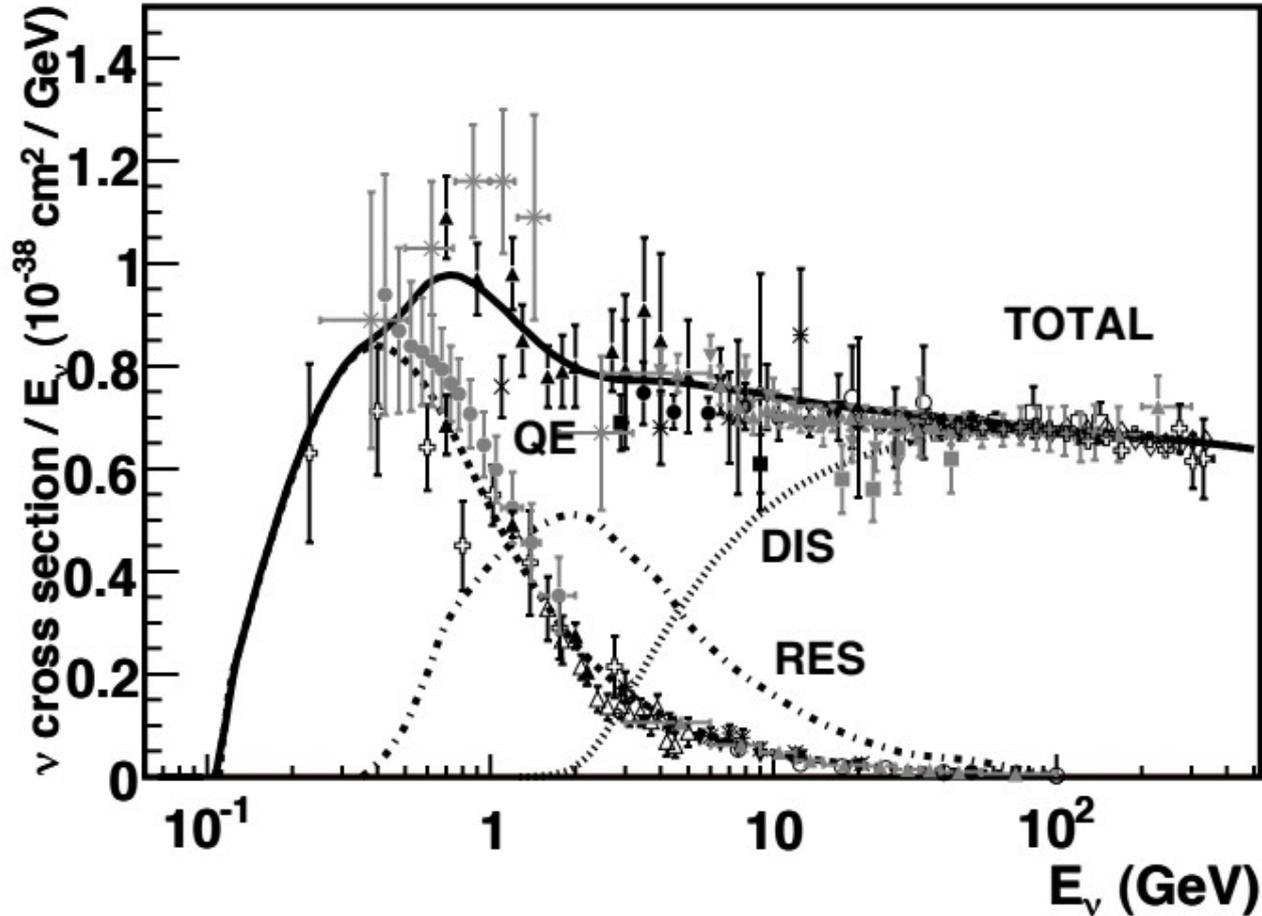
←
One recent
measurement
(COHERENT)



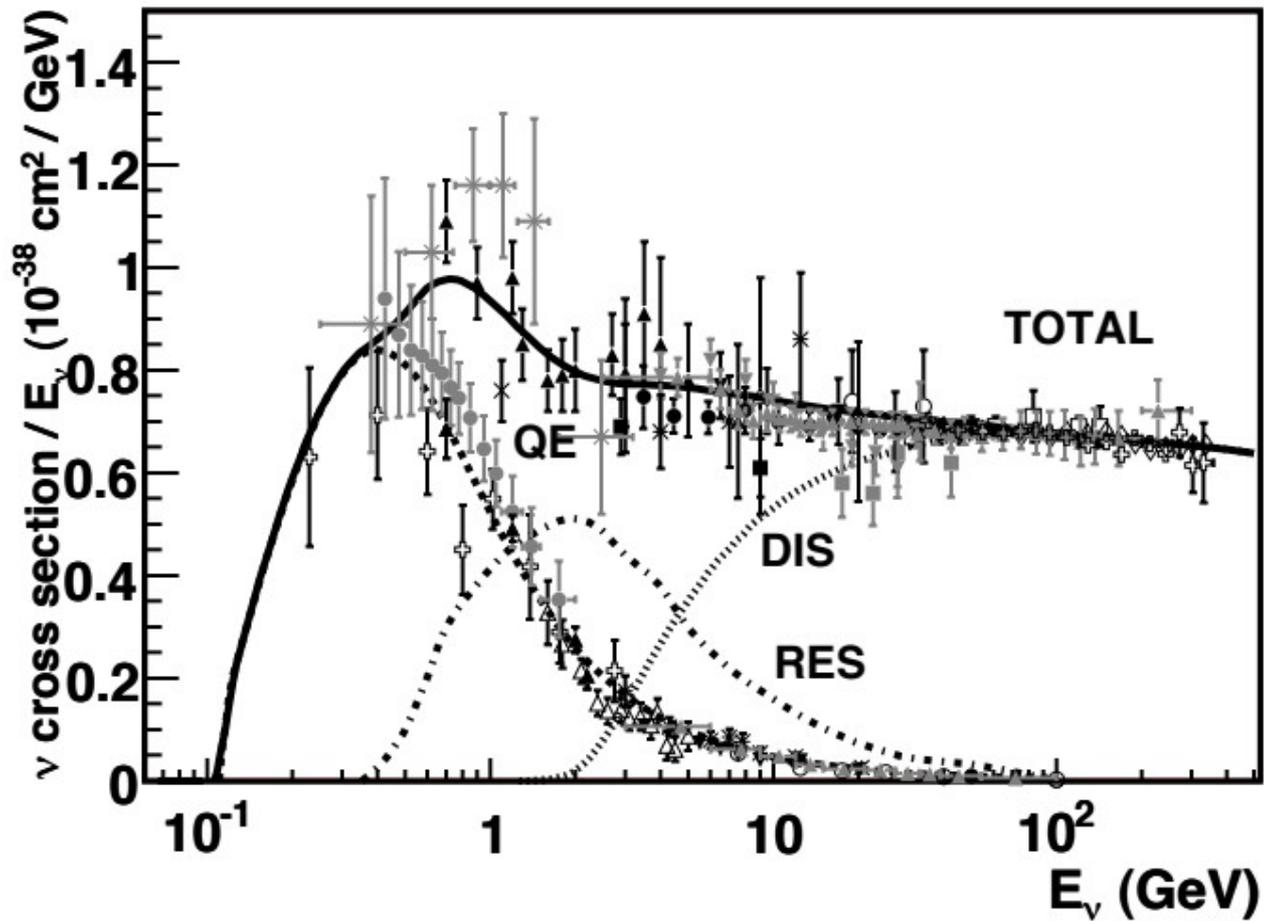
→
No
measurements
... until recently!

Accelerator experiments

← One recent measurement (COHERENT)

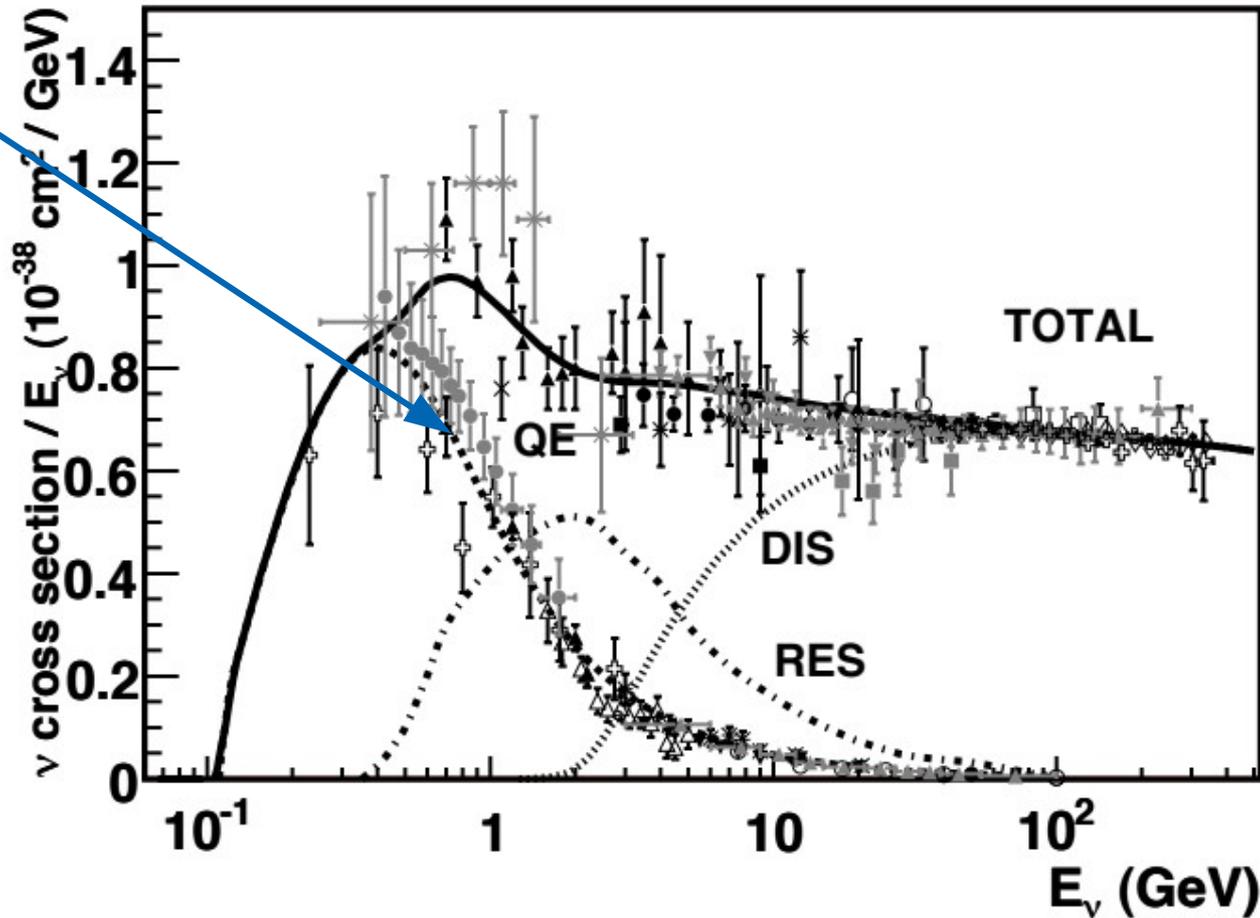
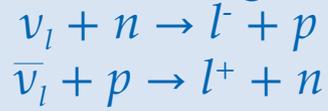


Particle Data Group



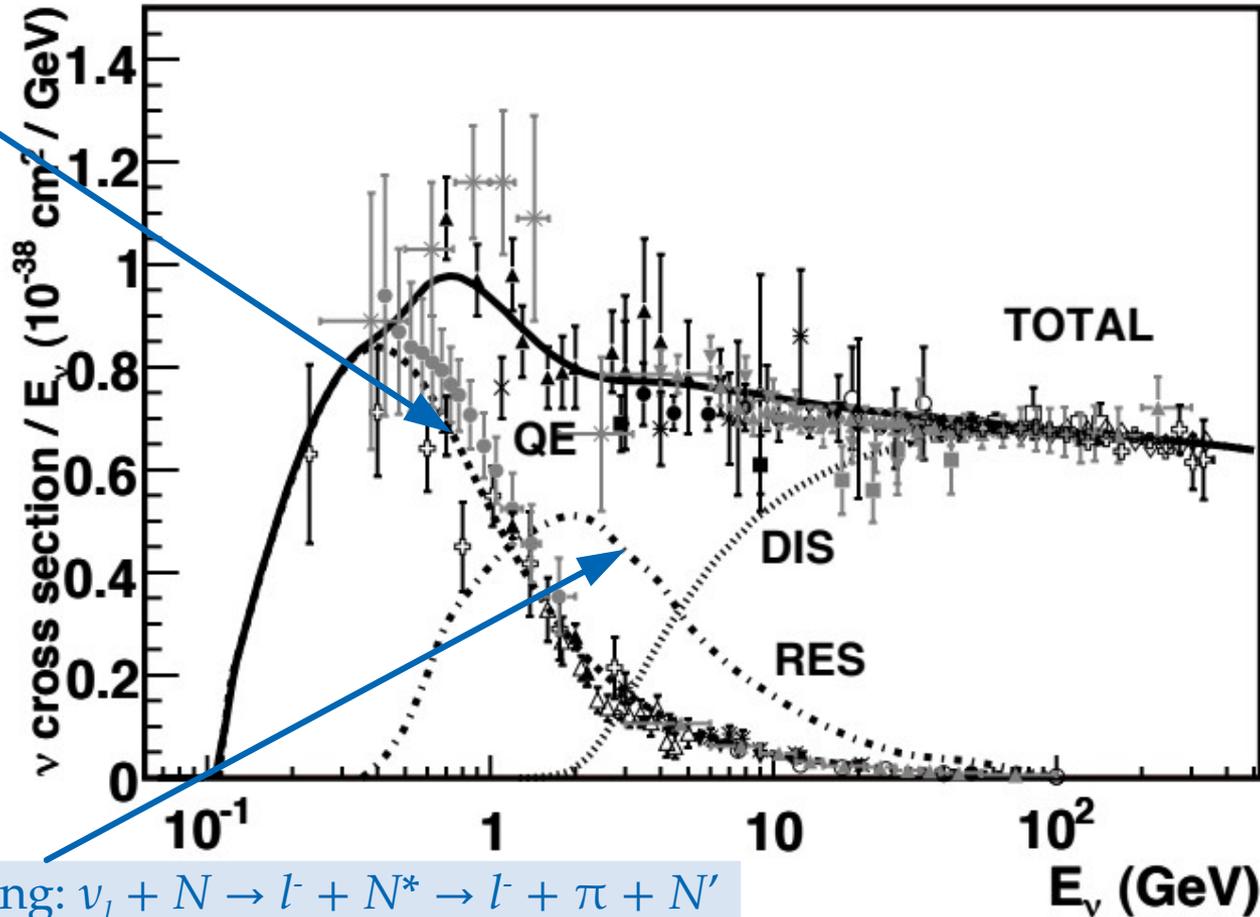
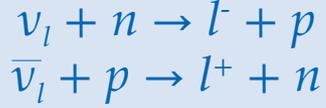
Particle Data Group

Quasi-elastic
scattering:



Particle Data Group

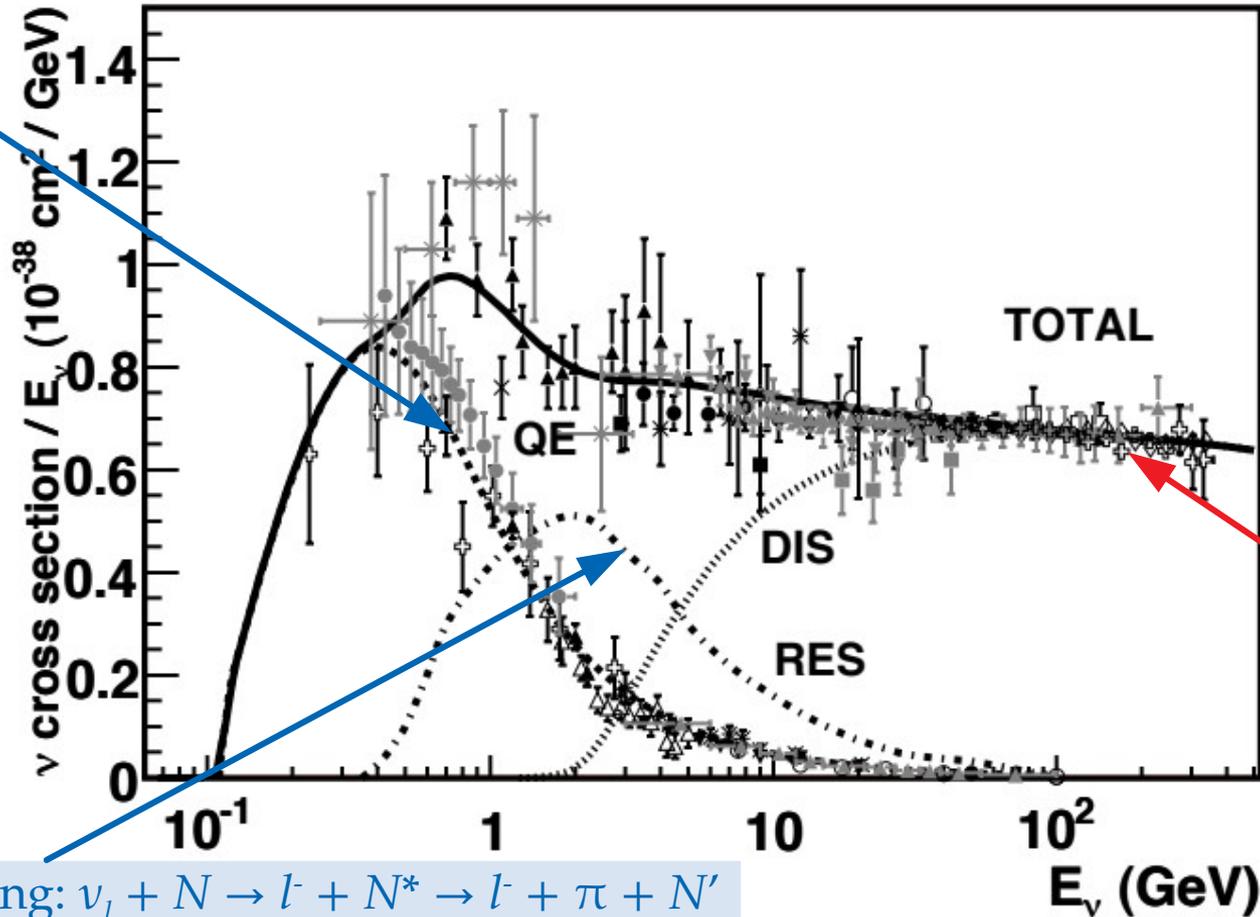
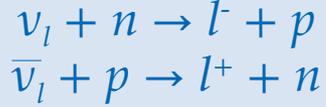
Quasi-elastic scattering:



Resonant scattering: $\nu_l + N \rightarrow l' + N^* \rightarrow l' + \pi + N'$

Particle Data Group

Quasi-elastic scattering:



Deep inelastic scattering:

$$\nu_l + N \rightarrow l^- + X$$
$$\bar{\nu}_l + N \rightarrow l^+ + X$$

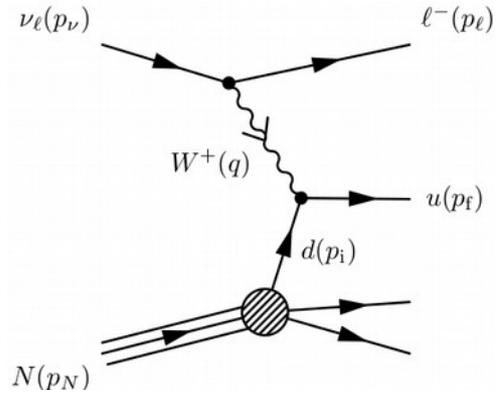
Resonant scattering: $\nu_l + N \rightarrow l^- + N^* \rightarrow l^- + \pi + N'$

Particle Data Group

Extrapolating the cross section to high energies

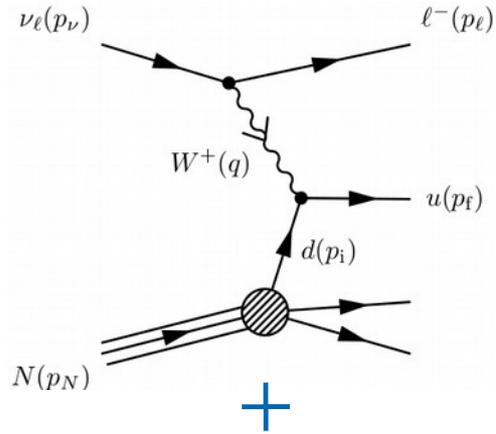
Extrapolating the cross section to high energies

SM

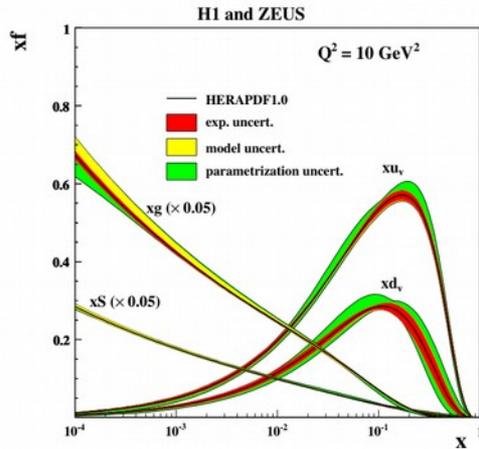


Extrapolating the cross section to high energies

SM

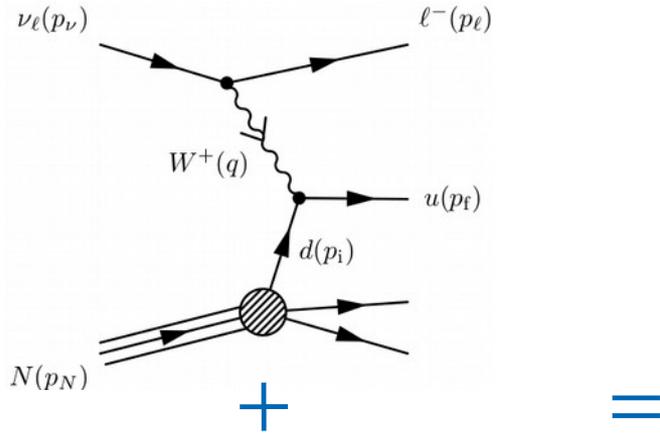


PDFs

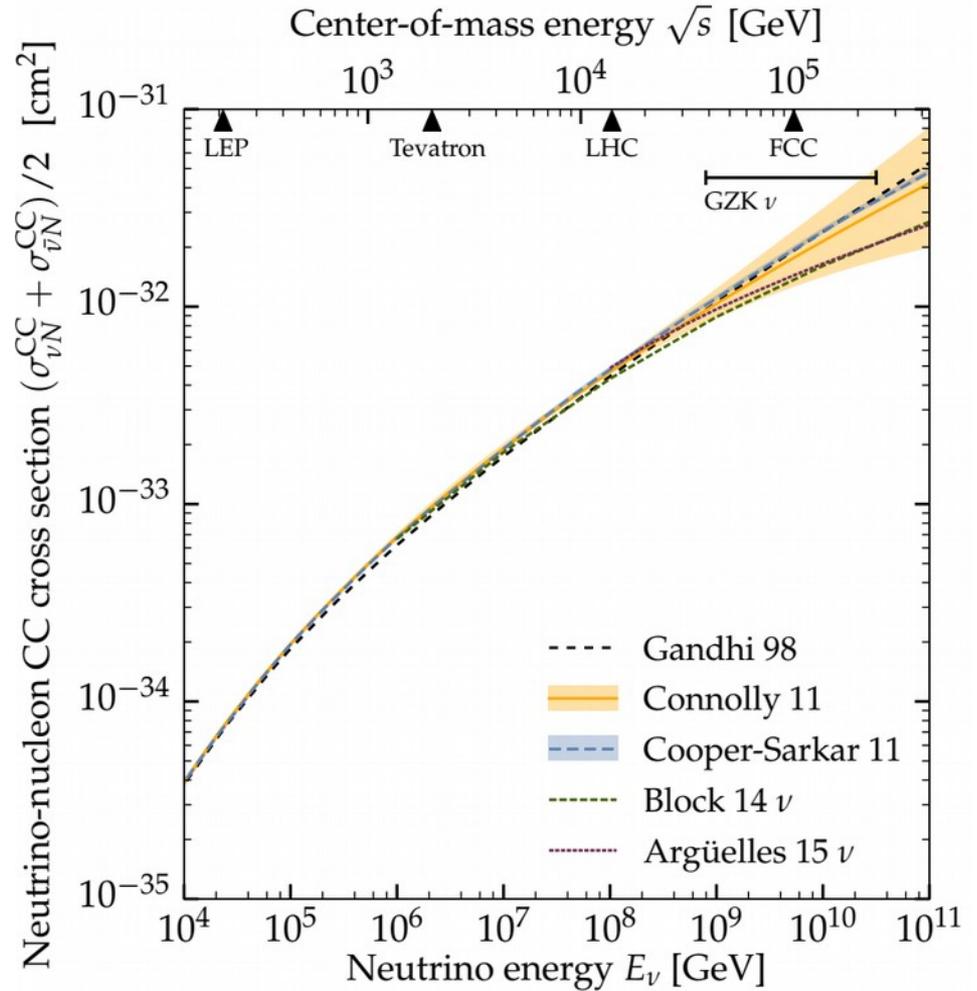
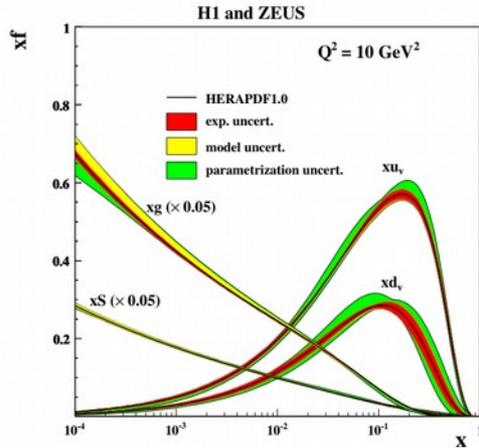


Extrapolating the cross section to high energies

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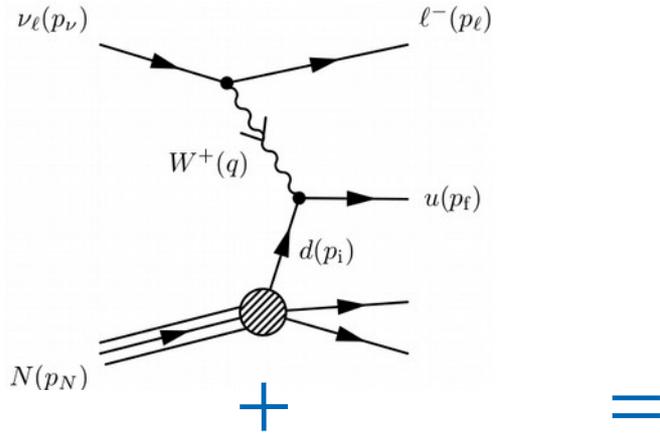


PDFs

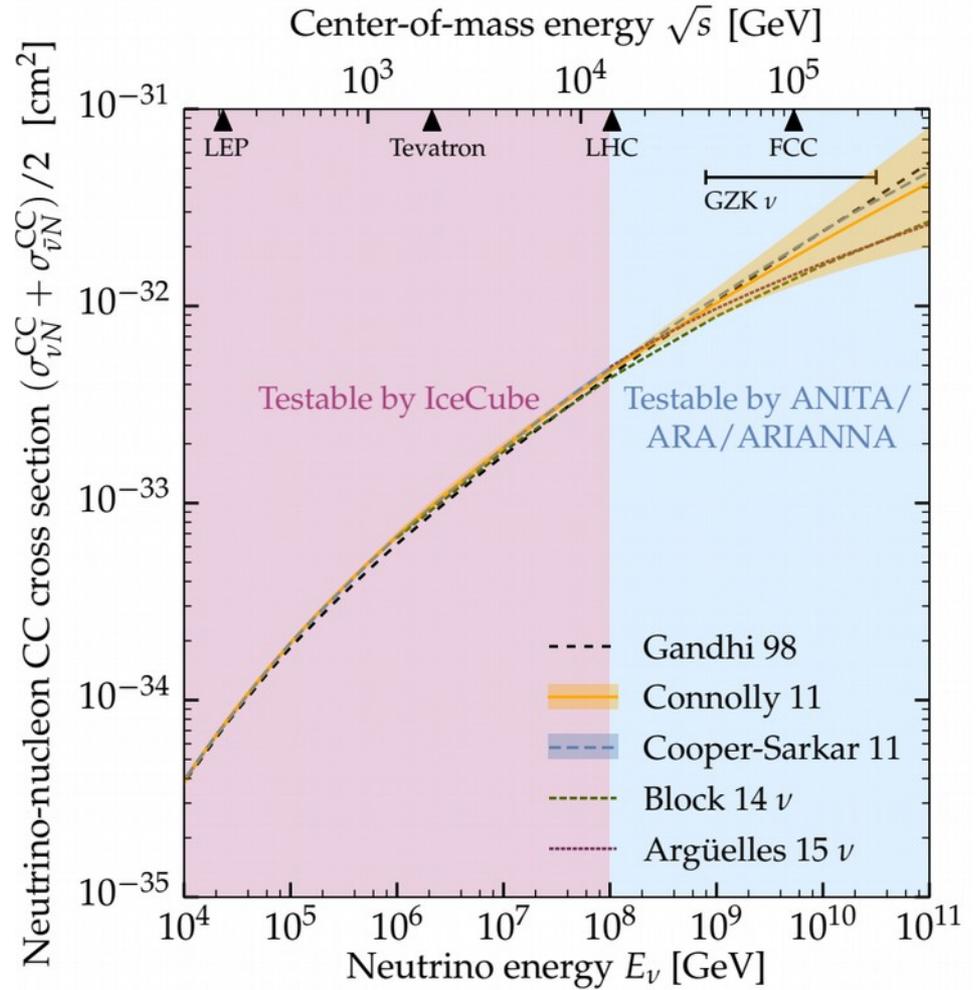
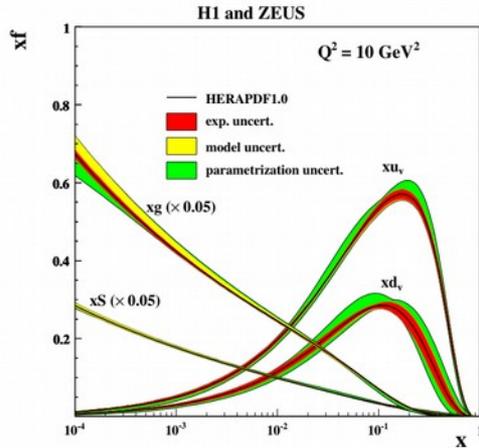


Extrapolating the cross section to high energies

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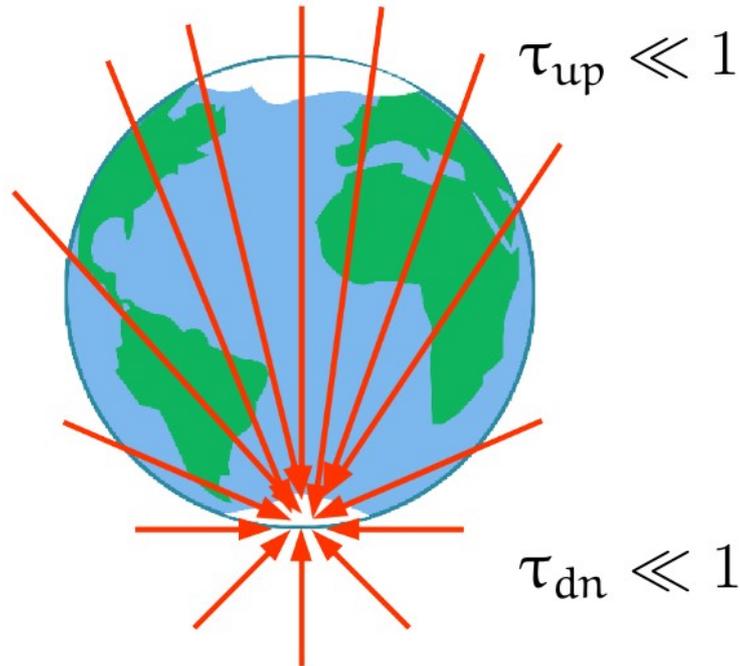
PDFs



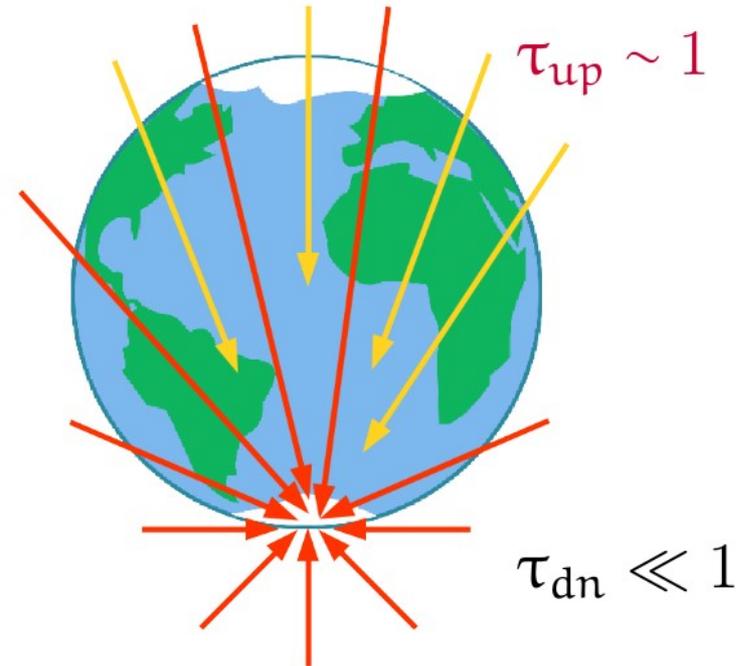
Measuring the high-energy cross section

$$\text{Optical depth to } \nu N \text{ int's} = \frac{\text{Distance from Earth's surface to IceCube}}{\text{Mean free path inside Earth}} \equiv \tau(E_\nu, \theta_z) \propto \sigma_{\nu N}$$

Below ~ 10 TeV: Earth is transparent



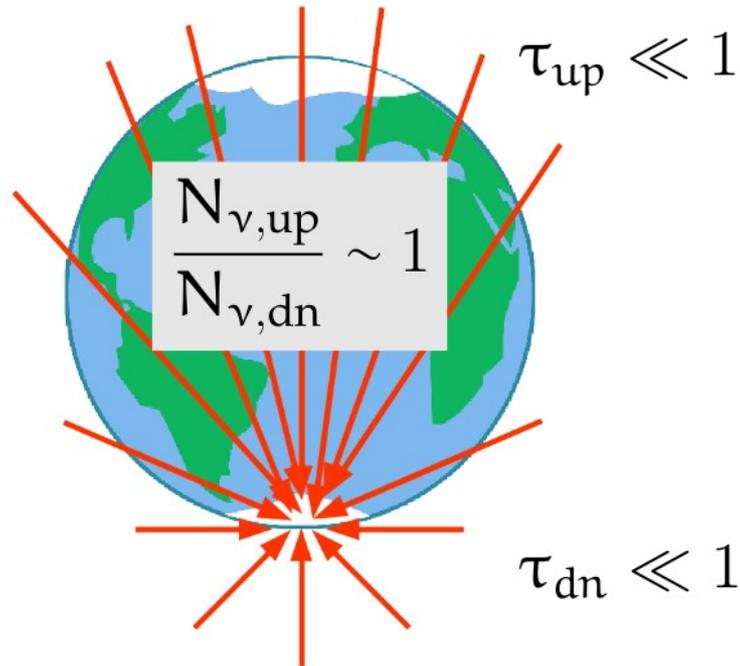
Above ~ 10 TeV: Earth is opaque



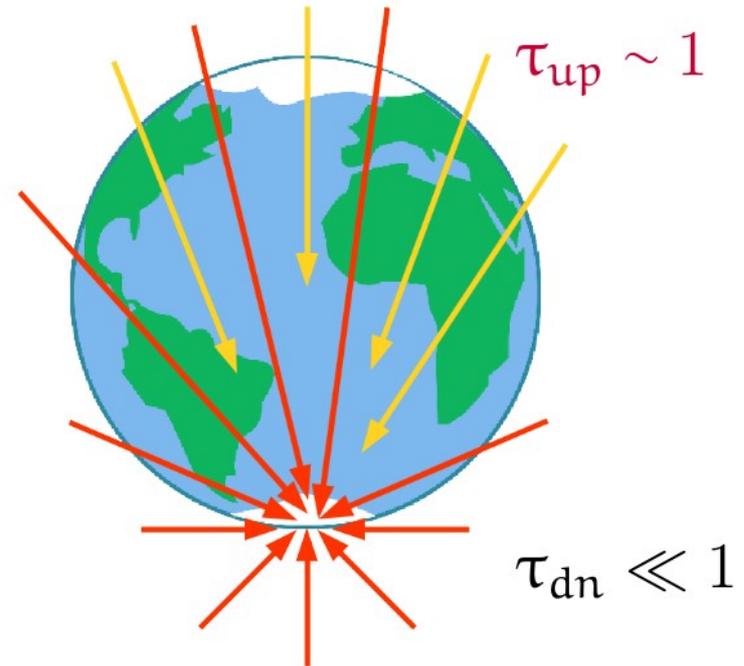
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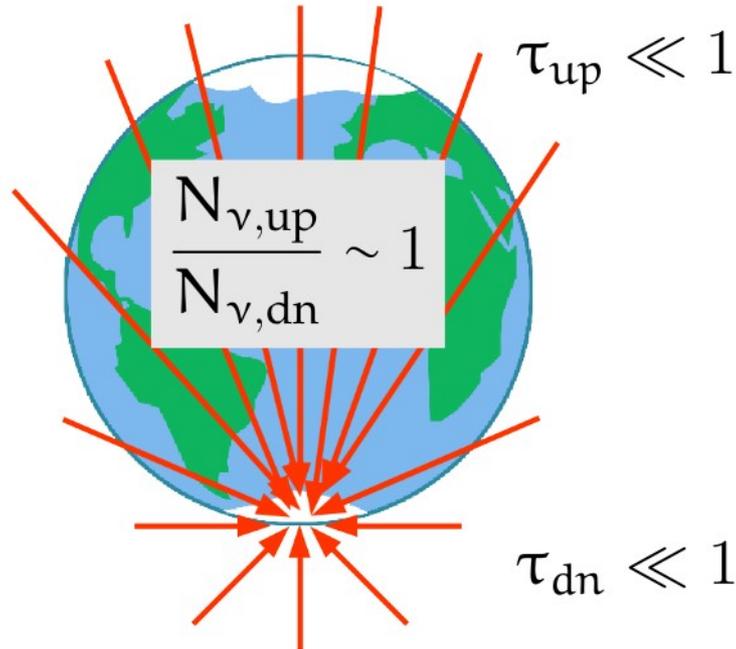
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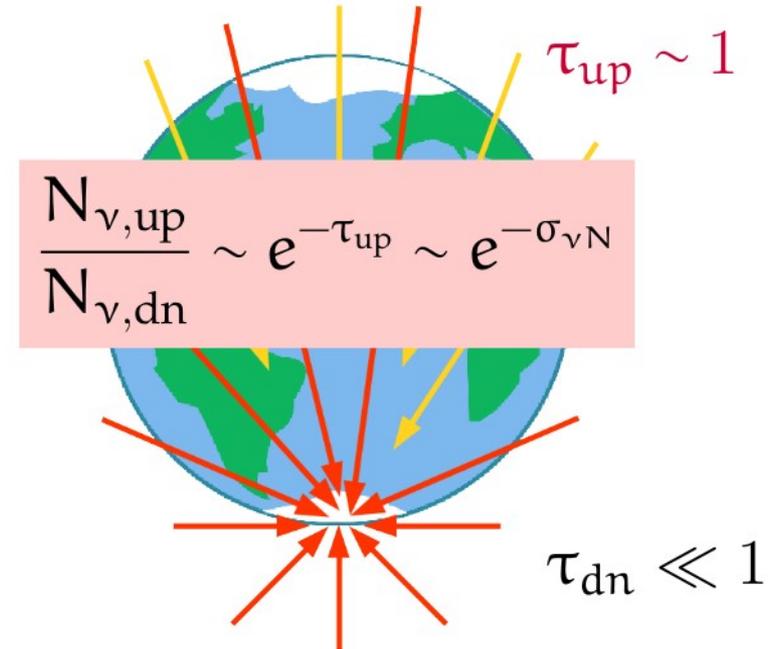
Measuring the high-energy cross section

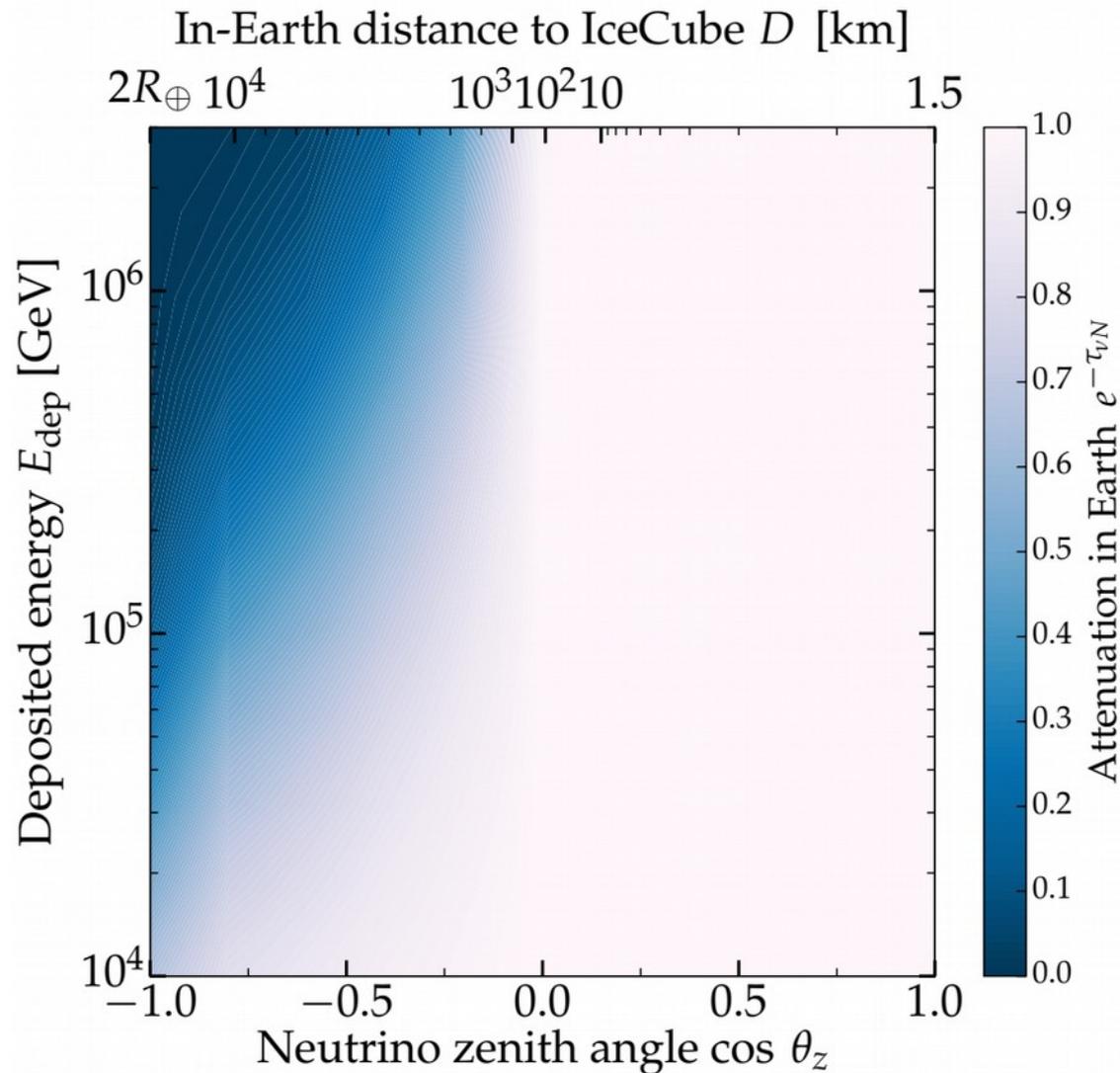
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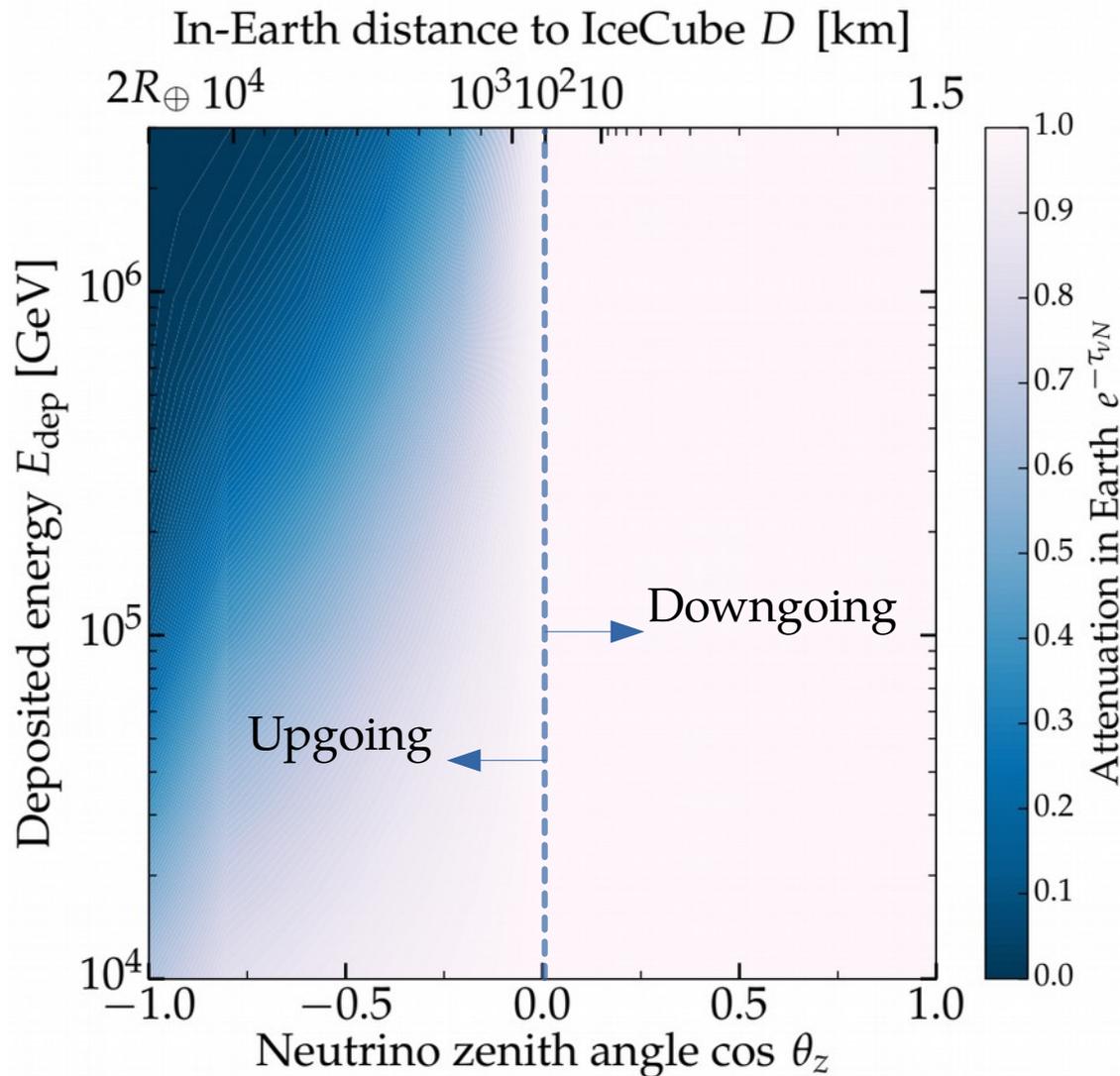
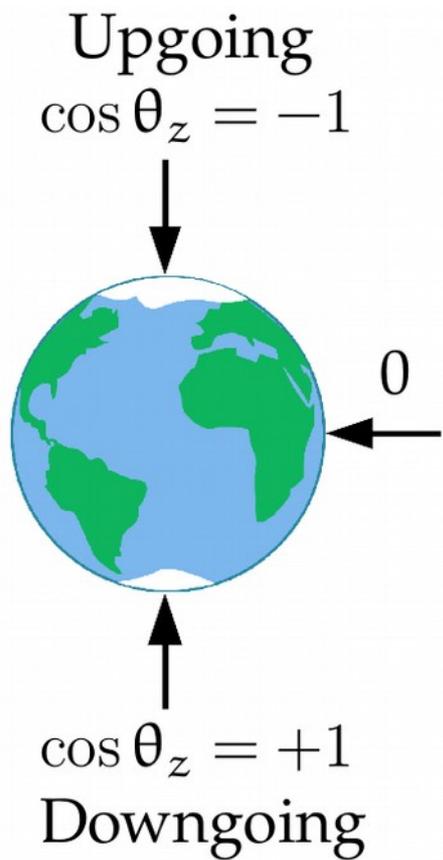
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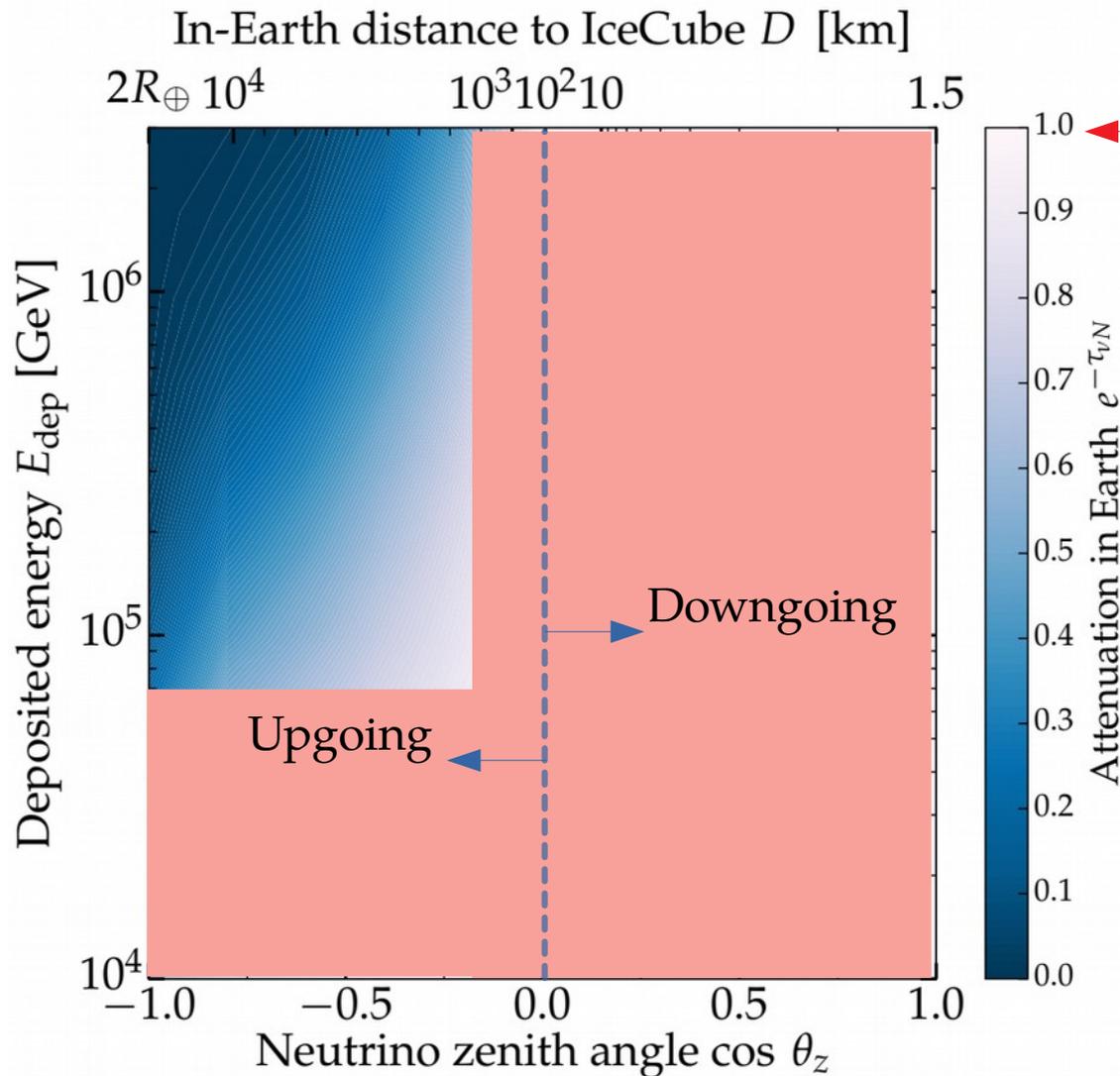
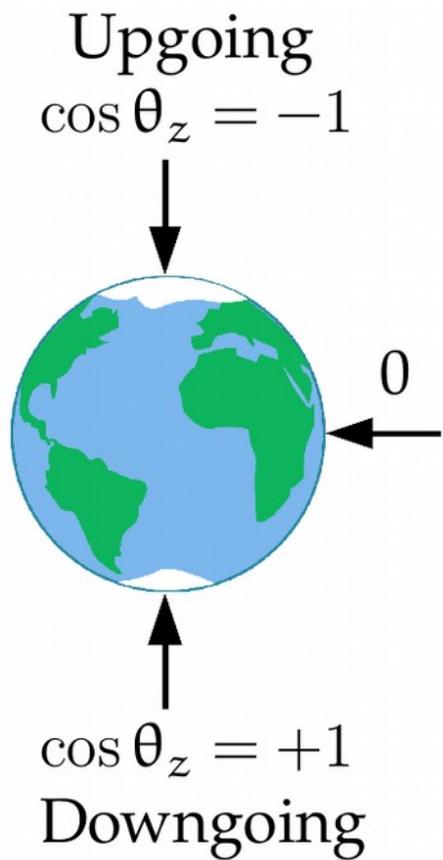


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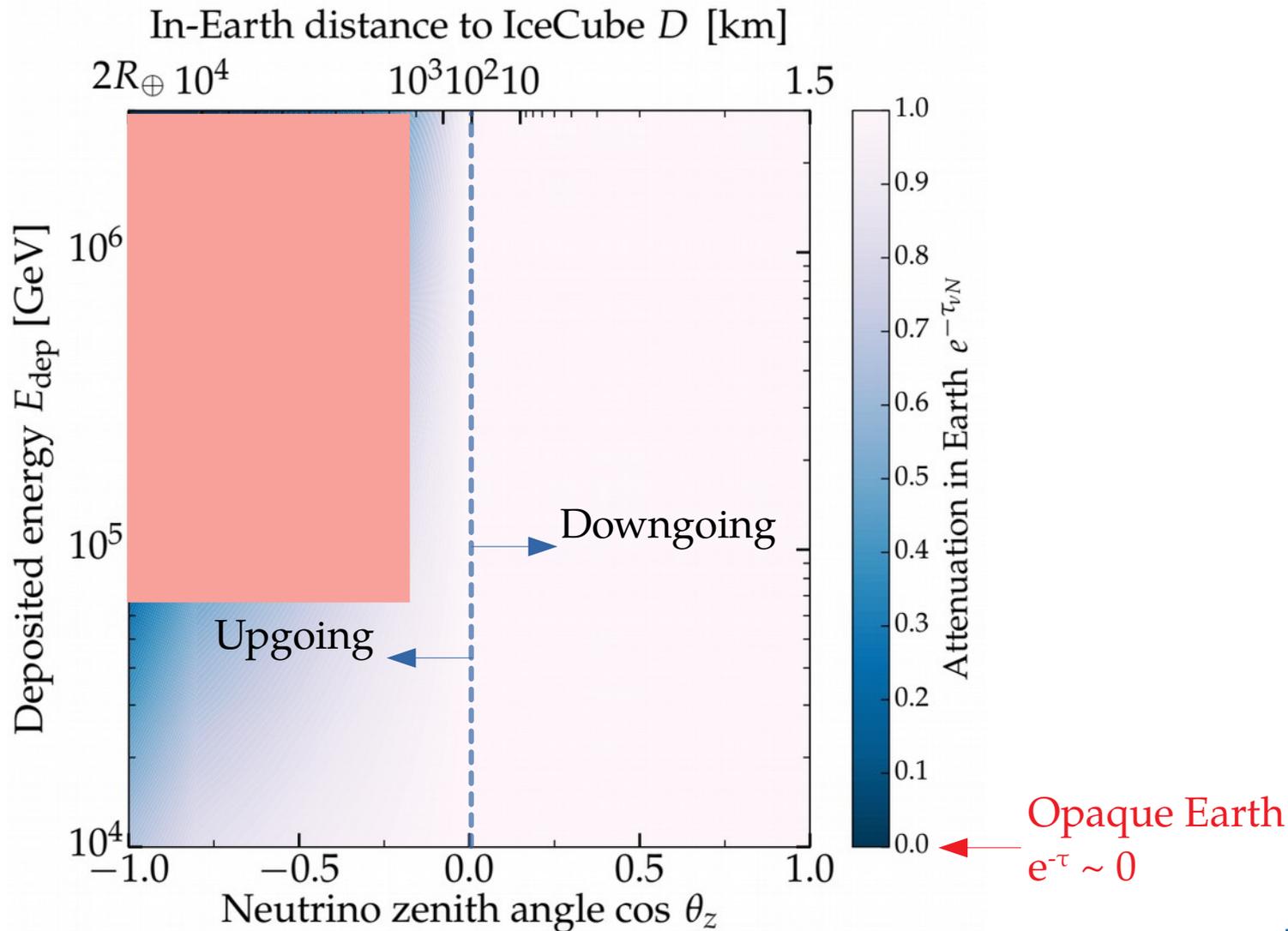
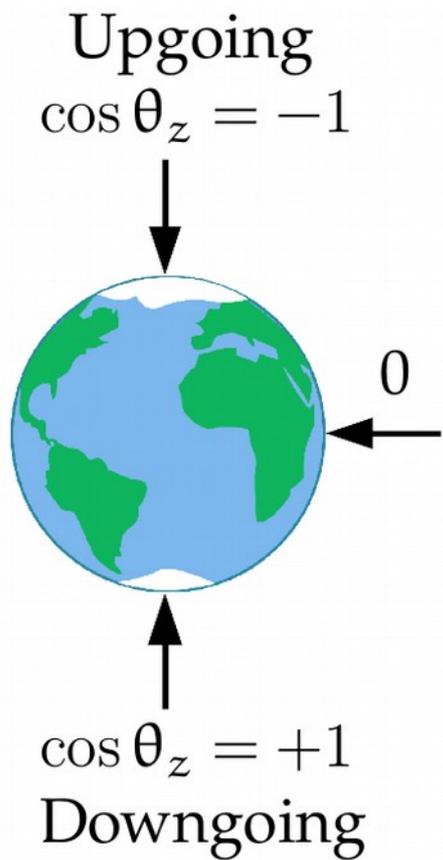


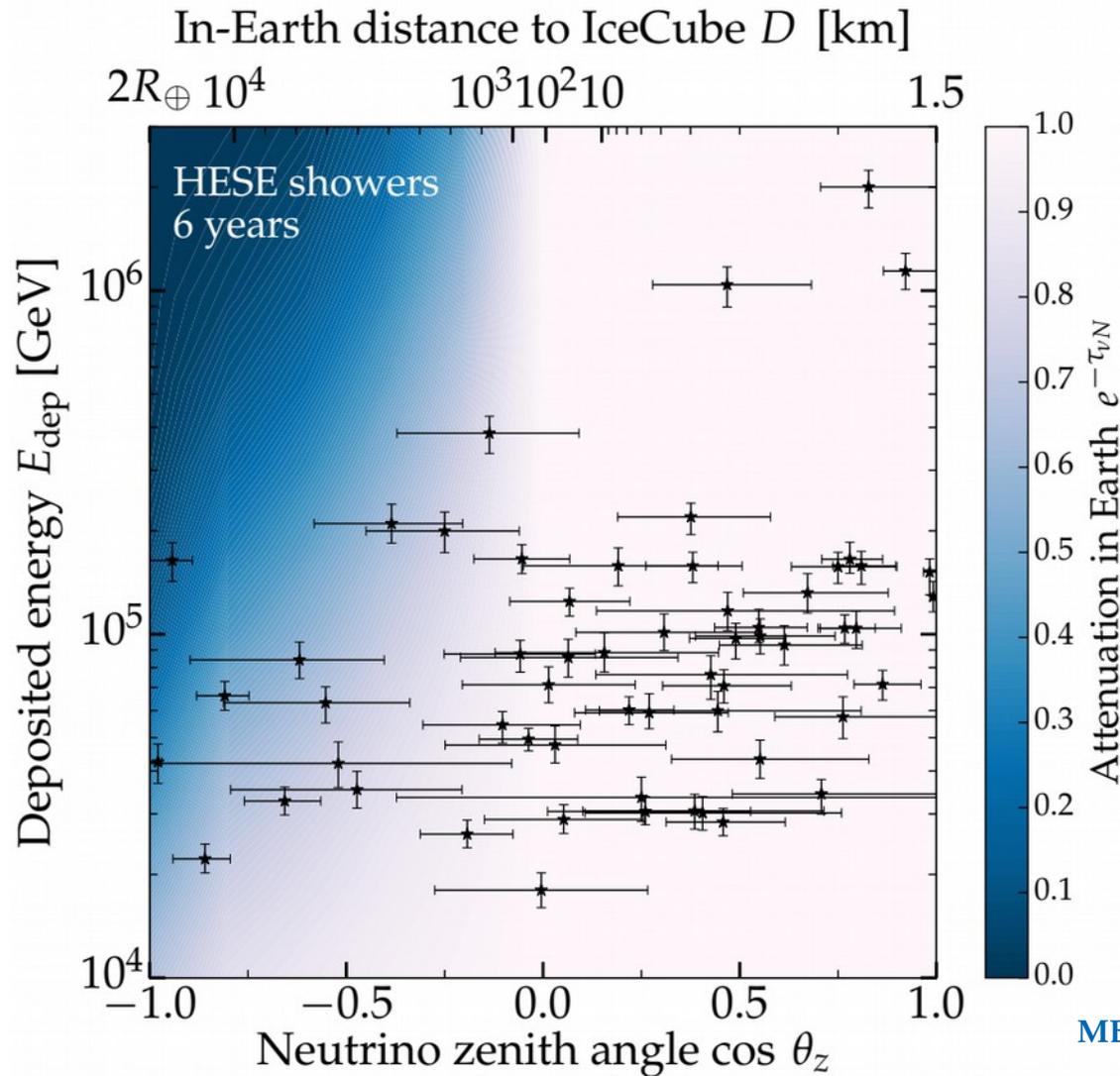




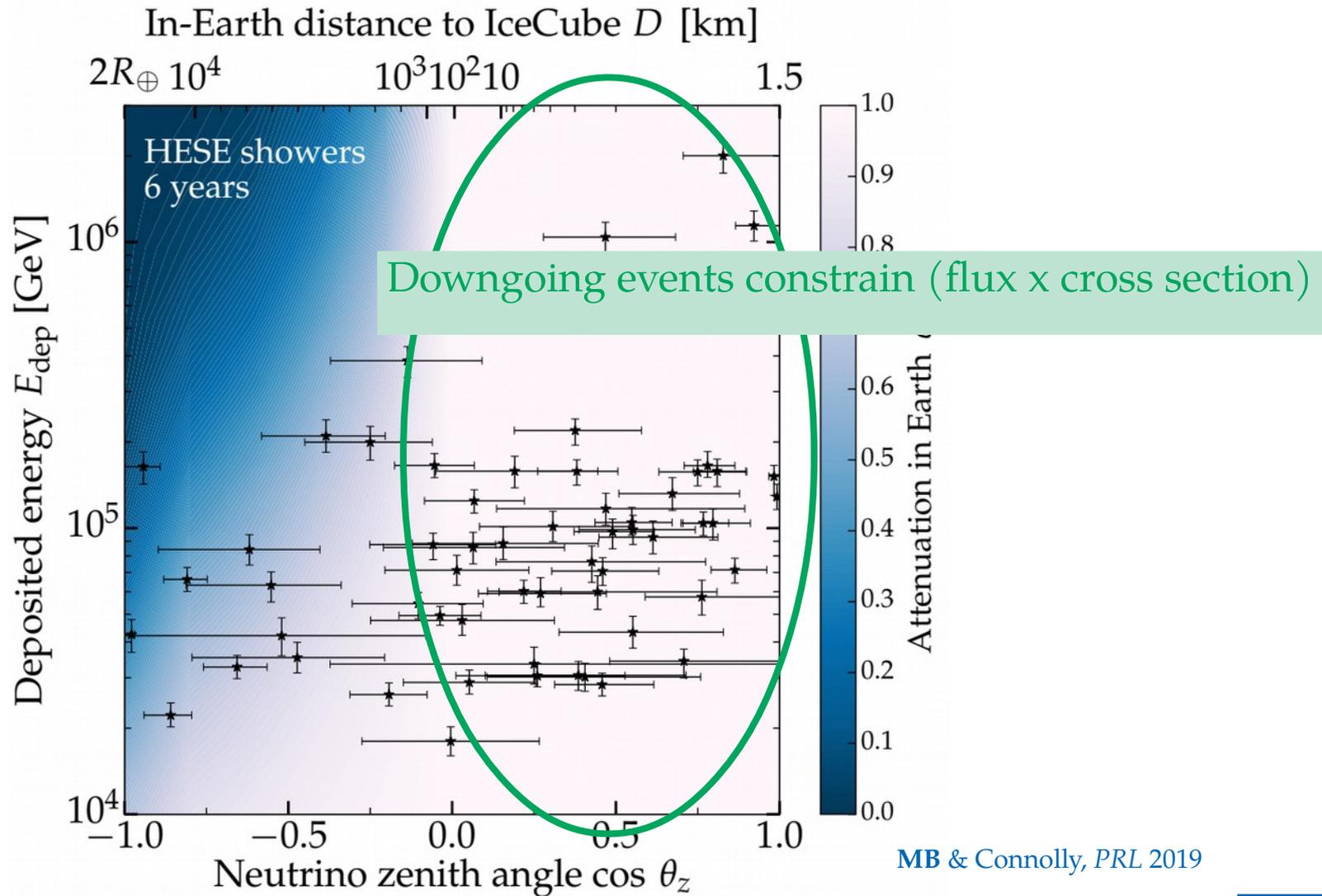


Transparent Earth
 $e^{-\tau} \sim 1$





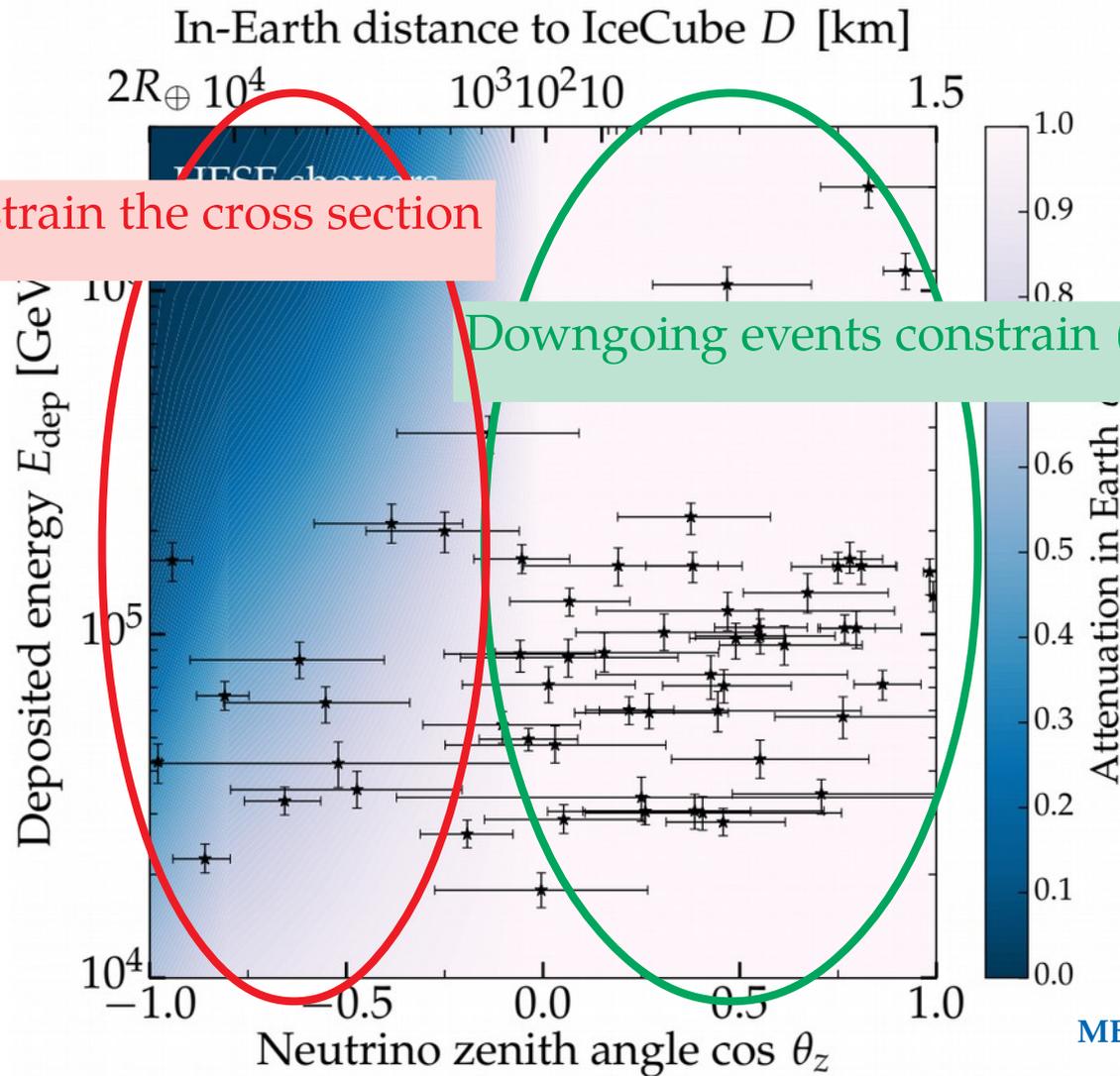
MB & Connolly, *PRL* 2019



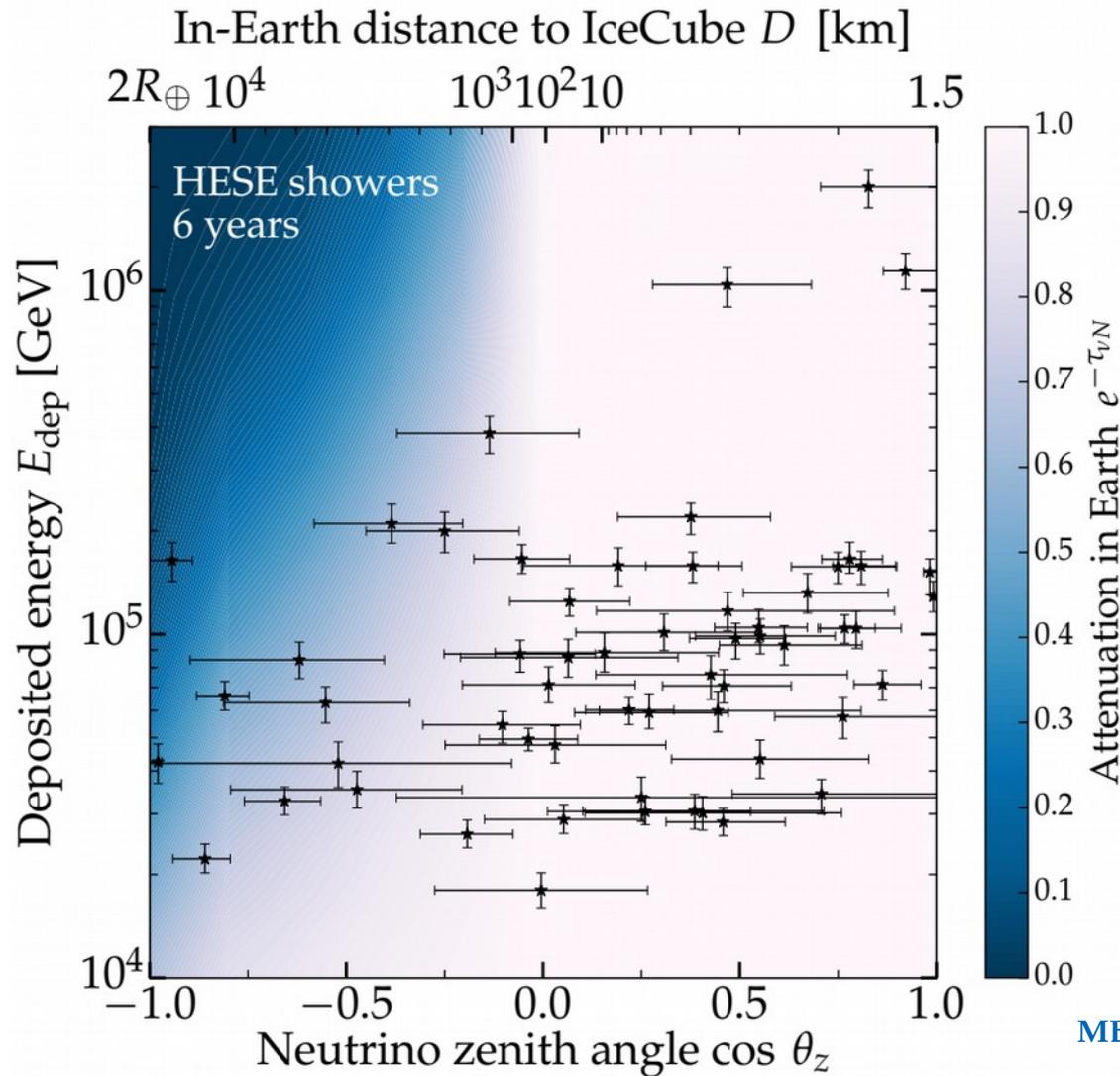
MB & Connolly, *PRL* 2019

Upgoing events constrain the cross section

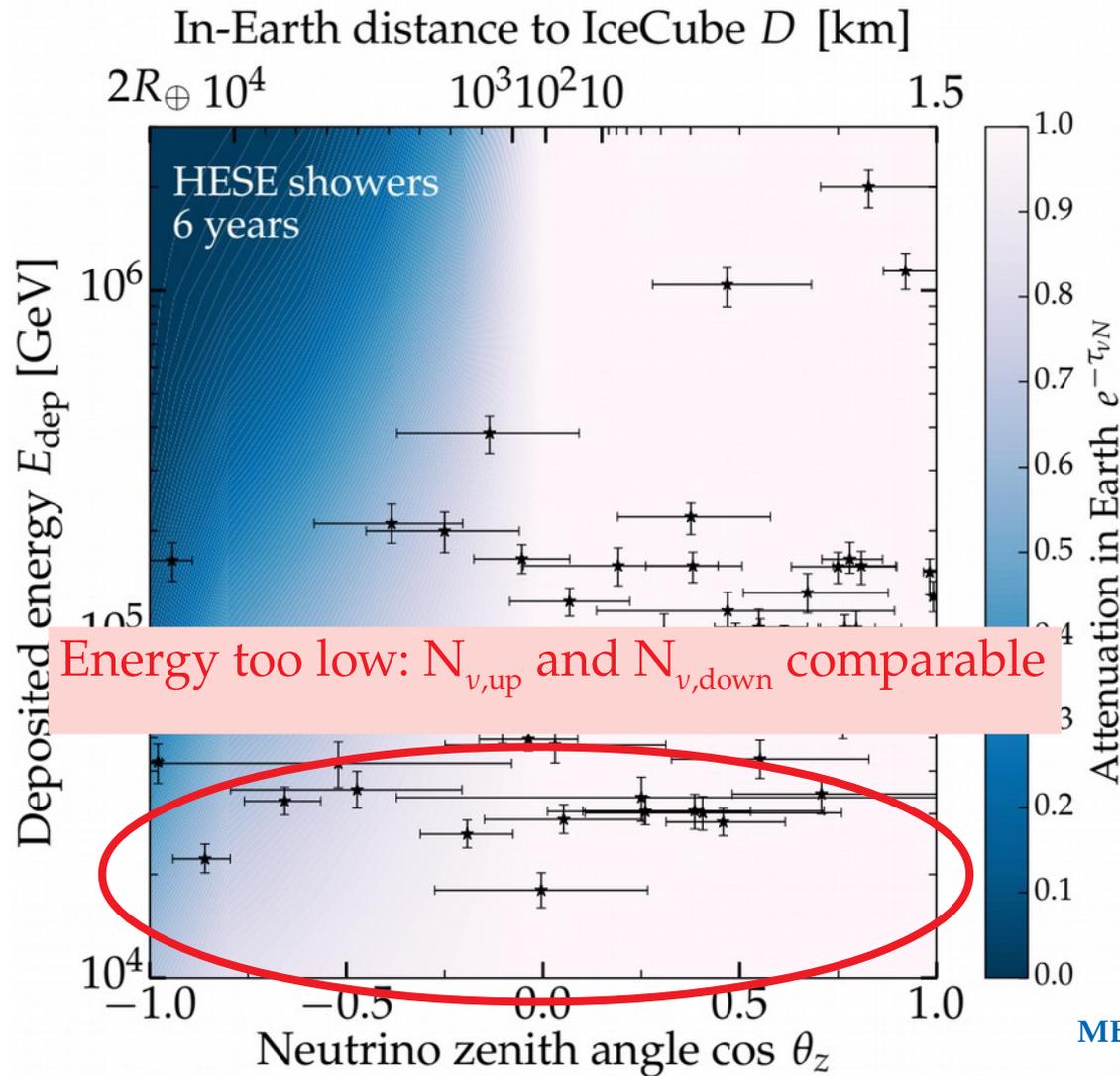
Downgoing events constrain (flux x cross section)



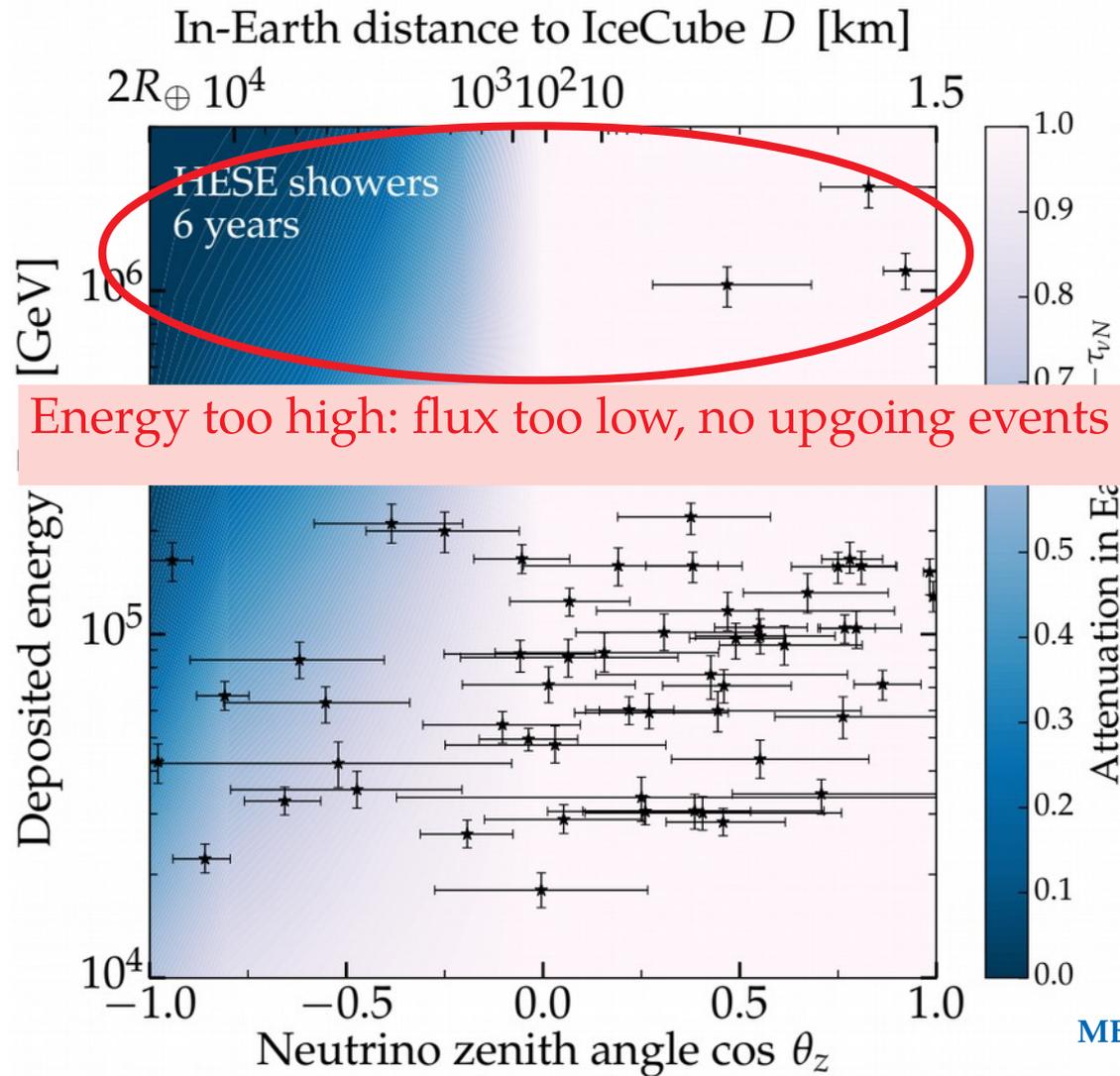
MB & Connolly, *PRL* 2019



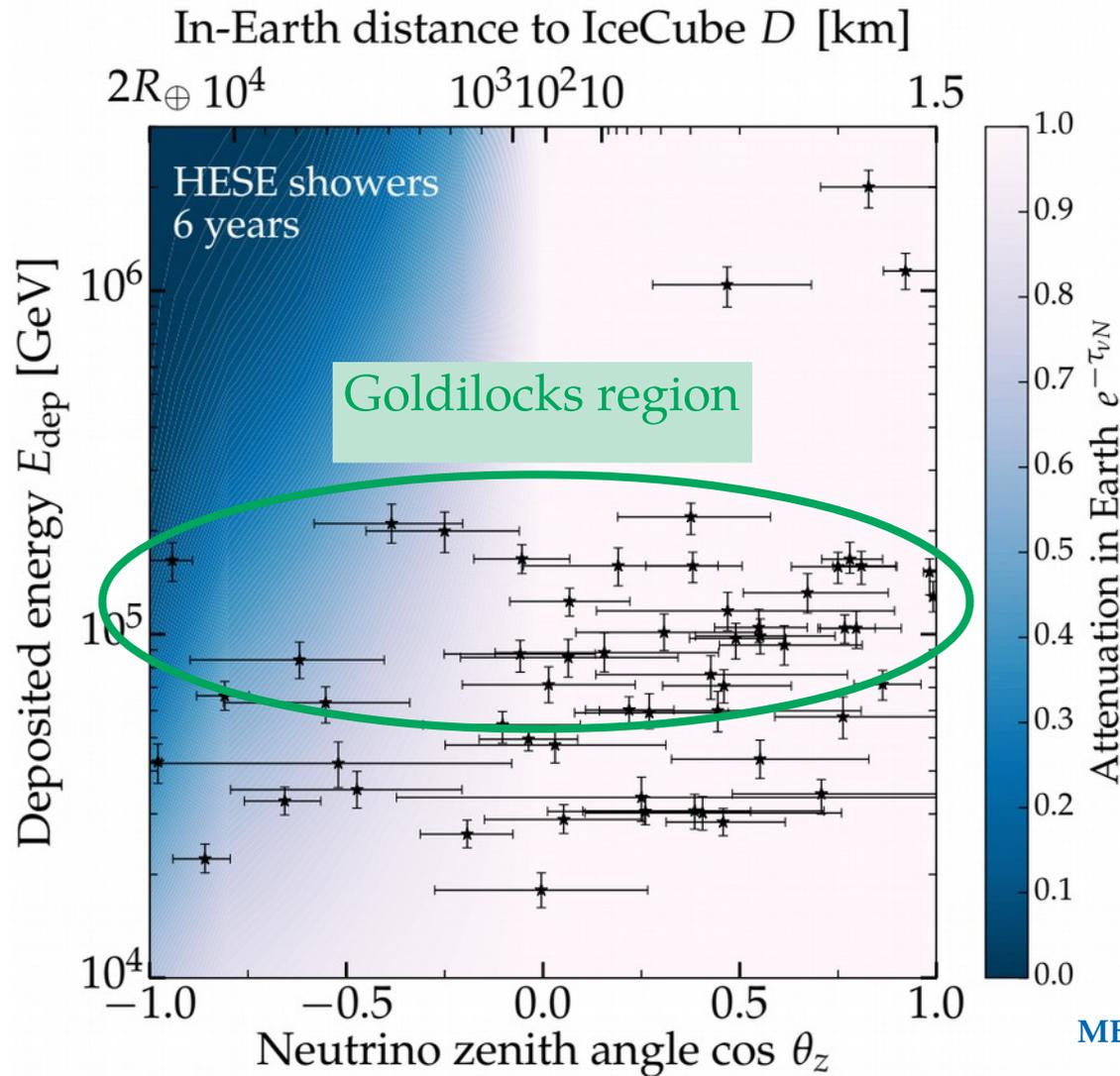
MB & Connolly, *PRL* 2019



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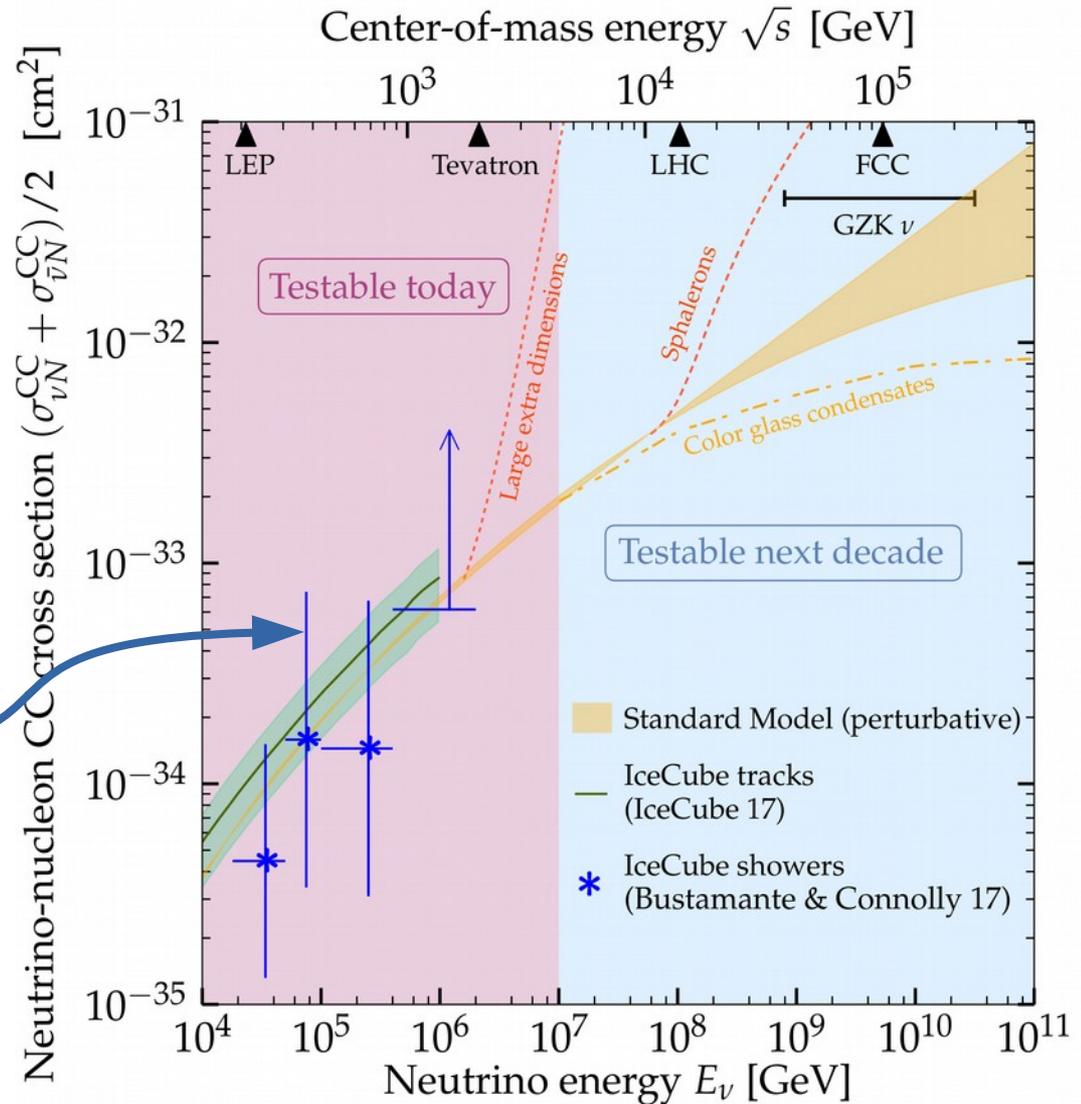


MB & Connolly, *PRL* 2019



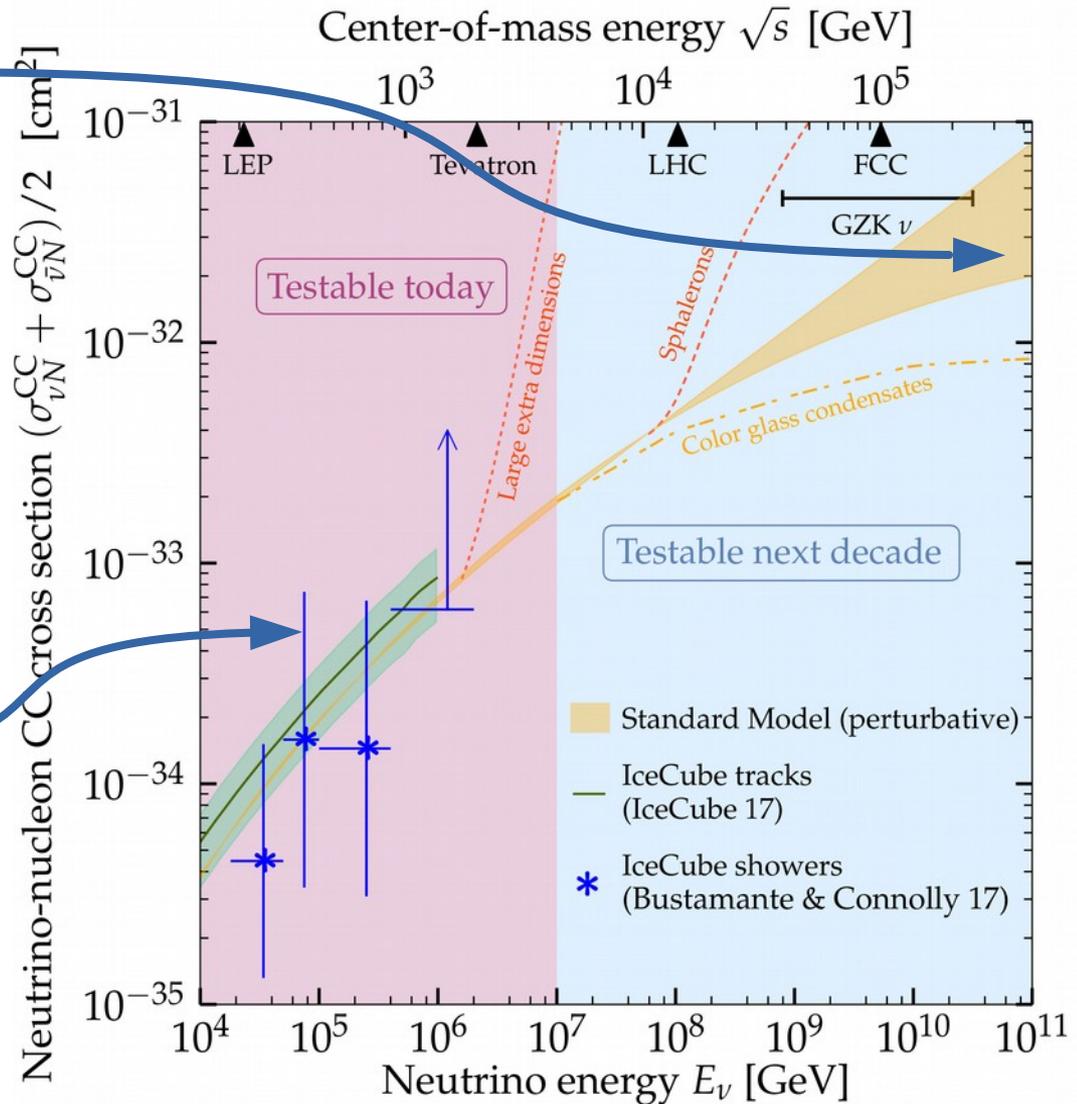
MB & Connolly, *PRL* 2019

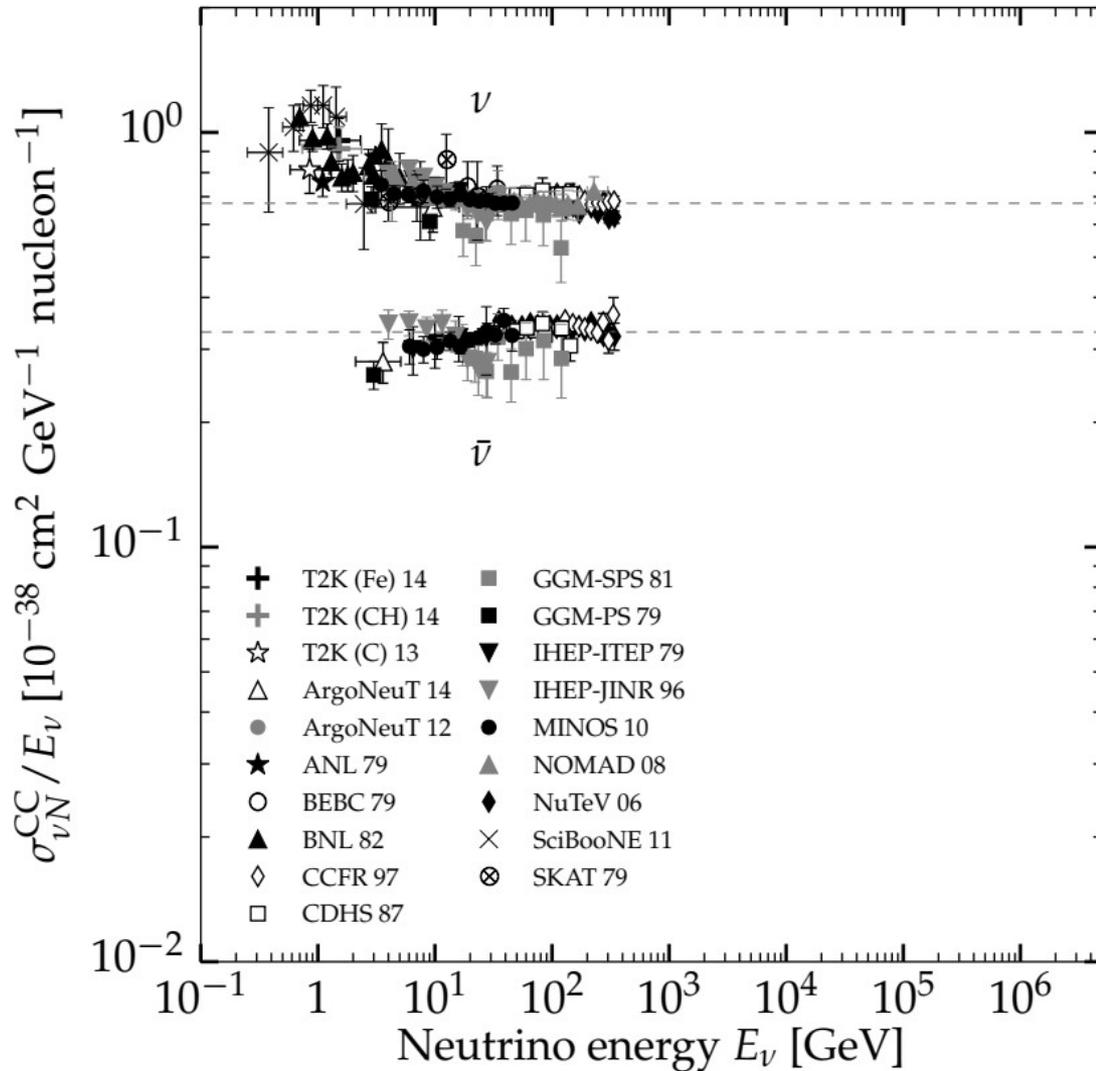
- ▶ Fold in astrophysical unknowns (spectral index, normalization)
- ▶ Compatible with SM predictions
- ▶ Still room for new physics
- ▶ Today, using IceCube:
 - ▶ Extracted from ~ 60 showers in 6 yr
 - ▶ Limited by statistics
- ▶ Future, using IceCube-Gen2:
 - ▶ $\times 5$ volume \Rightarrow 300 showers in 6 yr
 - ▶ Reduce statistical error by 40%

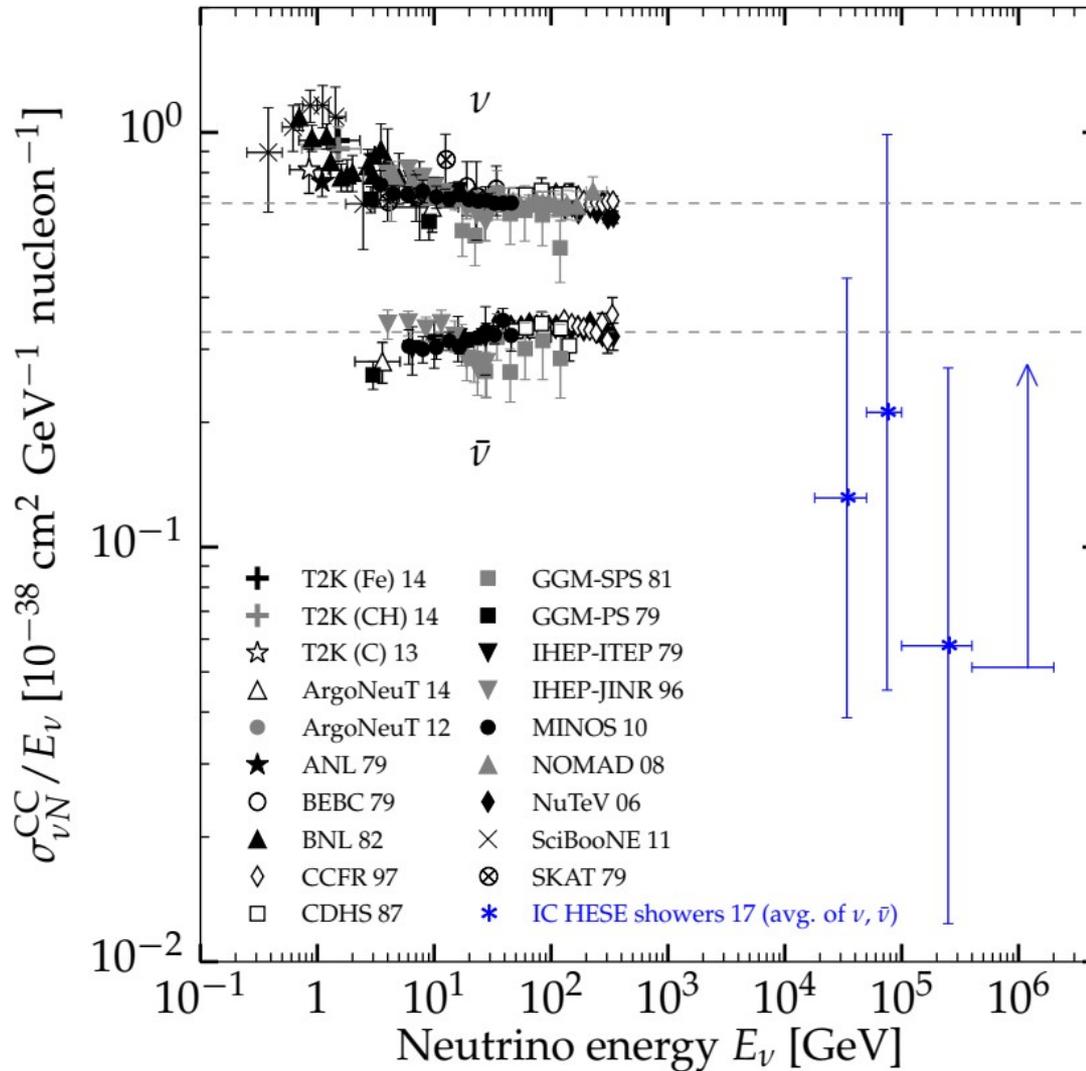


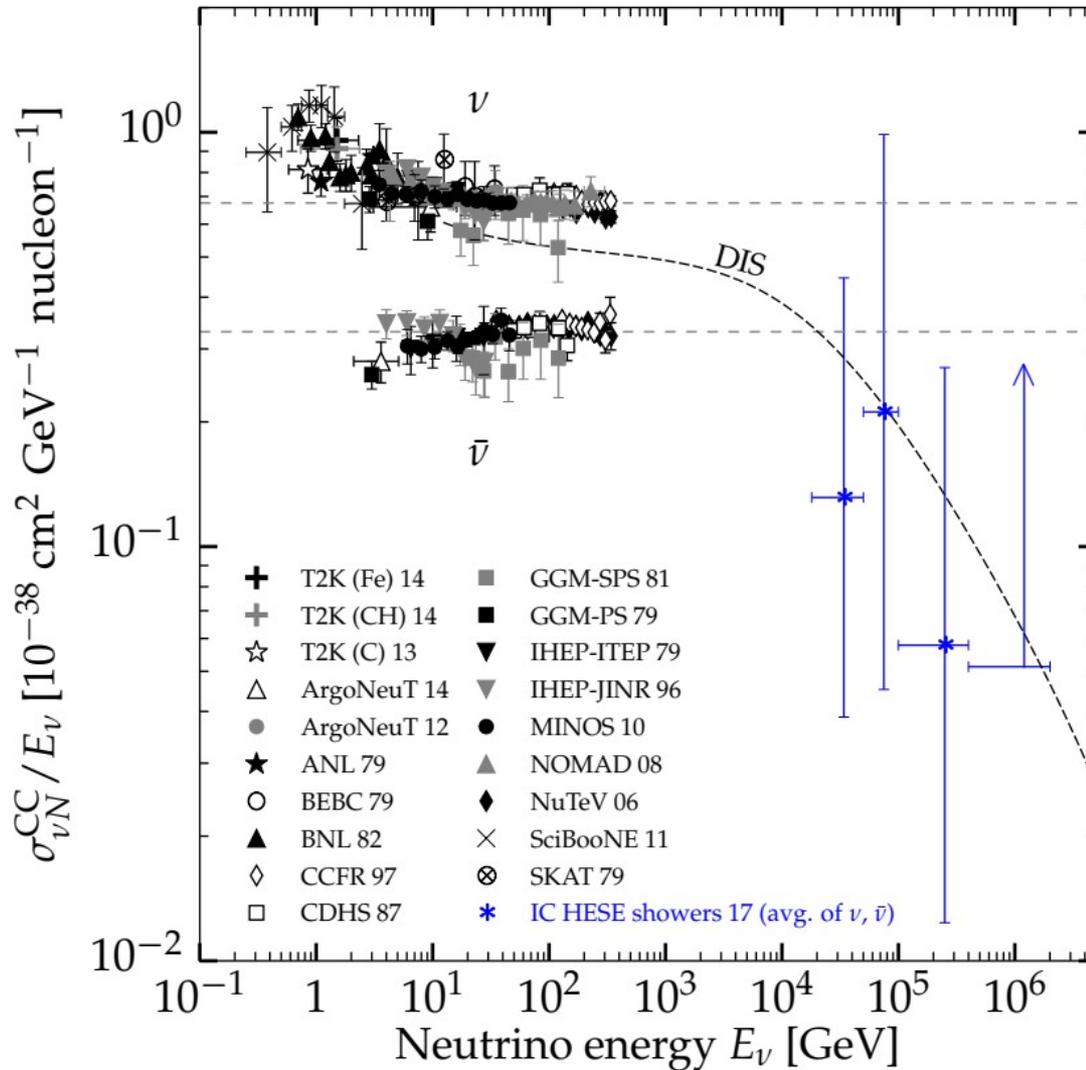
UHE uncertainties are actually smaller:
 Cooper-Sarkar, Mertsch, Sarkar *et al.*, *JHEP* 2011

- ▶ Fold in astrophysical unknowns (spectral index, normalization)
- ▶ Compatible with SM predictions
- ▶ Still room for new physics
- ▶ Today, using IceCube:
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 - ▶ × 5 volume ⇒ 300 showers in 6 yr
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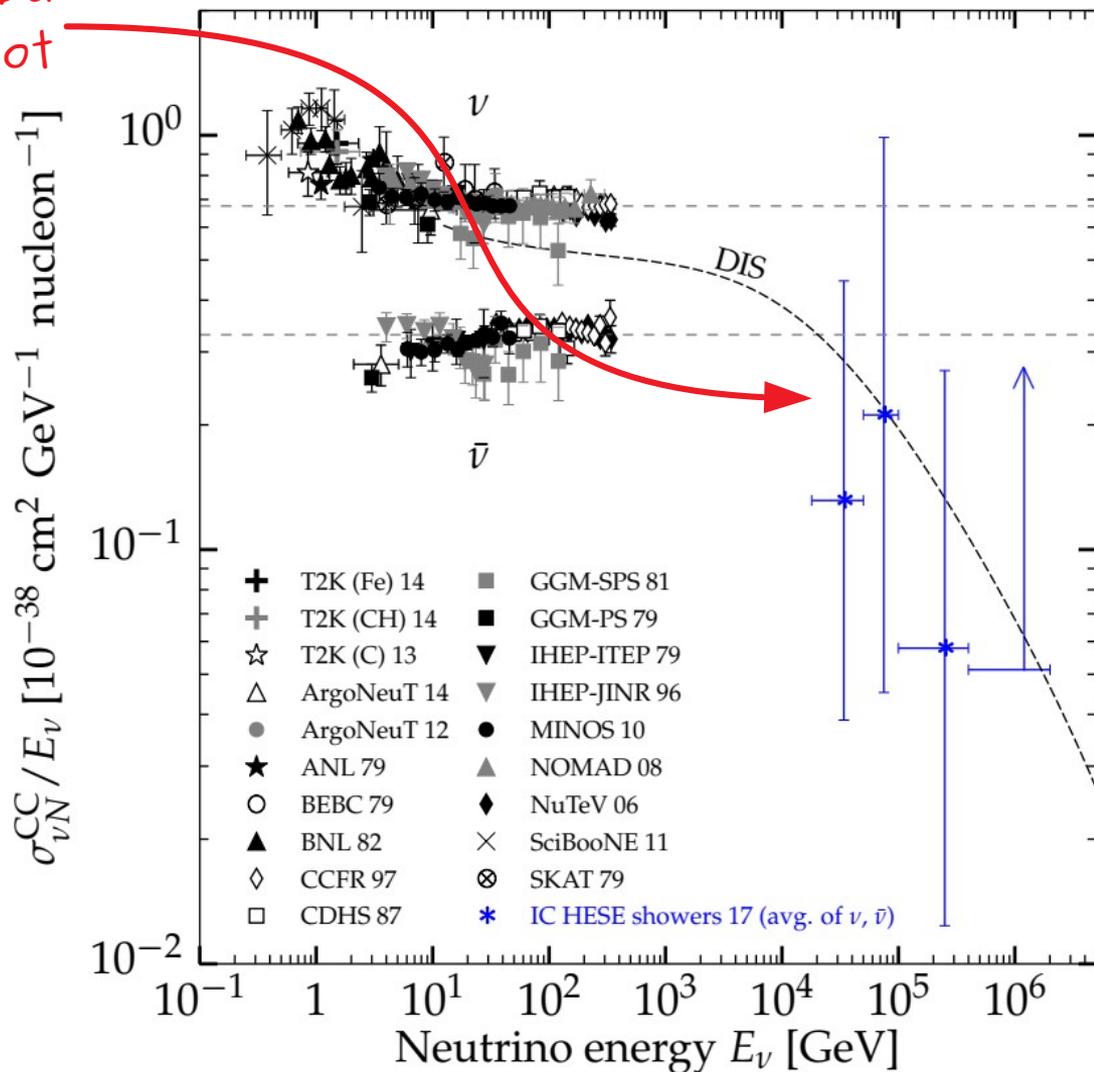








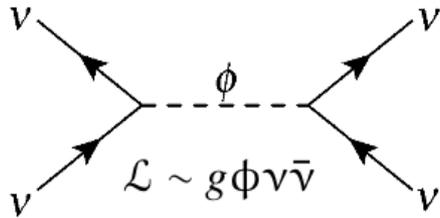
Extending the PDG cross-section plot



MB & Connolly PRL 2019
See also: IceCube, Nature 2017

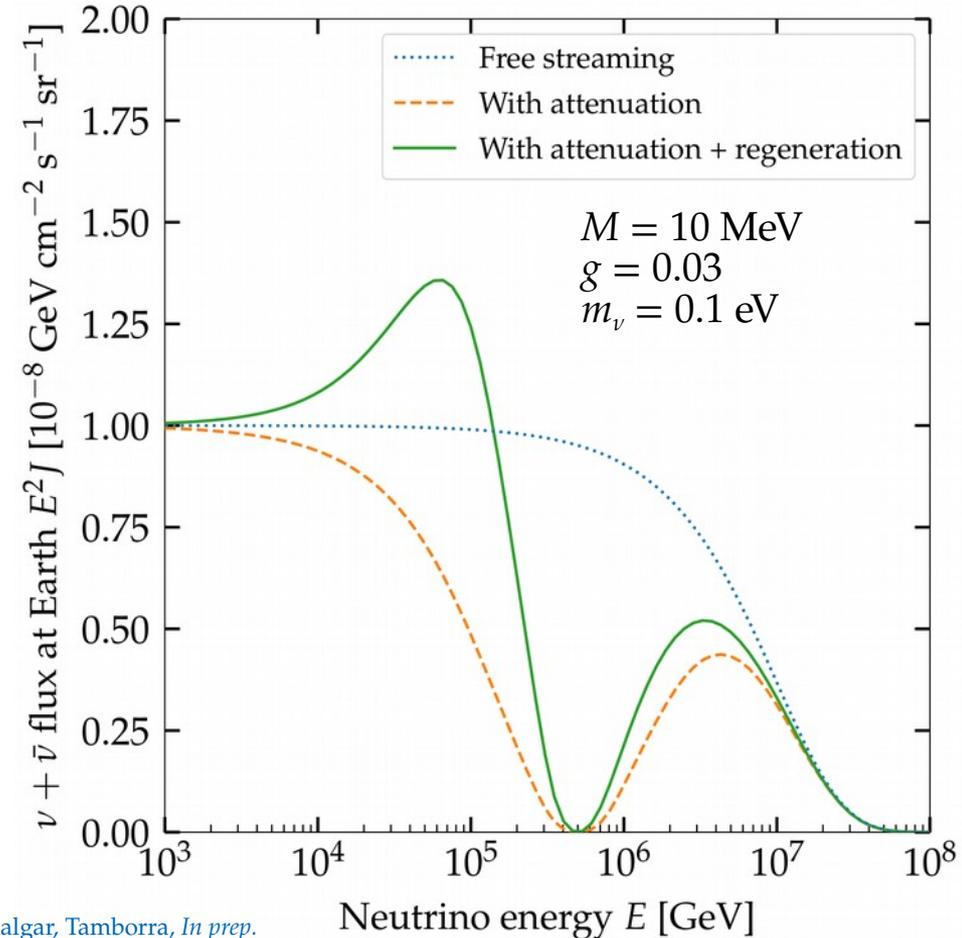
New physics in the spectral shape: $\nu\nu$ interactions

“Secret” neutrino interactions between astrophysical ν (PeV) and relic ν (0.1 meV):



Cross section:
$$\sigma = \frac{g^4}{4\pi} \frac{s}{(s - M^2)^2 + M^2\Gamma^2}$$

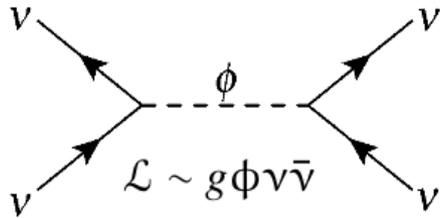
Resonance energy:
$$E_{\text{res}} = \frac{M^2}{2m_\nu}$$



MB, Rosenstroem, Shalgar, Tamborra, *In prep.*
Ng & Beacom, *PRD* 2014
Cherry, Friedland, Shoemaker, 1411.1071
Blum, Hook, Murase, 1408.3799

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New coupling

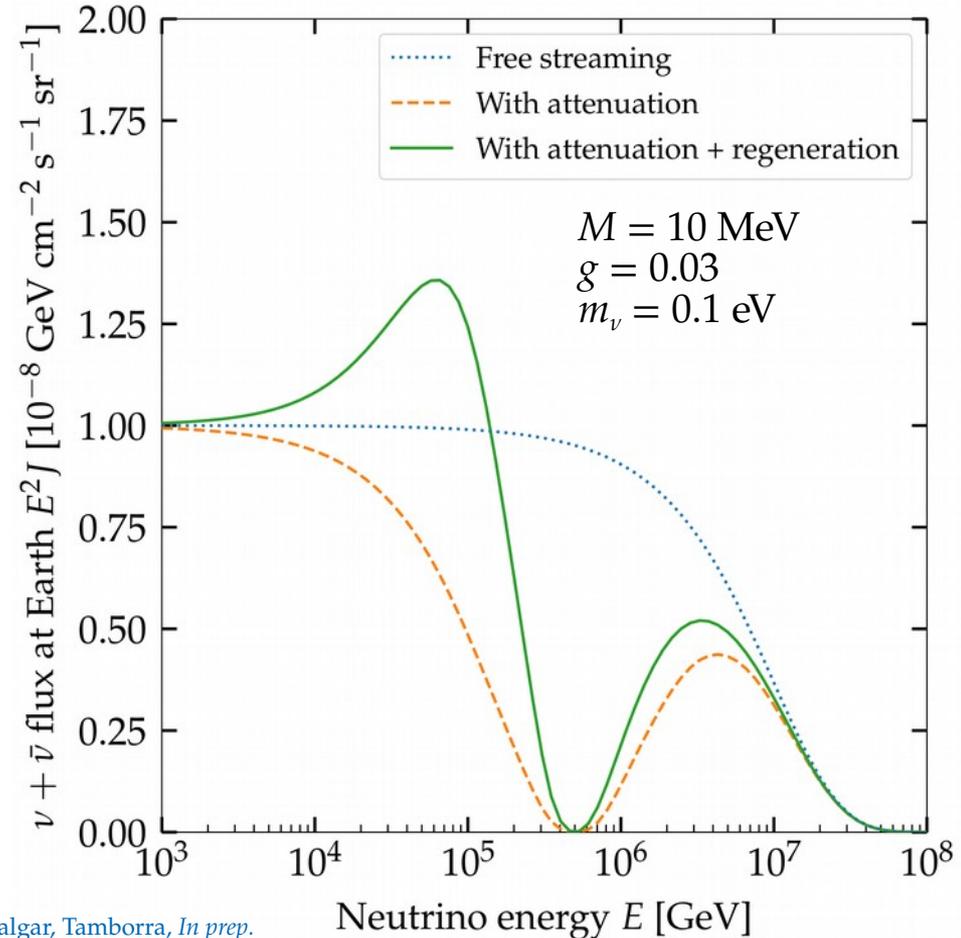
Cross section:

$$\sigma = \frac{g^4 s}{4\pi (s - M^2)^2 + M^2\Gamma^2}$$

Mediator mass

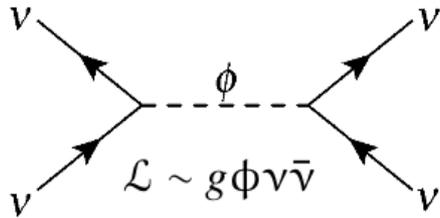
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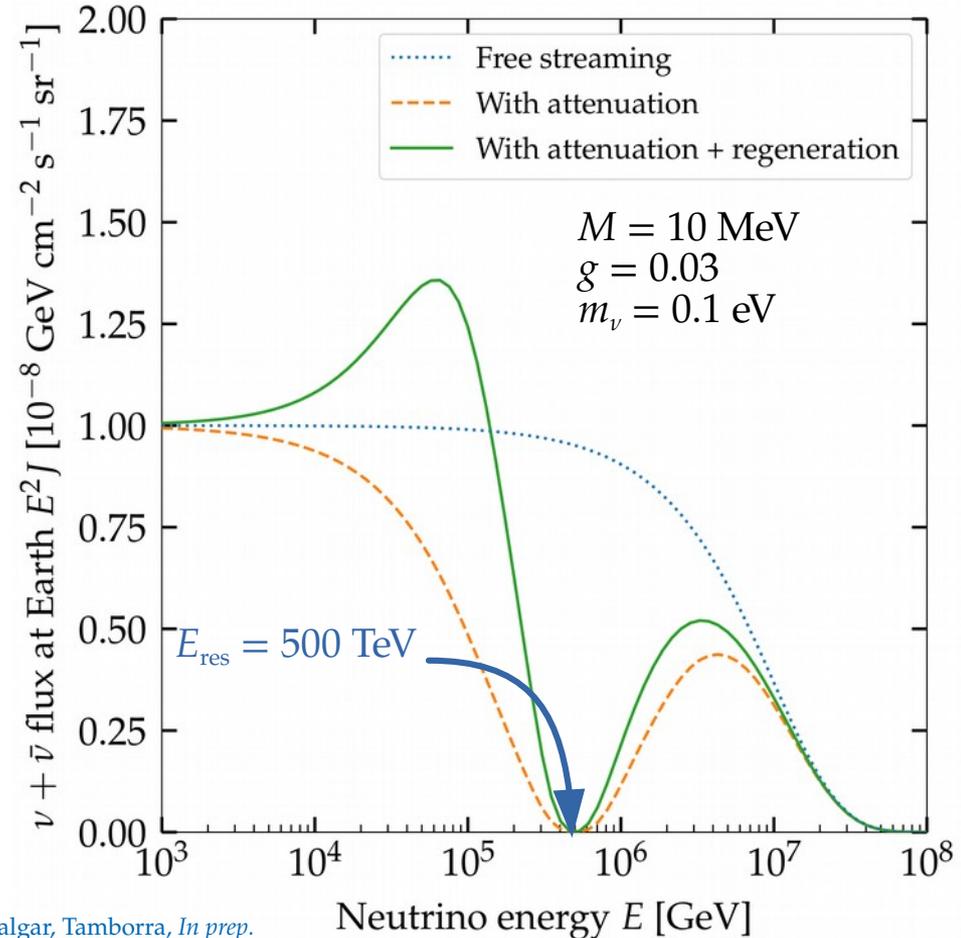
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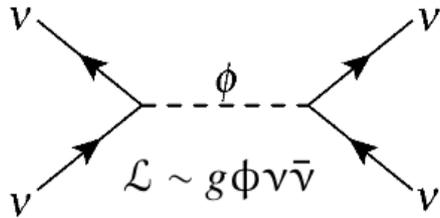
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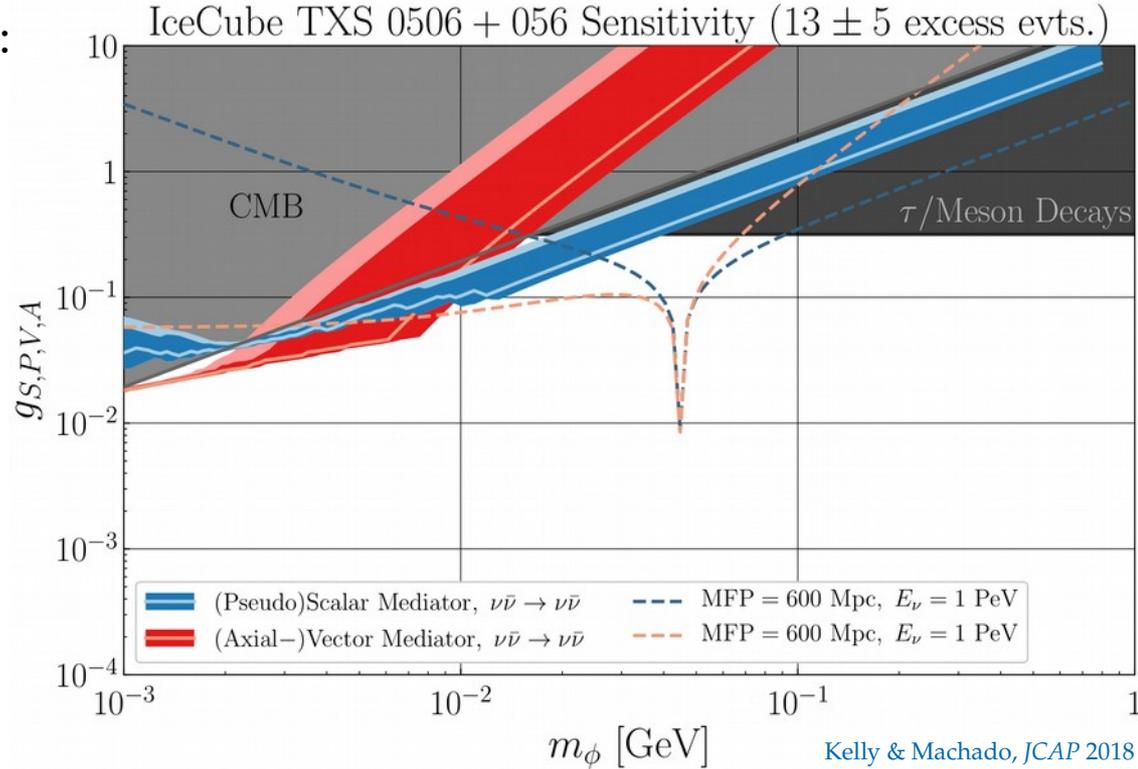
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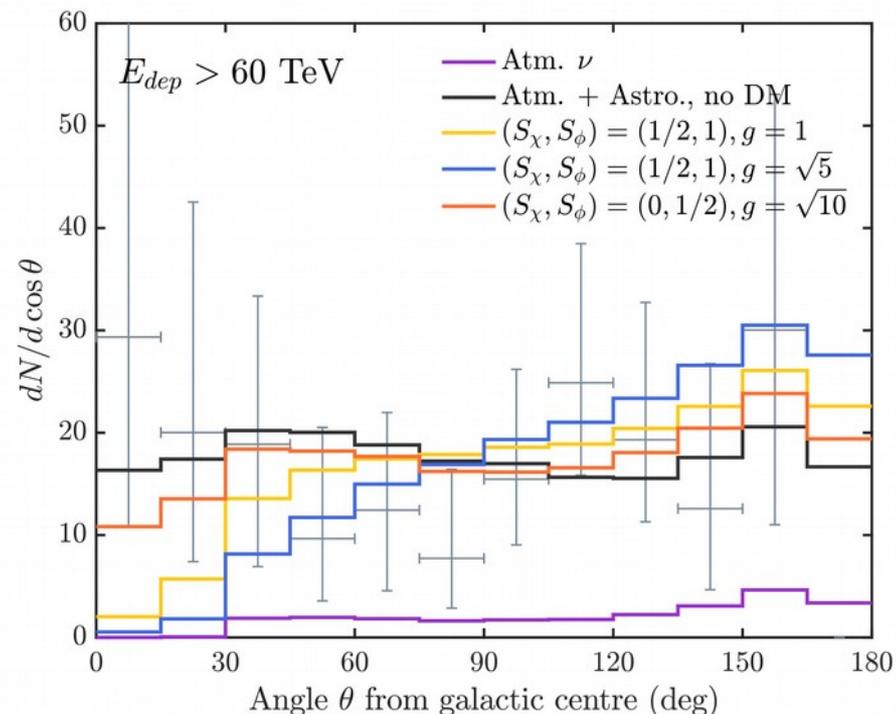
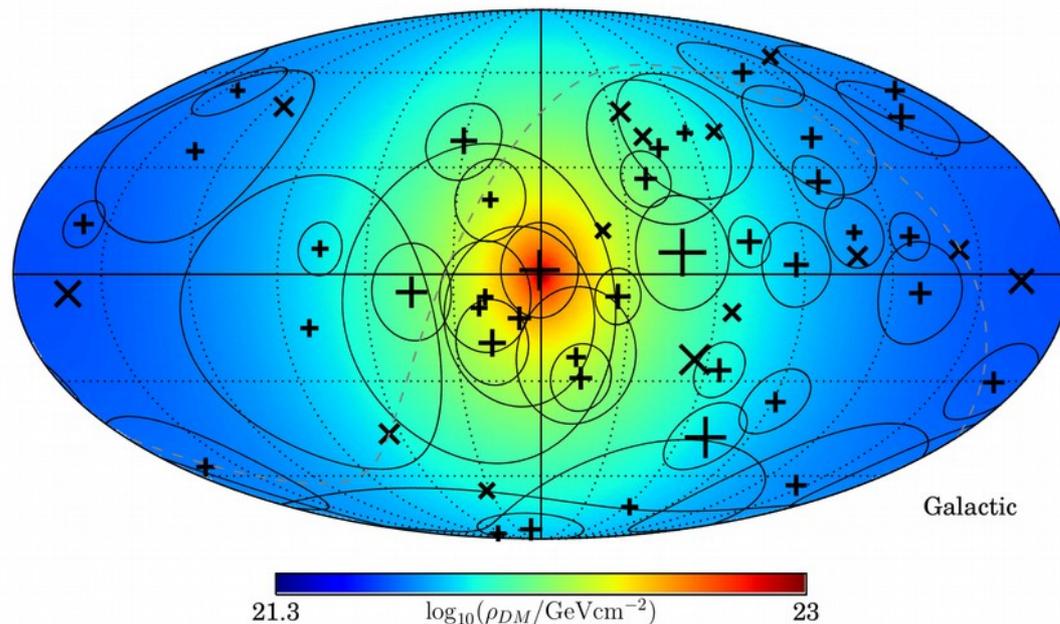
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Resonance energy:
$$E_{\text{res}} = \frac{M^2}{2m_\nu}$$



New physics in the angular distribution: ν -DM interactions

Interaction between astrophysical neutrinos and the Galactic dark matter profile –

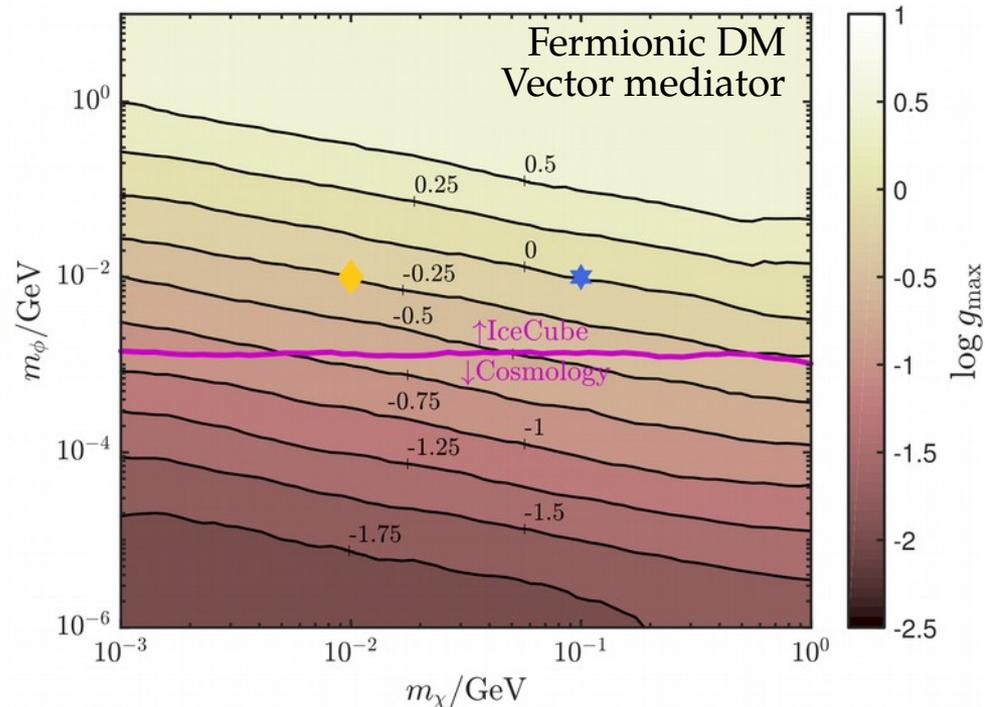
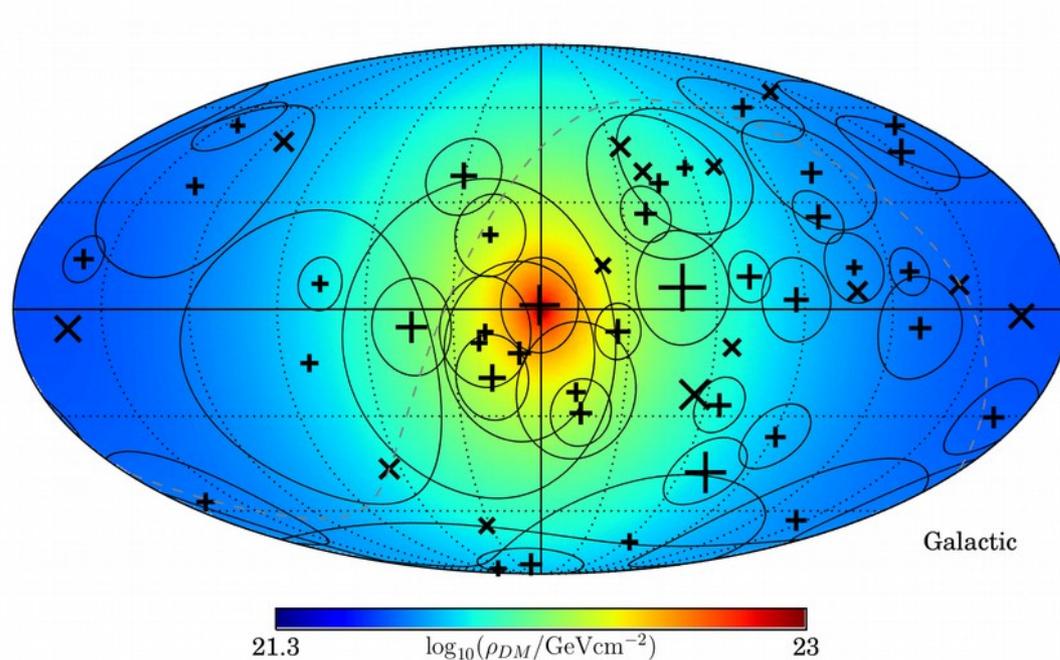


Expected: Fewer neutrinos coming from the Galactic Center

Observed: Isotropy

New physics in the angular distribution: ν -DM interactions

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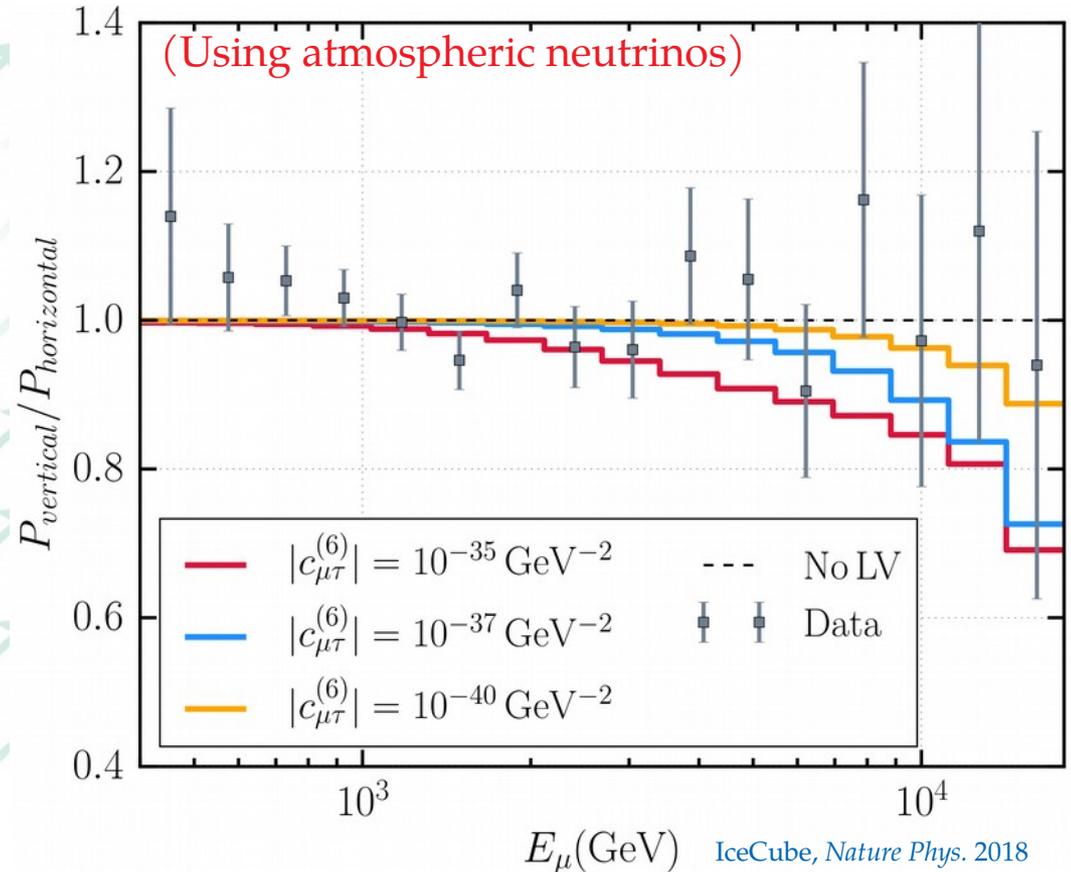
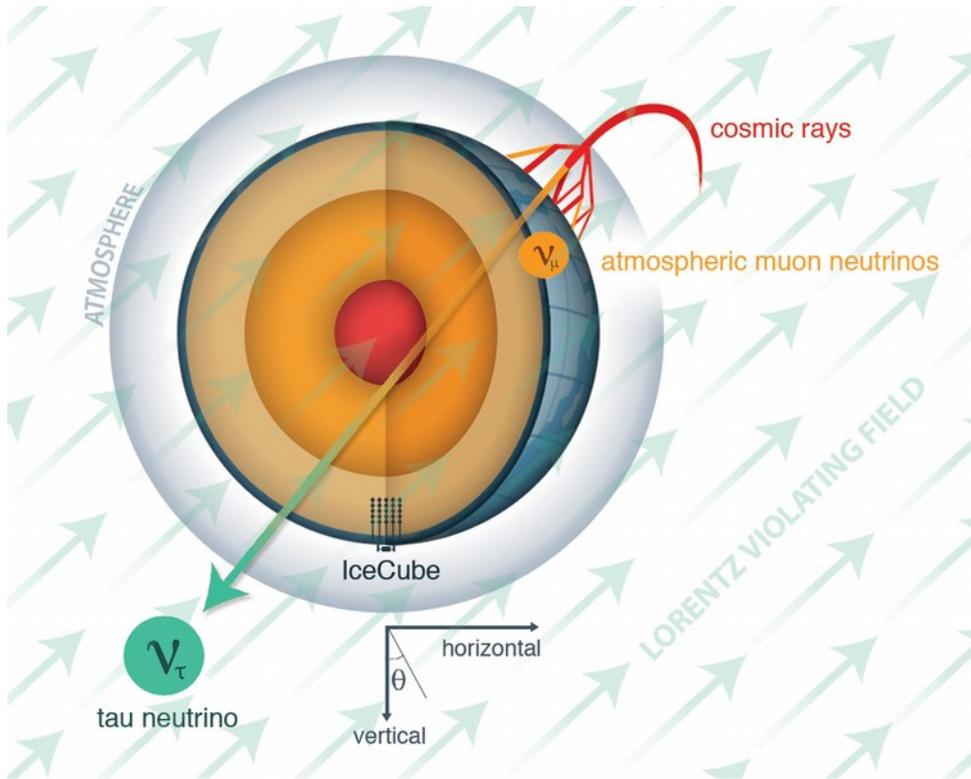


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New physics in the energy & angular distribution

Lorentz invariance violation – Hamiltonian: $H \sim m^2/(2E) + \hat{a}^{(3)} - E \cdot \hat{c}^{(4)} + E^2 \cdot \hat{a}^{(5)} - E^3 \cdot \hat{c}^{(6)}$

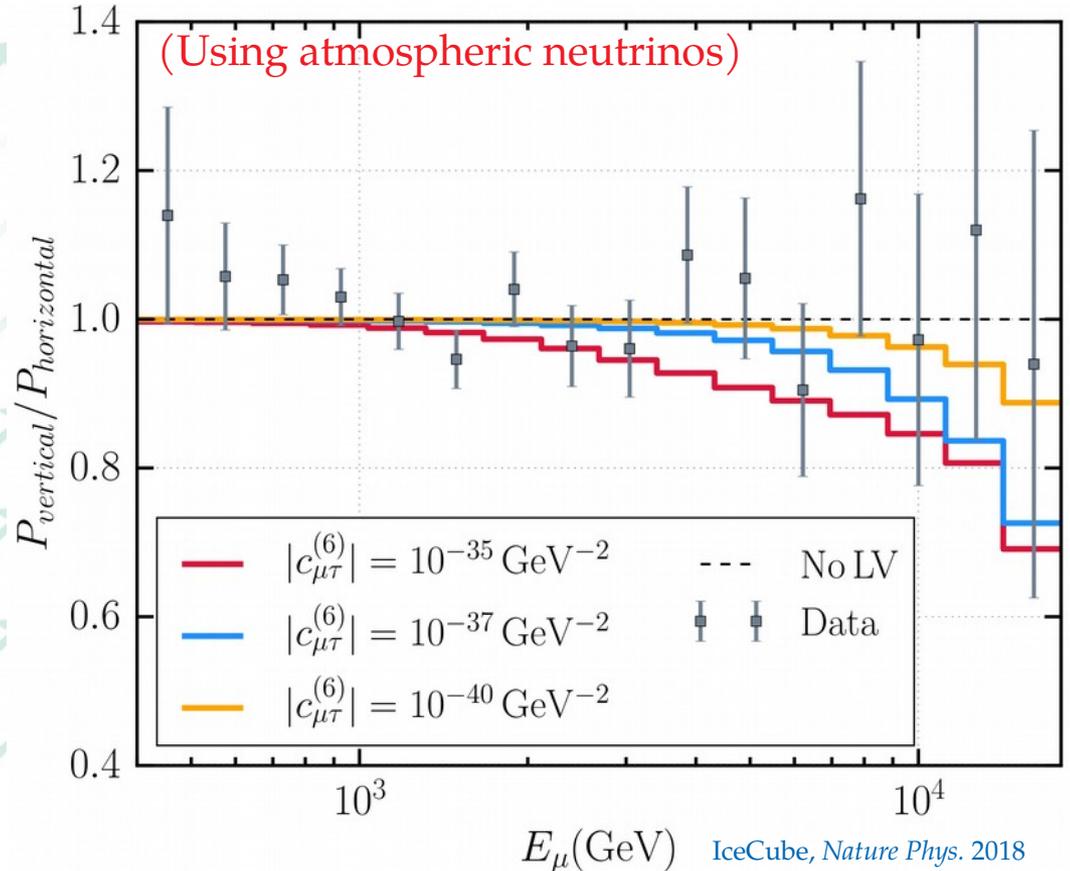
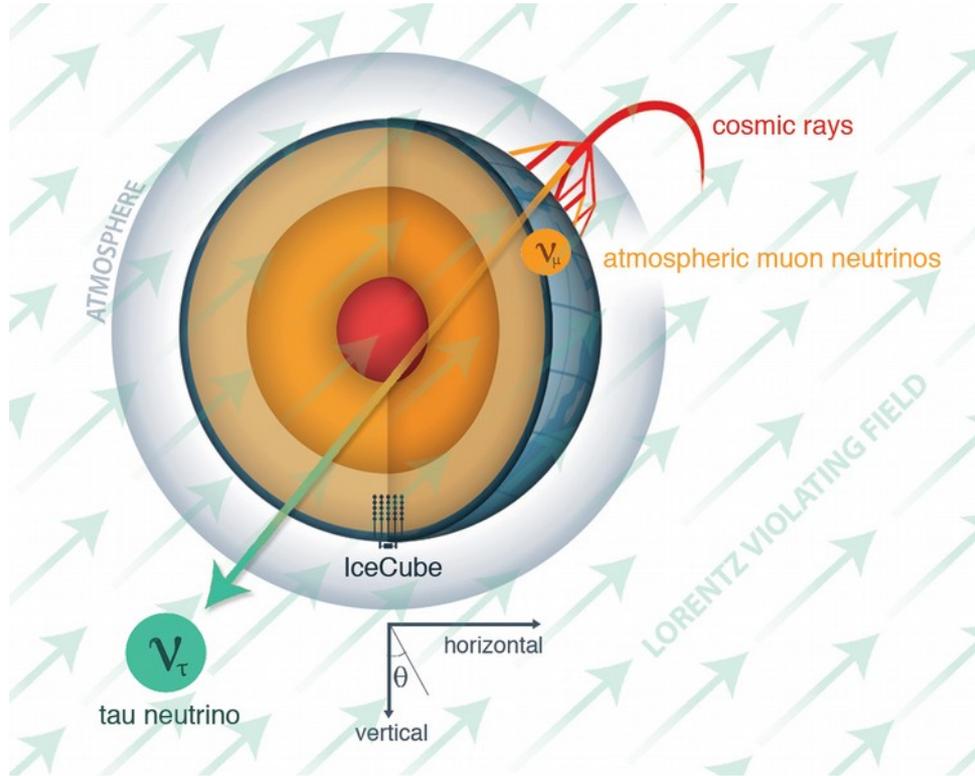


New physics in the energy & angular distribution

Lorentz violation

Standard oscillations

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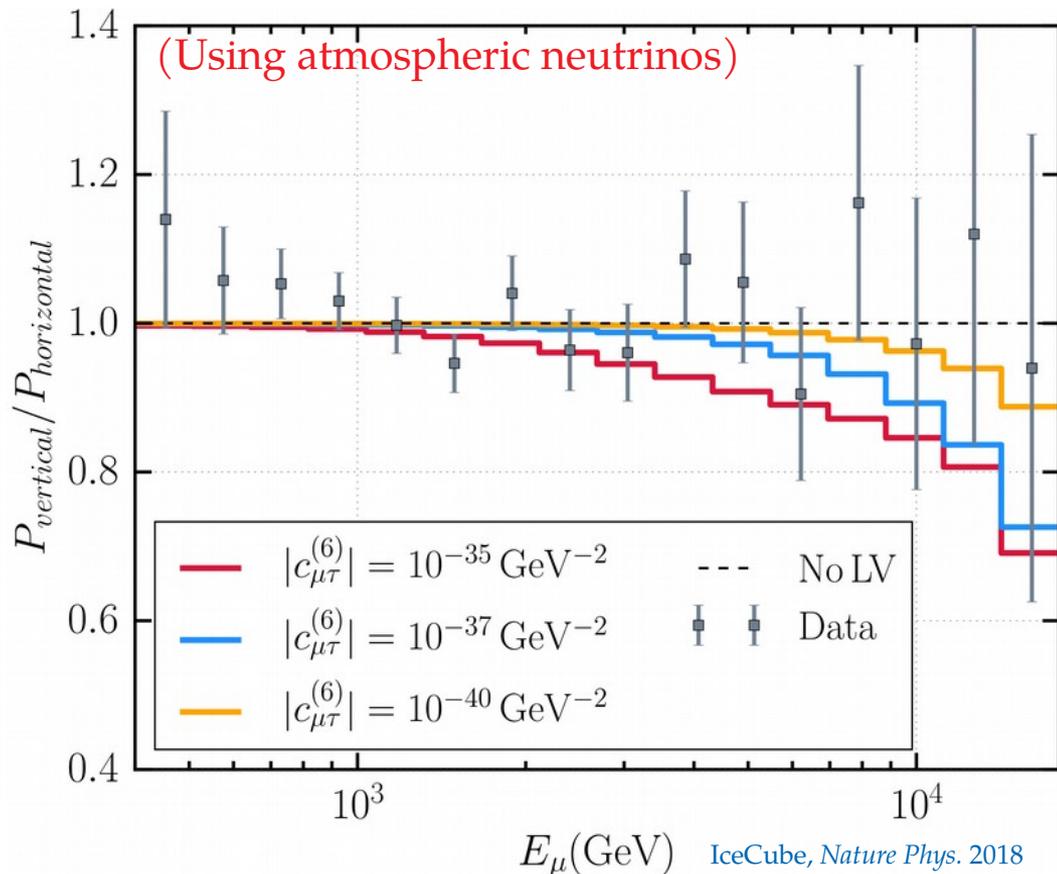
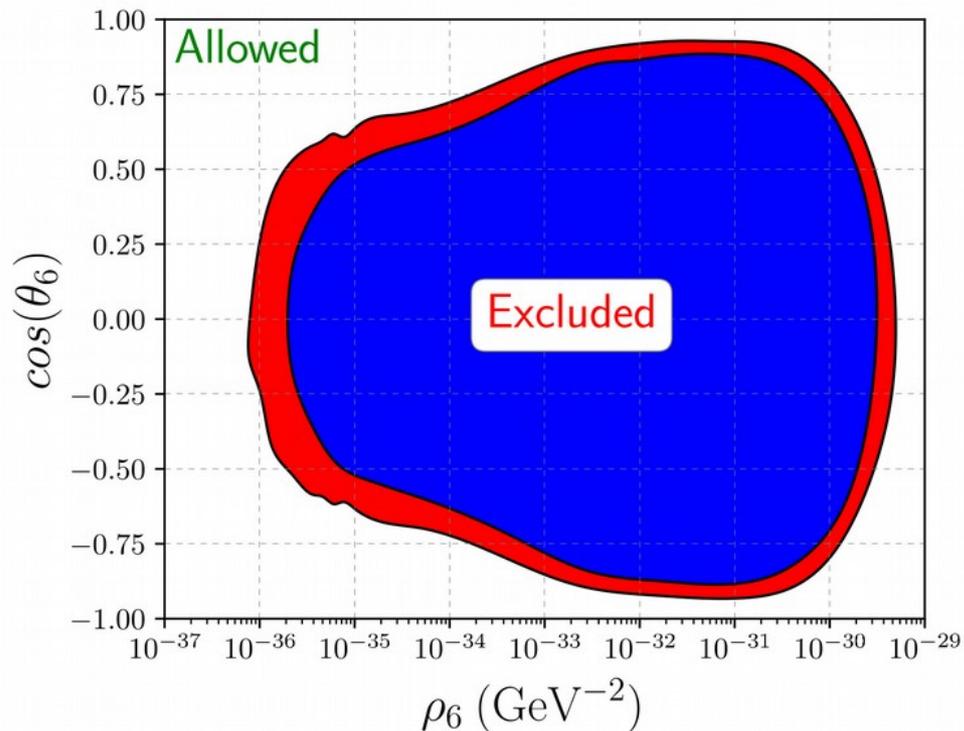
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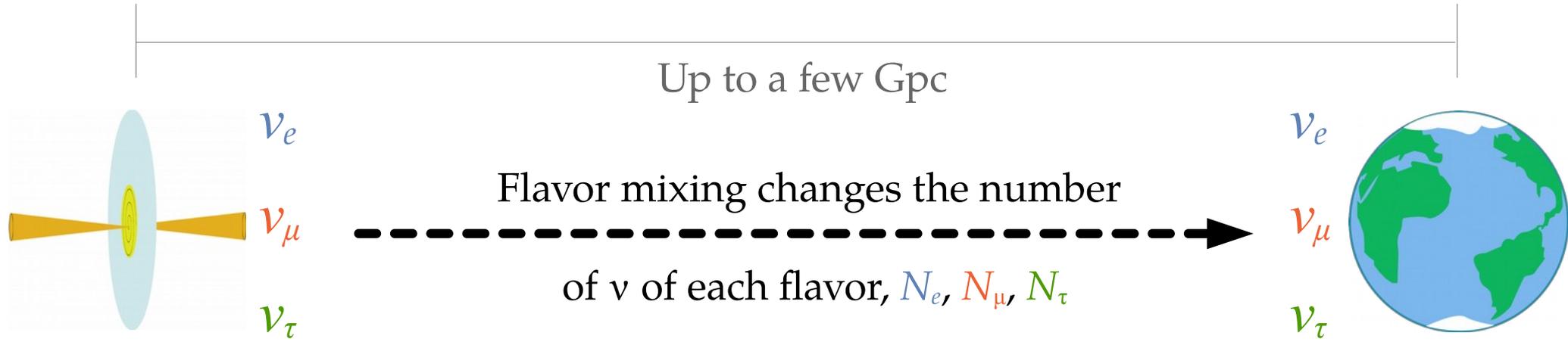
Best bounds come from IceCube



Flavor composition

Astrophysical neutrino sources

Earth



- ▶ Different processes yield different ratios of neutrinos of each flavor:

$$(f_{e,S}, f_{\mu,S}, f_{\tau,S}) \equiv (N_{e,S}, N_{\mu,S}, N_{\tau,S}) / N_{\text{tot}}$$

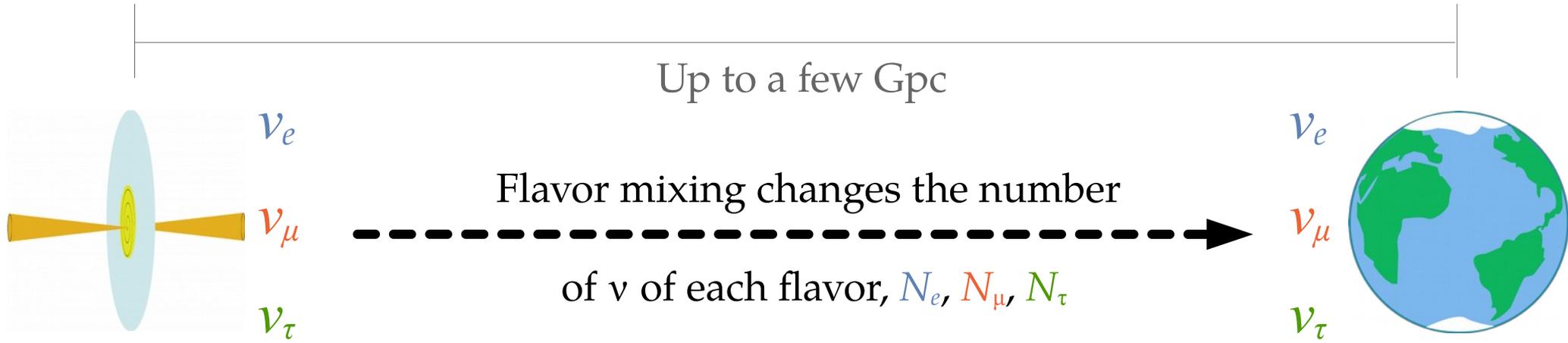
- ▶ Flavor ratios at Earth ($\alpha = e, \mu, \tau$):

$$f_{\alpha,\oplus} = \sum_{\beta=e,\mu,\tau} P_{\nu_\beta \rightarrow \nu_\alpha} f_{\beta,S}$$

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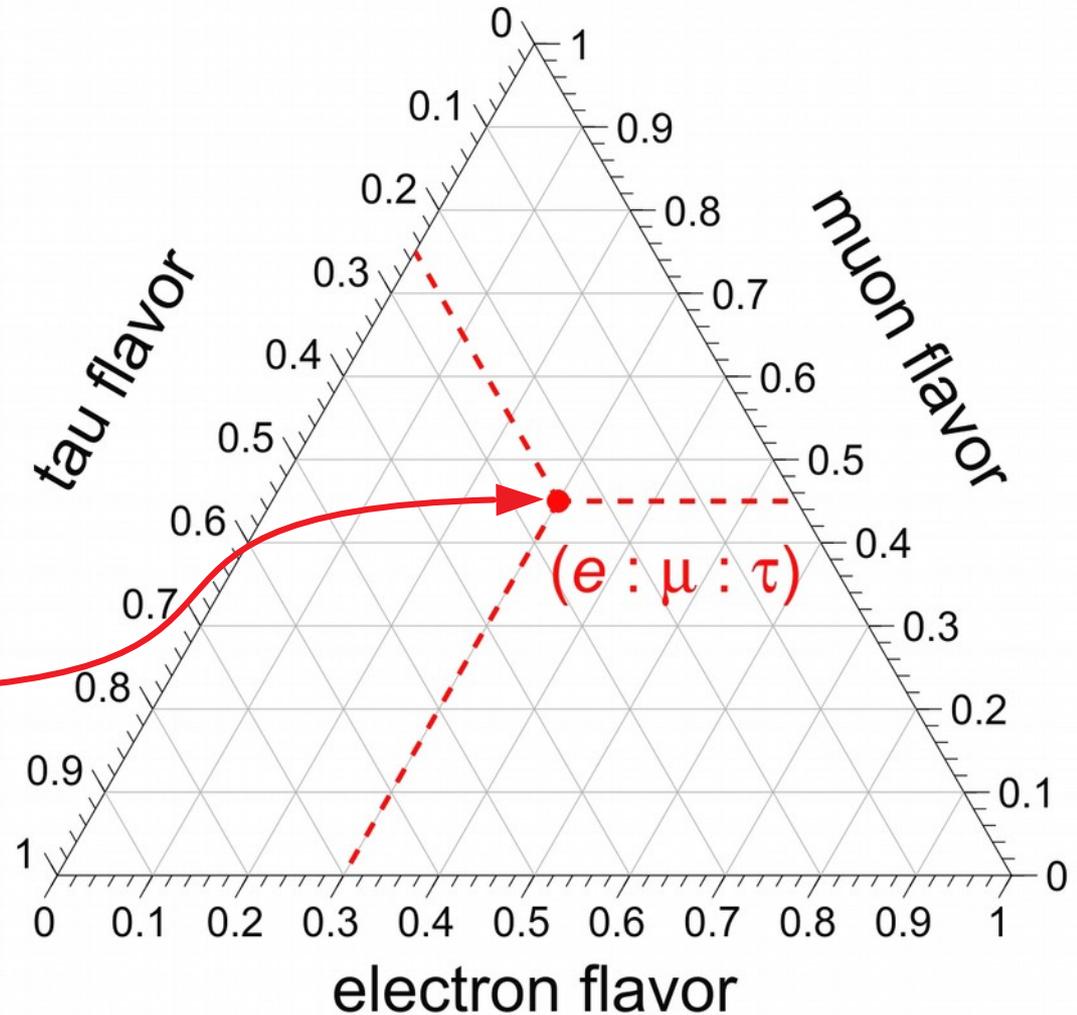
Standard oscillations
or
new physics

Reading a ternary plot

Assumes underlying unitarity –
sum of projections on each axis is 1

How to read it: Follow the tilt of
the tick marks, *e.g.*,

$$(e:\mu:\tau) = (0.30:0.45:0.25)$$



One likely TeV–PeV ν production scenario:

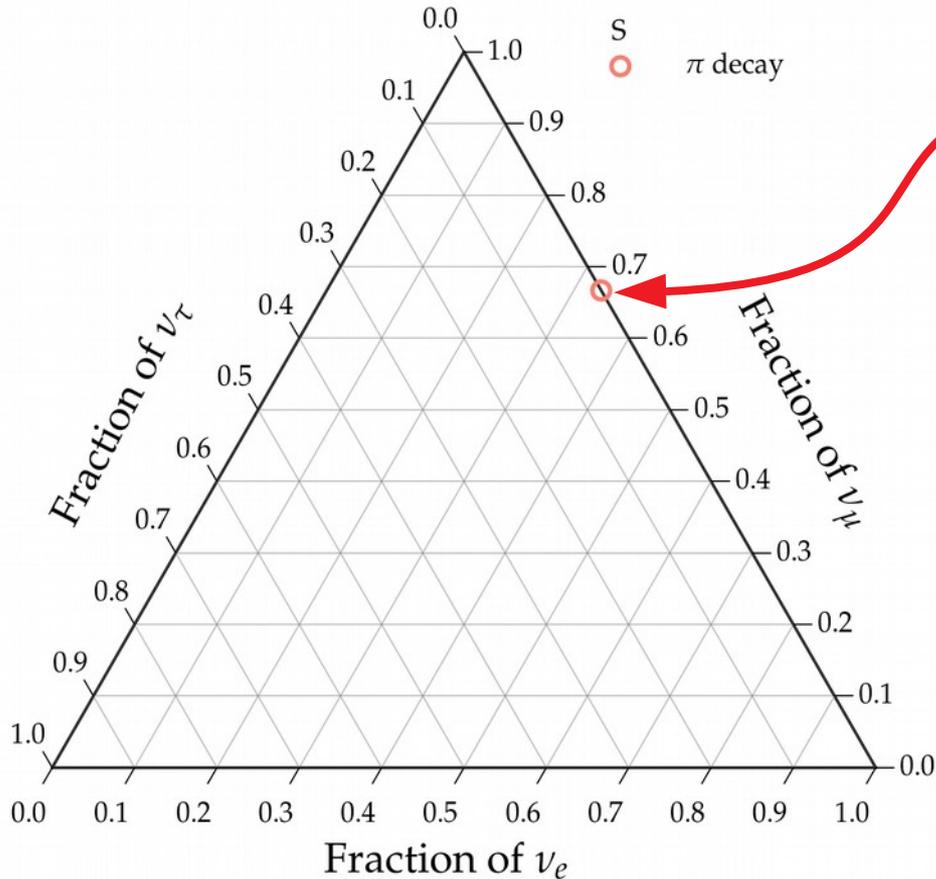
$$p + \gamma \rightarrow \pi^+ \rightarrow \mu^+ + \nu_\mu \quad \text{followed by} \quad \mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$$

Full π decay chain

$$(1/3:2/3:0)_S$$

Note: ν and $\bar{\nu}$ are (so far) indistinguishable
in neutrino telescopes

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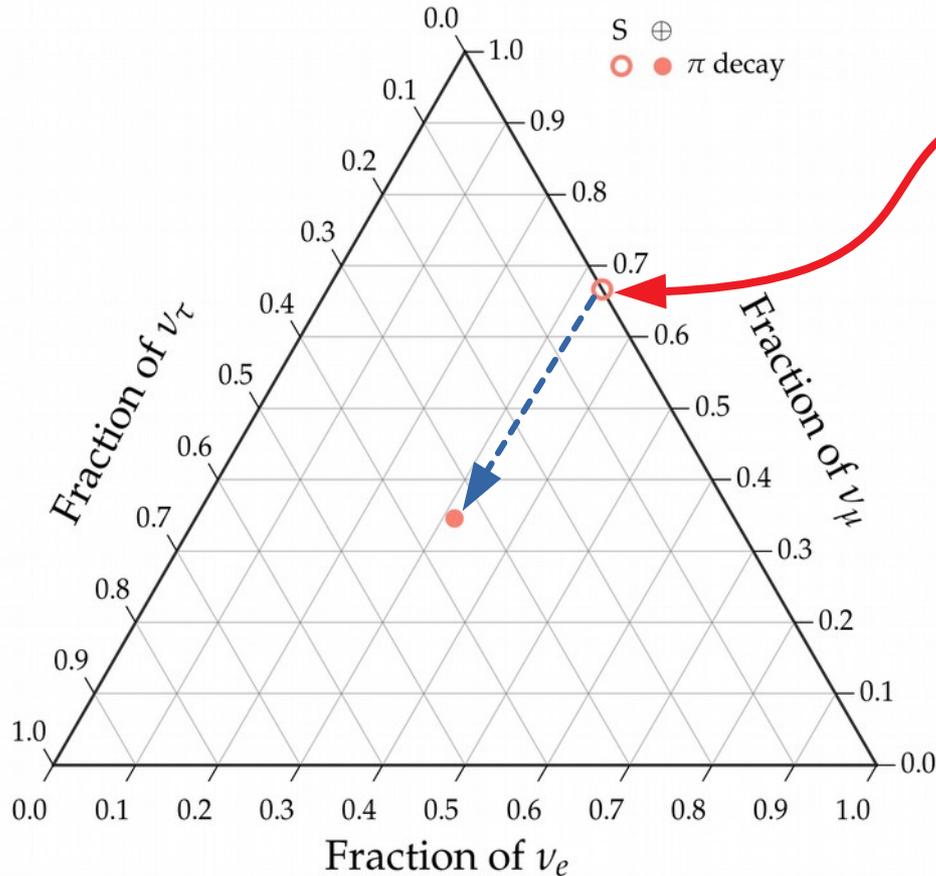


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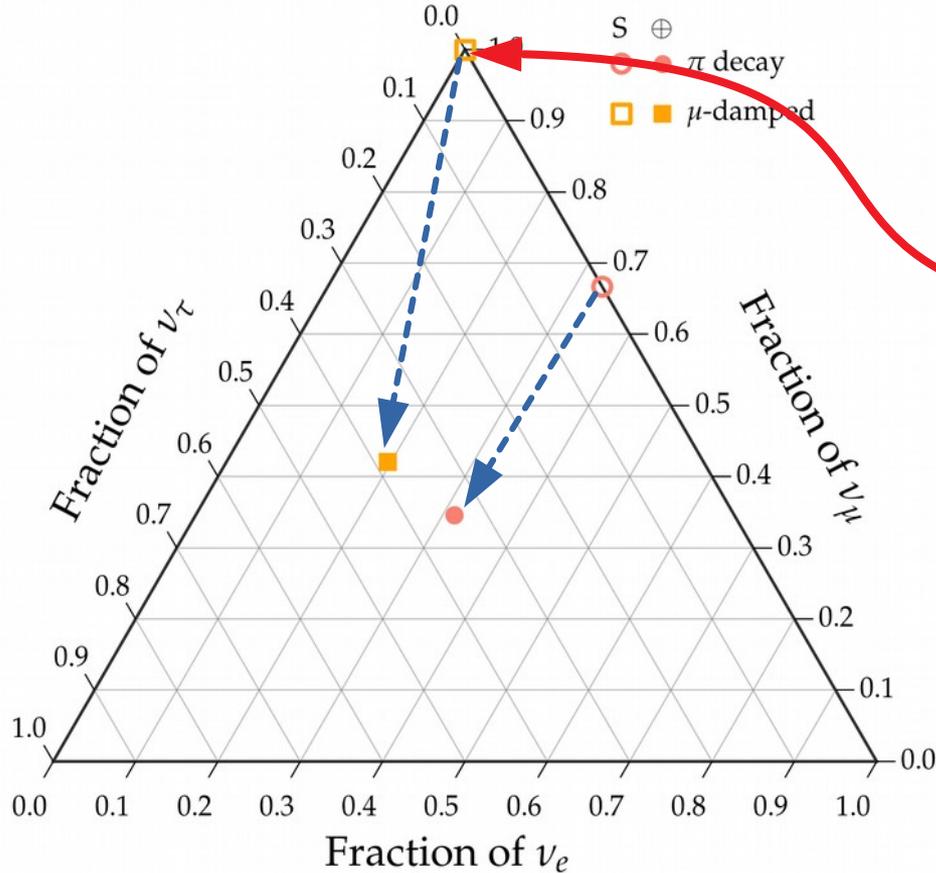


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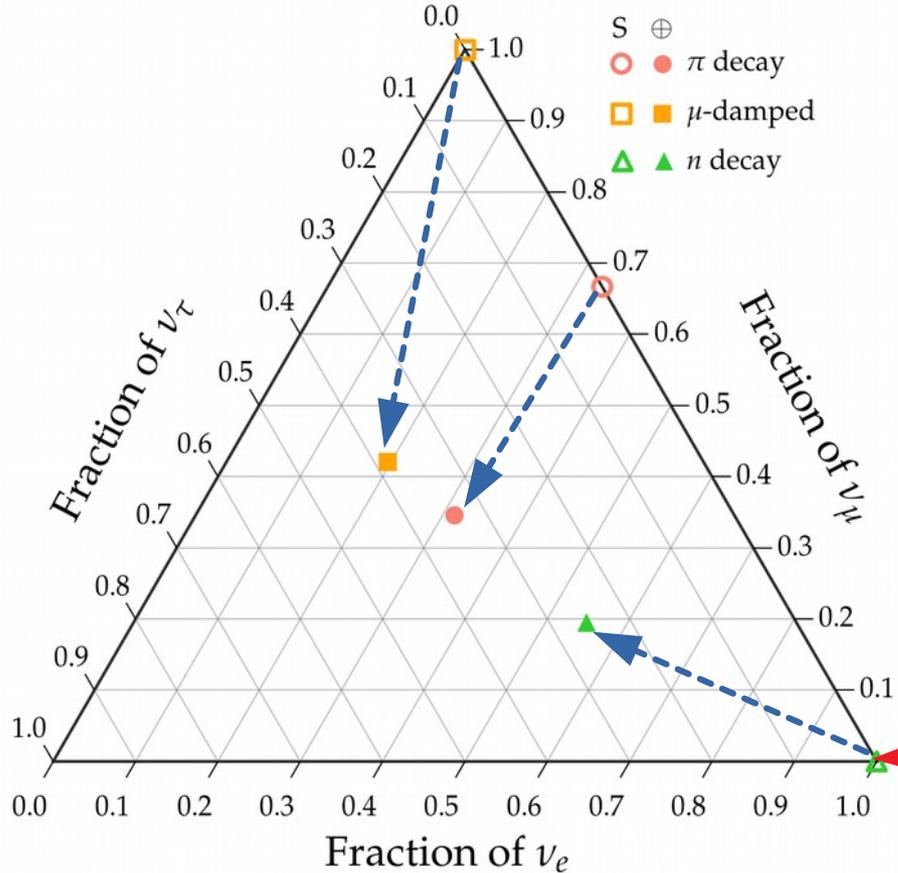
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Muon damped

$(0:1:0)_S$

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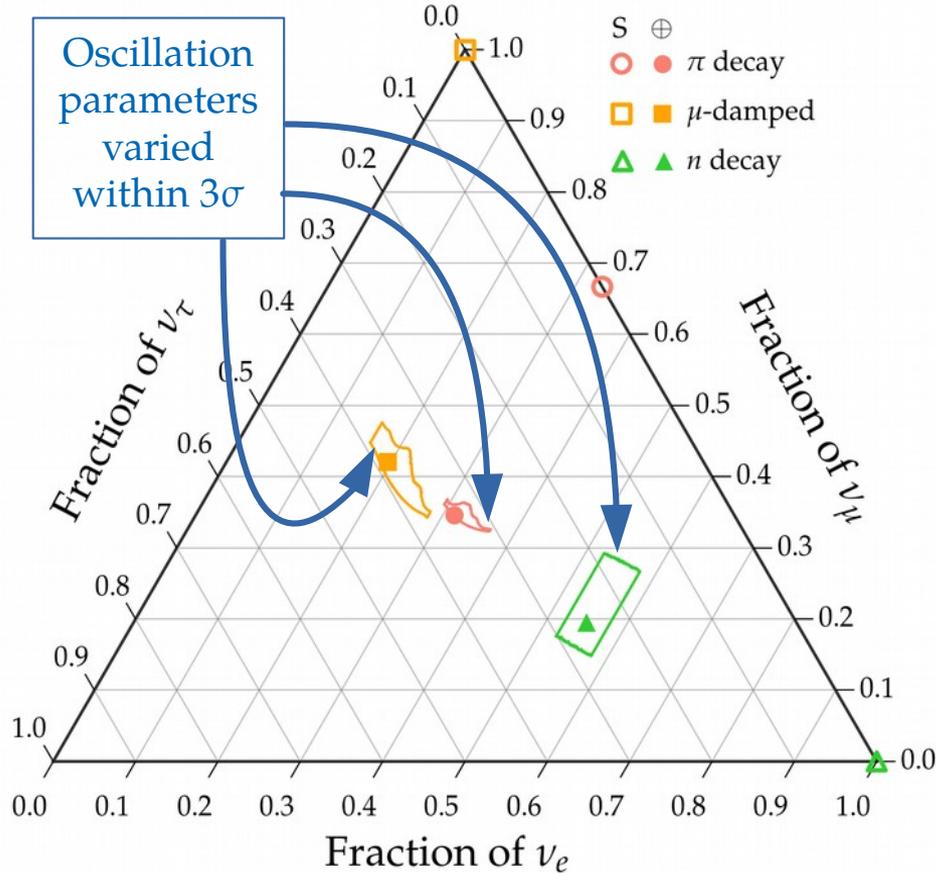
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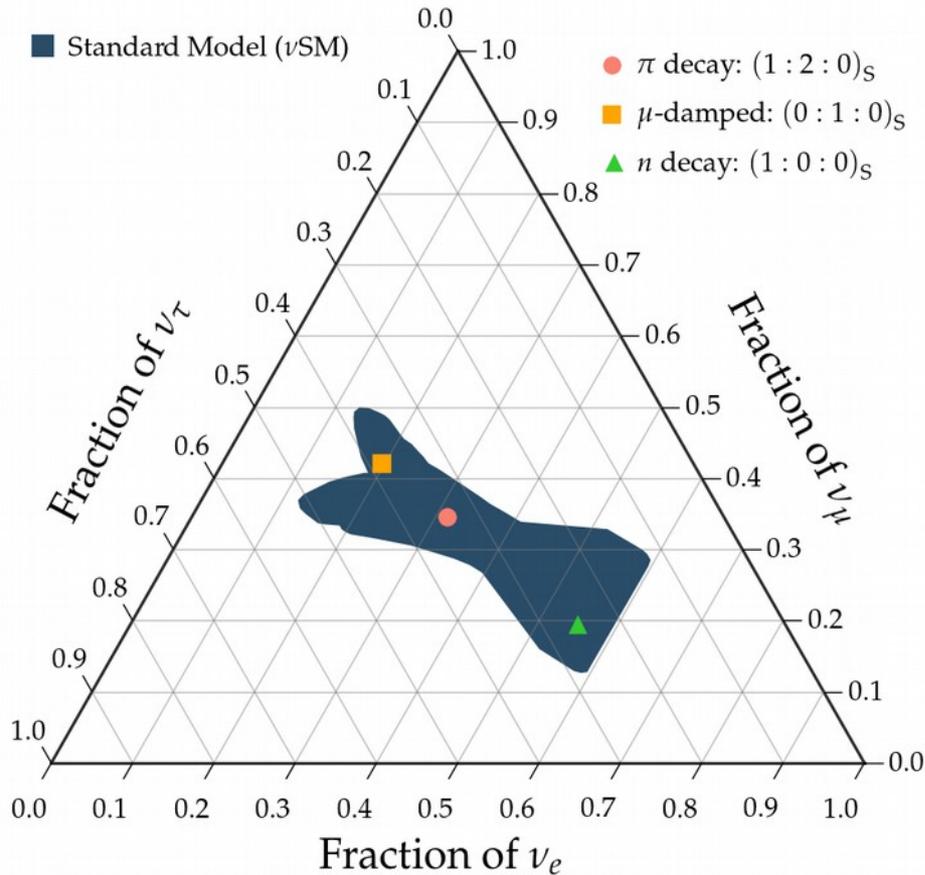
Muon damped

$$(0:1:0)_S$$

Neutron decay

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Note: ν and $\bar{\nu}$ are (so far) indistinguishable in neutrino telescopes



All possible flavor ratios at the sources

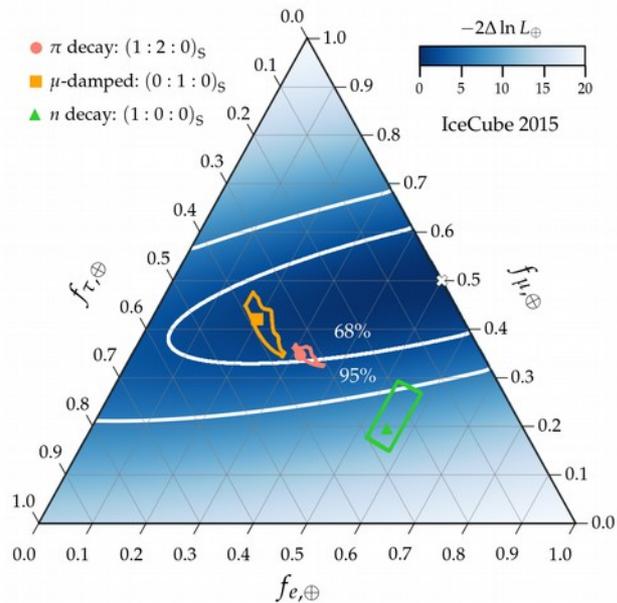
+

Vary oscillation parameters within 3σ

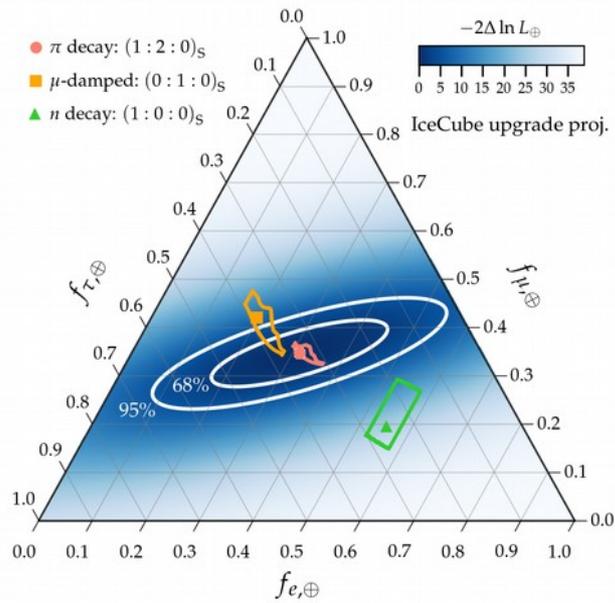
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IceCube flavor composition

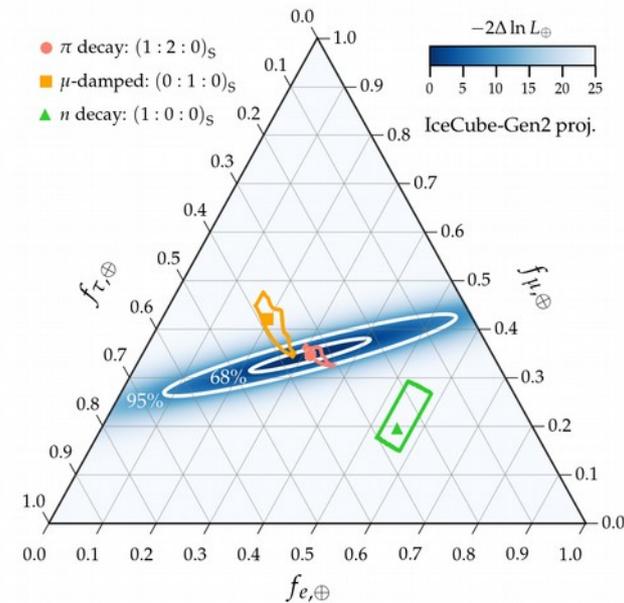
Today
IceCube



Near future (2022)
IceCube upgrade



In 10 years (2030s)
IceCube-Gen2

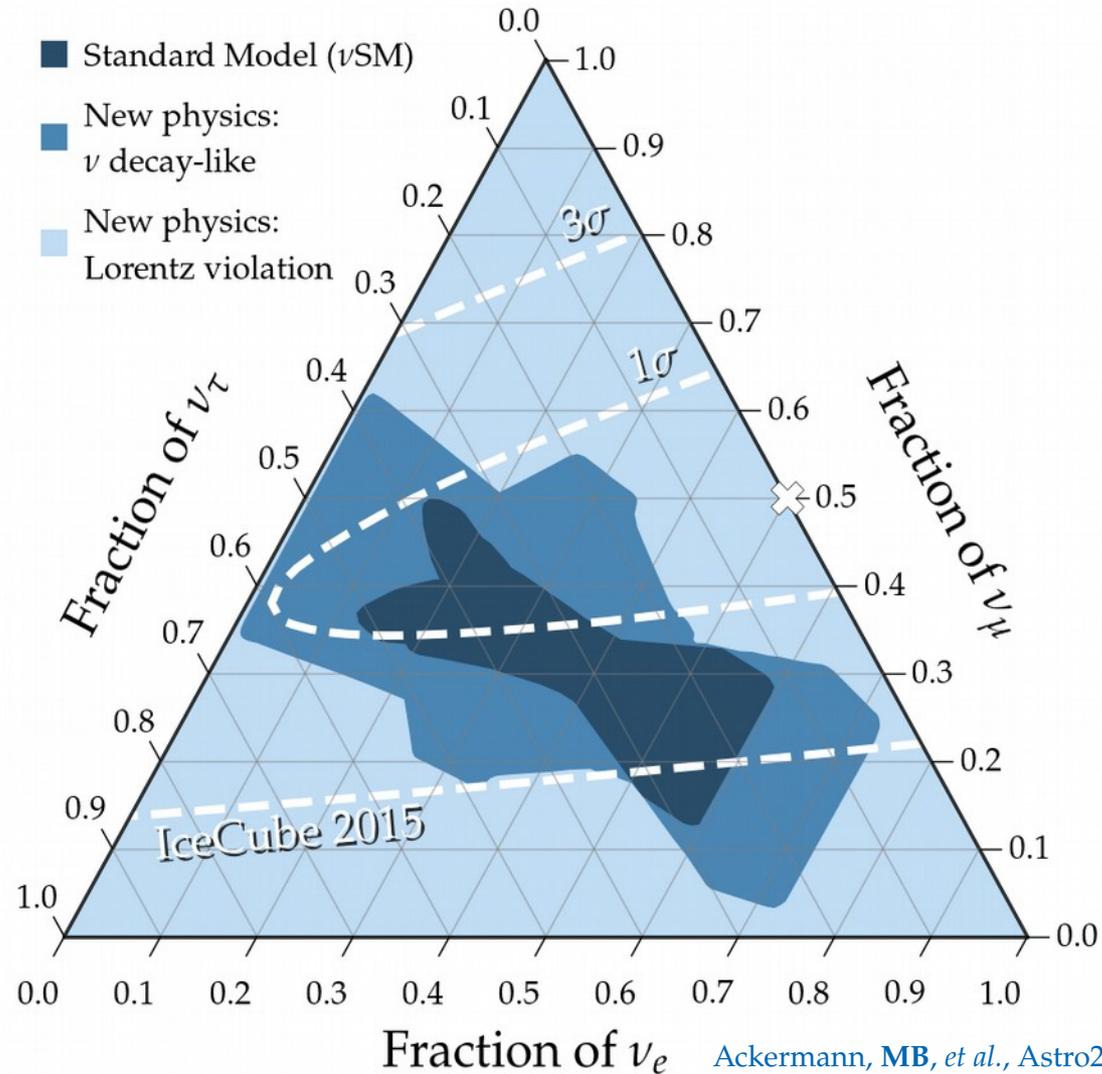


- ▶ Best fit:
 $(f_e : f_\mu : f_\tau)_\oplus = (0.49 : 0.51 : 0)_\oplus$
- ▶ Compatible with standard source compositions
- ▶ Hints of one ν_τ (not shown)

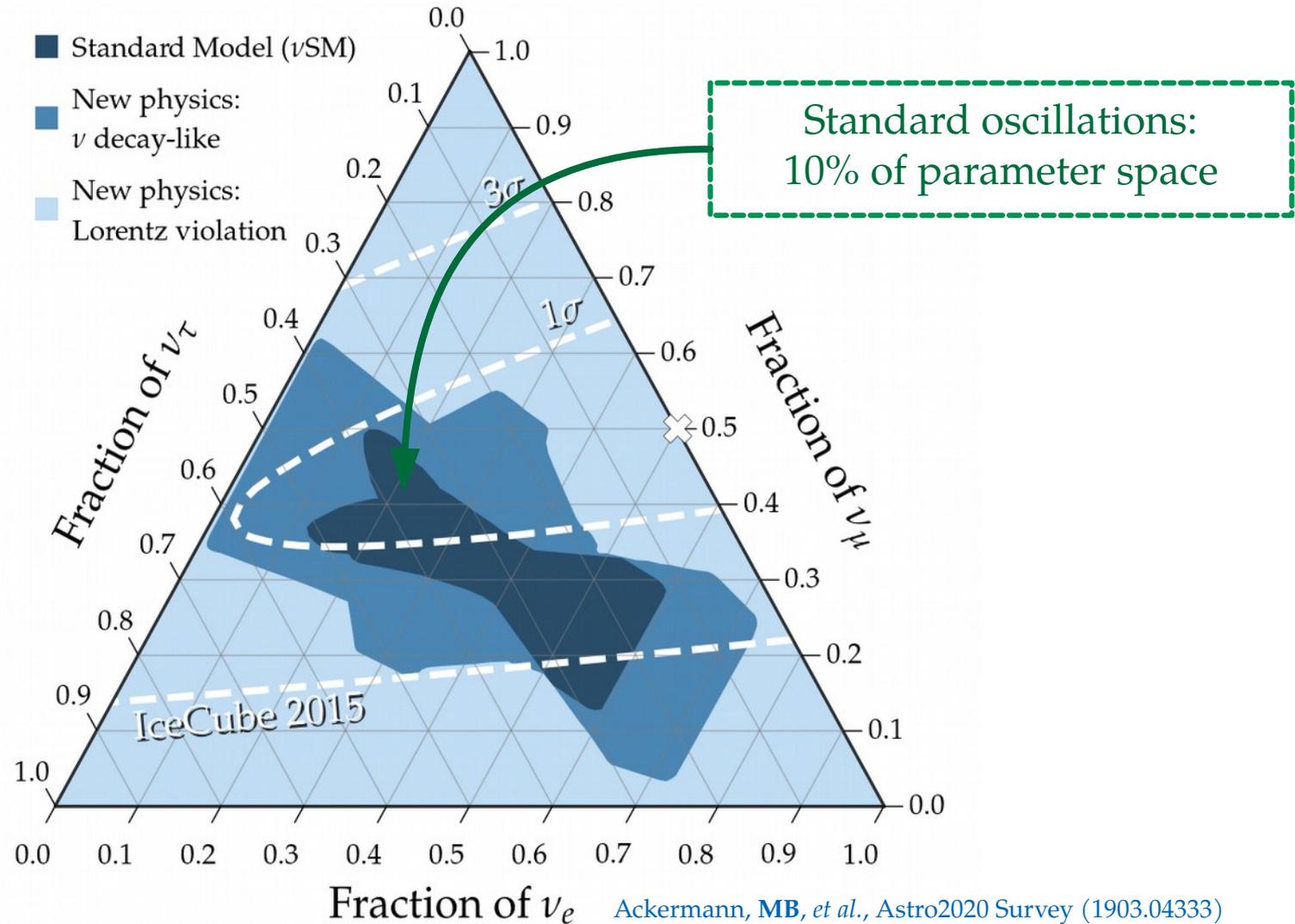
Assuming production by the full pion decay chain

Plus possibly better flavor-tagging, *e.g.*, muon and neutron echoes

[Li, MB, Beacom PRL 2019]



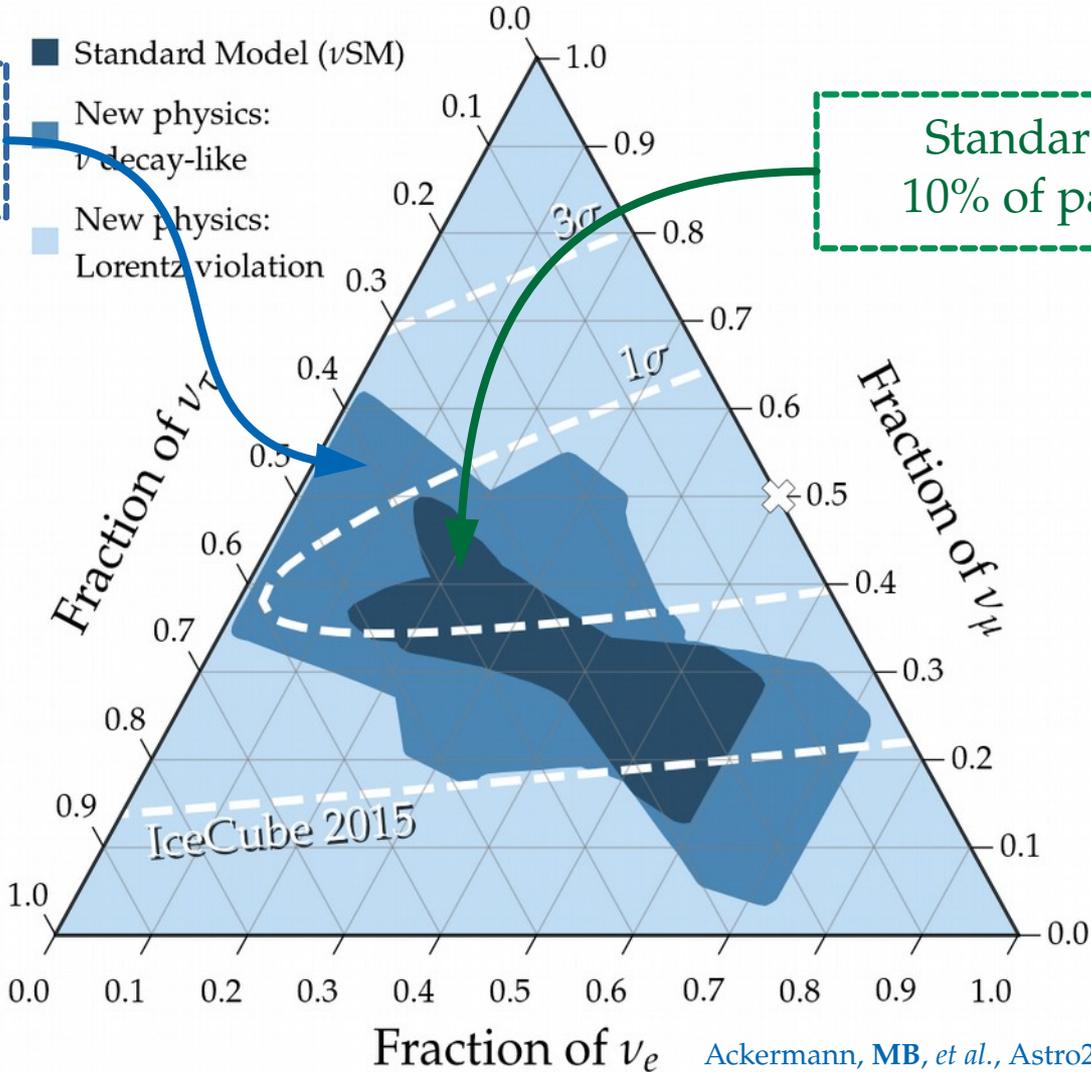
Ackermann, MB, et al., Astro2020 Survey (1903.04333)
 Based on: MB, Beacom, Winter PRL 2015



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Neutrino decay
30% of parameter space

Standard oscillations:
10% of parameter space



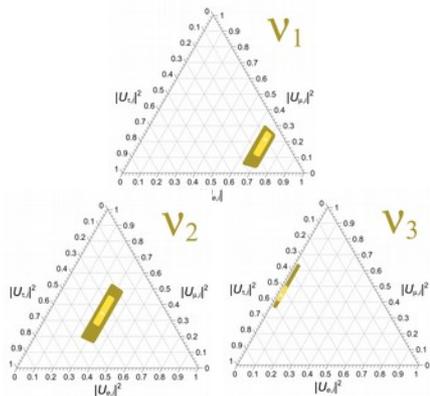
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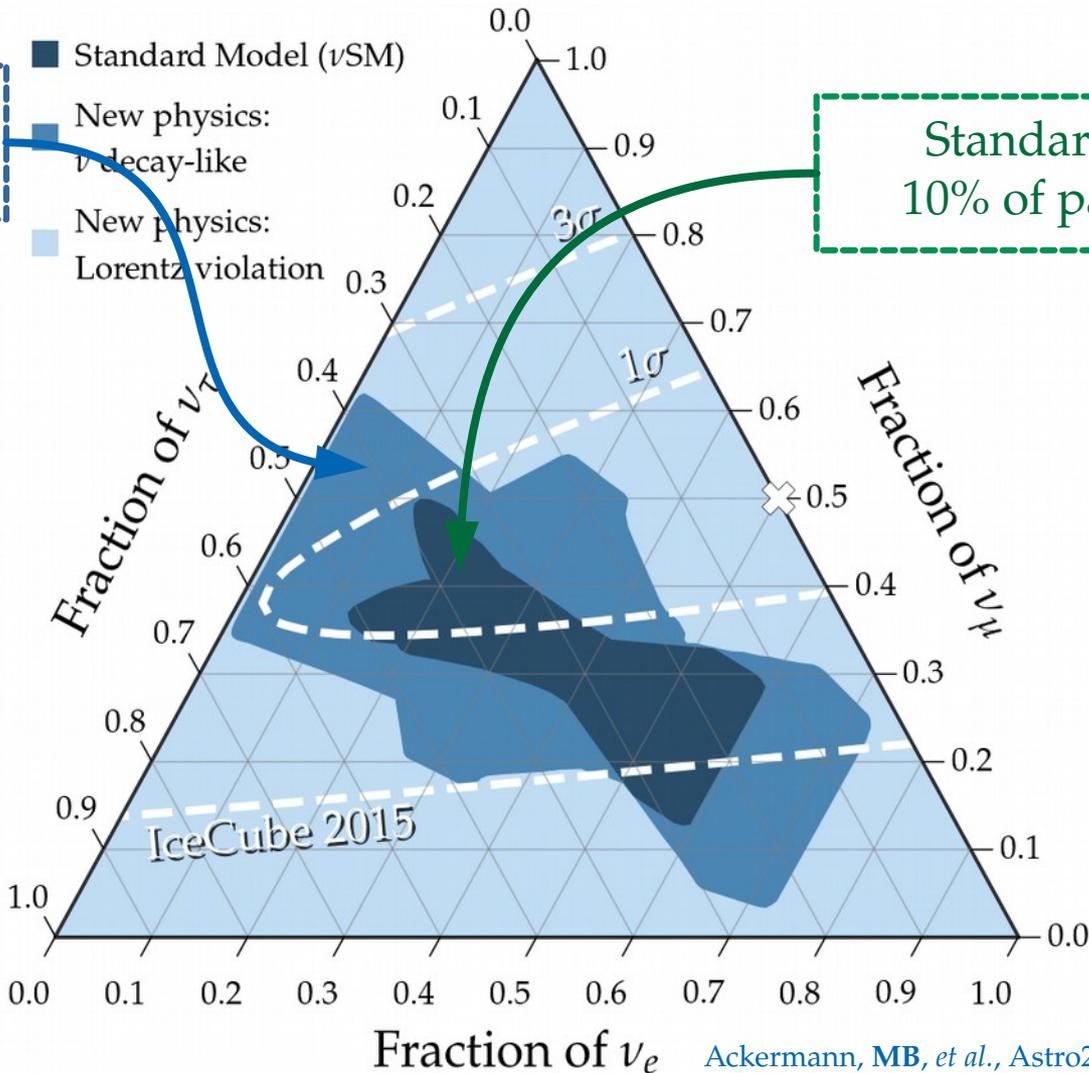
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$\nu_2, \nu_3 \rightarrow \nu_1$ OR $\nu_1, \nu_2 \rightarrow \nu_3$

Flavor ratios determined by
how many ν_1, ν_2, ν_3 survive:



$\tau_2/m_2, \tau_3/m_3 > 10 \text{ s eV}^{-1}$



MB, Beacom, Murase PRD 2017
Baerwald, MB, Winter JCAP 2012

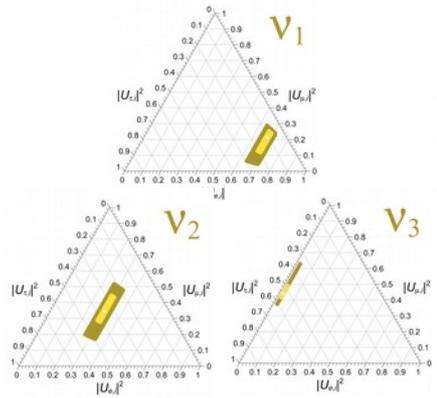
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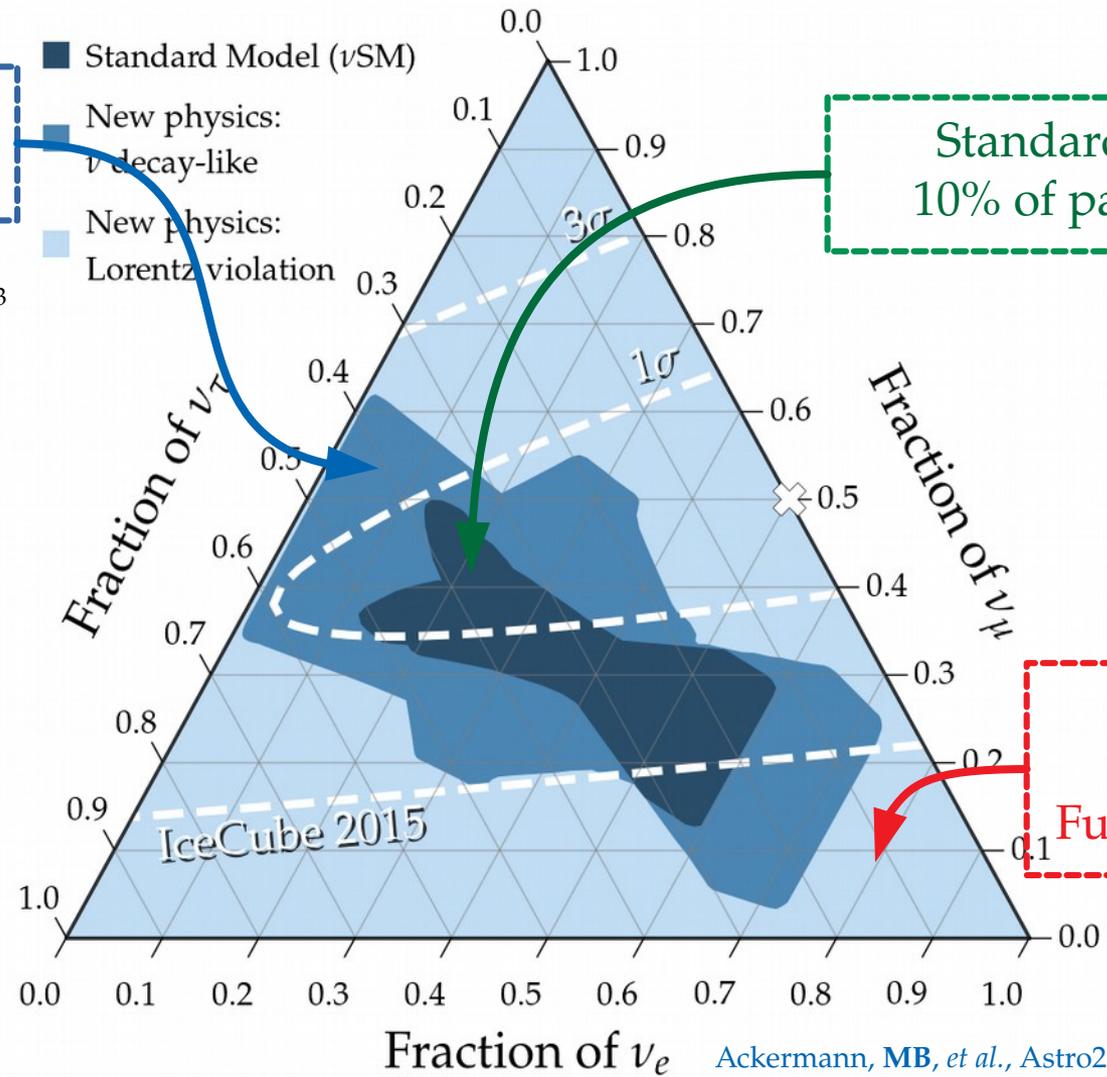
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- Standard Model (ν SM)
- New physics:
 ν decay-like
- New physics:
Lorentz violation



Lorentz, CPT
violation, etc:
Full parameter space

MB, Beacom, Murase PRD 2017
Baerwald, MB, Winter JCAP 2012

Fraction of ν_e

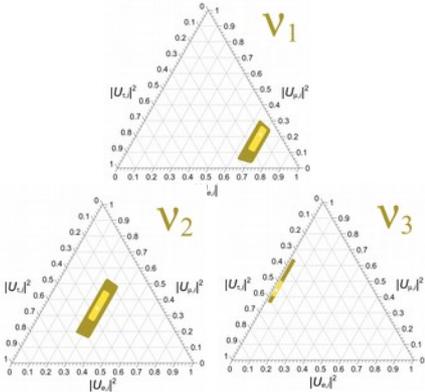
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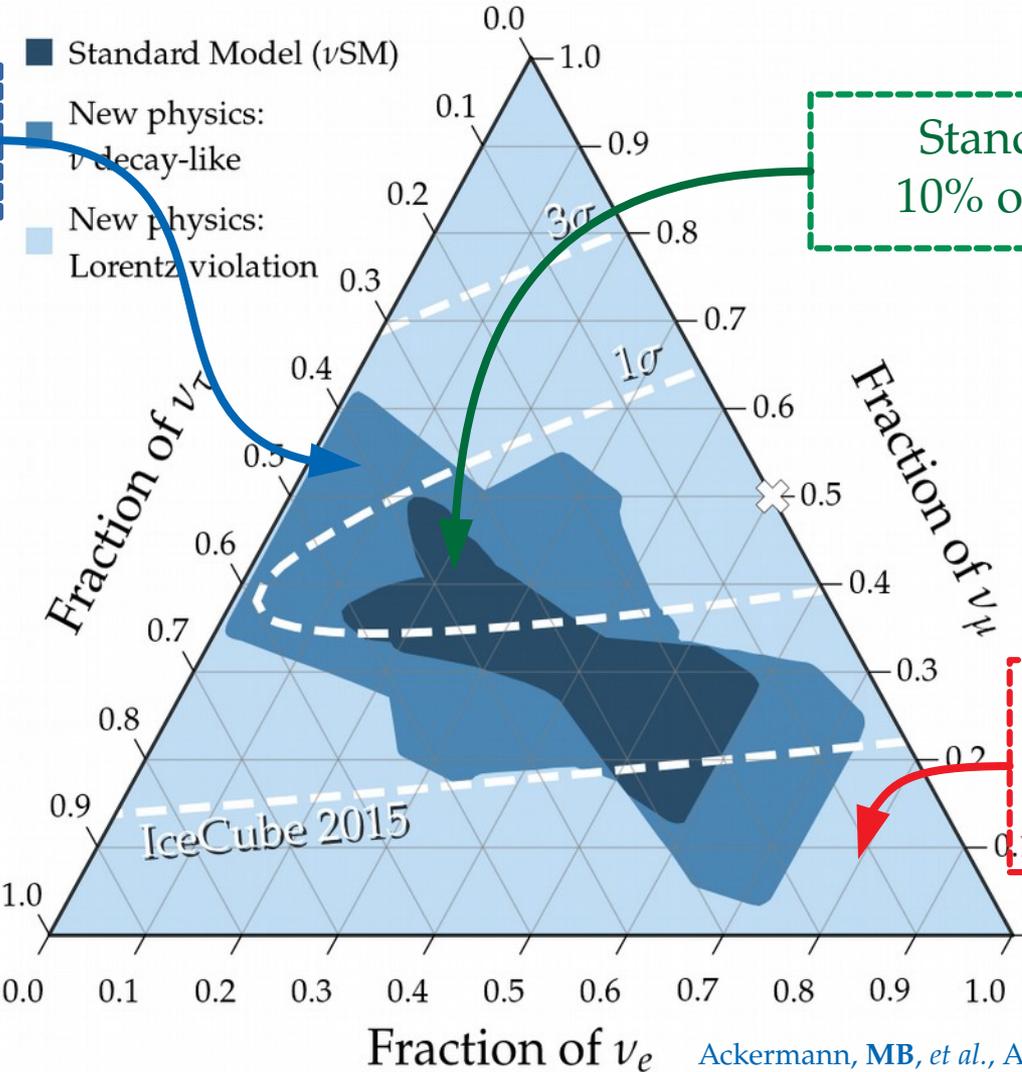
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How to fill out the flavor triangle?

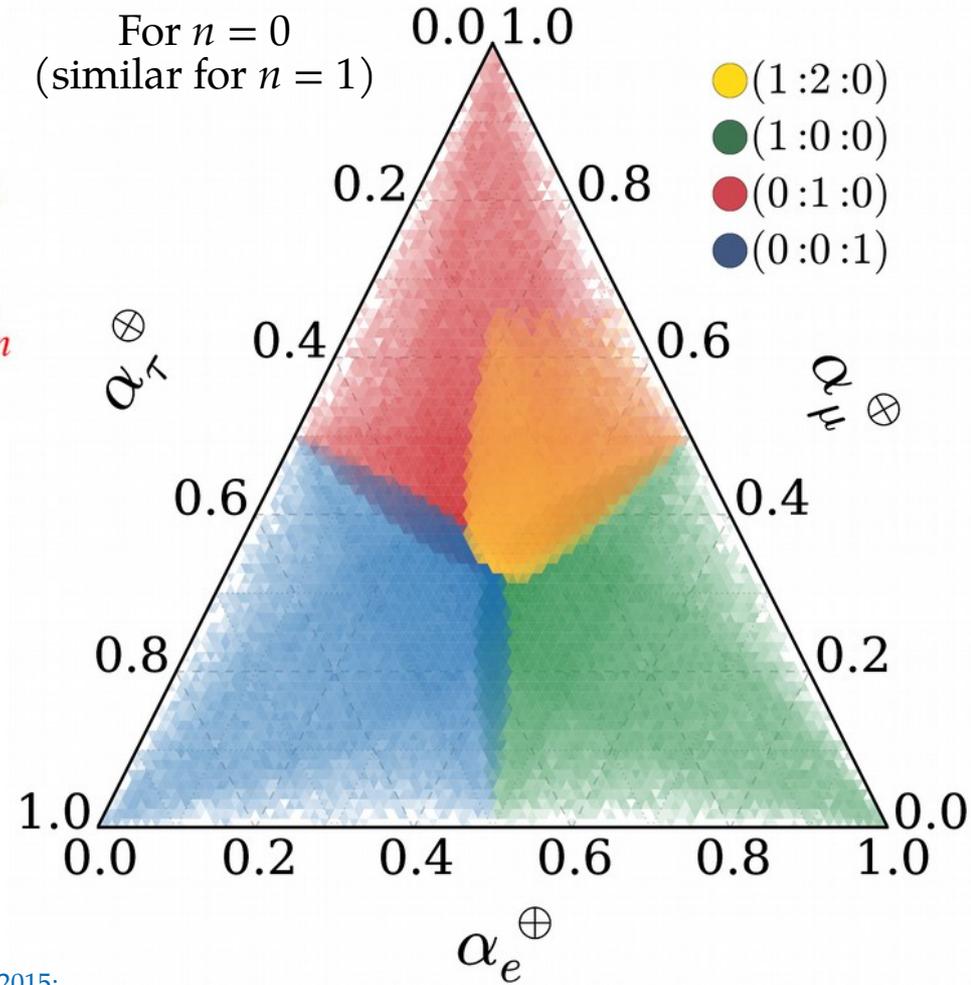
$$H_{\text{tot}} = H_{\text{std}} + H_{\text{NP}}$$

$$H_{\text{std}} = \frac{1}{2E} U_{\text{PMNS}}^\dagger \text{diag} (0, \Delta m_{21}^2, \Delta m_{31}^2) U_{\text{PMNS}}$$

$$H_{\text{NP}} = \sum_n \left(\frac{E}{\Lambda_n} \right)^n U_n^\dagger \text{diag} (O_{n,1}, O_{n,2}, O_{n,3}) U_n$$

This can populate *all* of the triangle –

- ▶ Use current atmospheric bounds on $O_{n,i}$:
 $O_0 < 10^{-23}$ GeV, $O_1/\Lambda_1 < 10^{-27}$ GeV
- ▶ Sample the unknown new mixing angles



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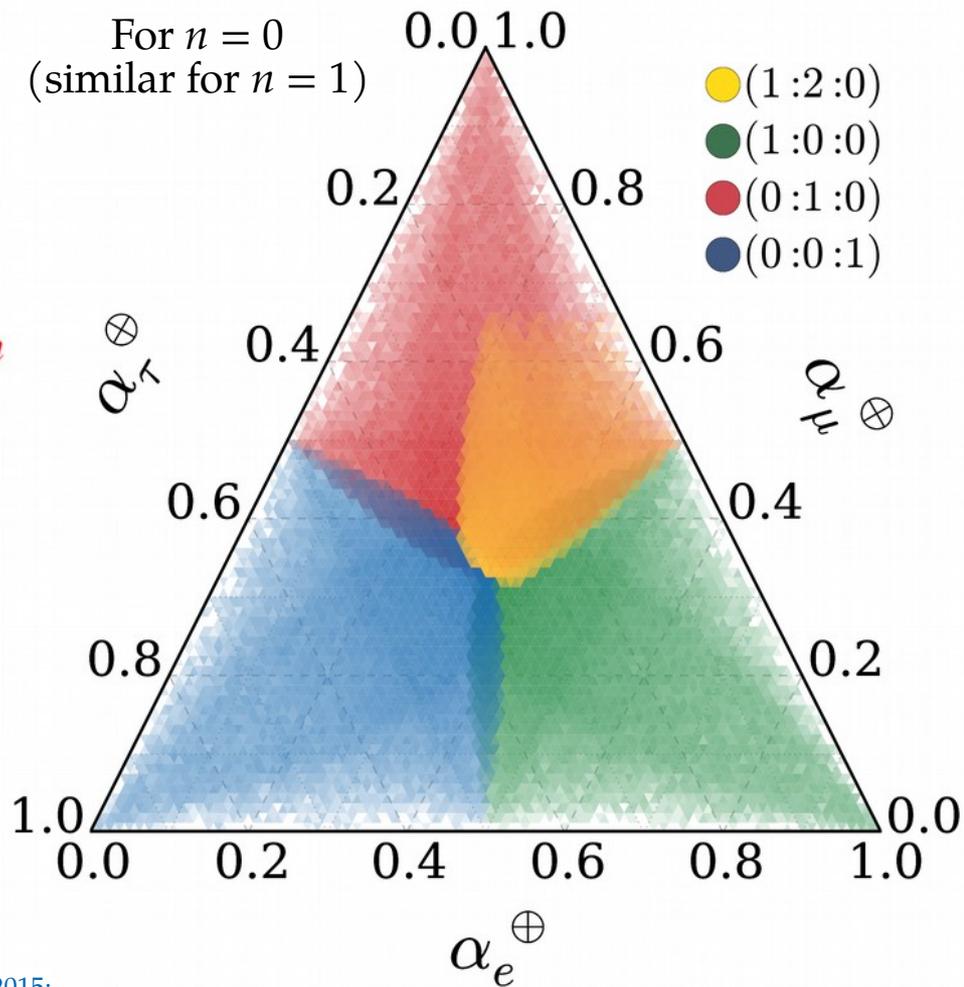
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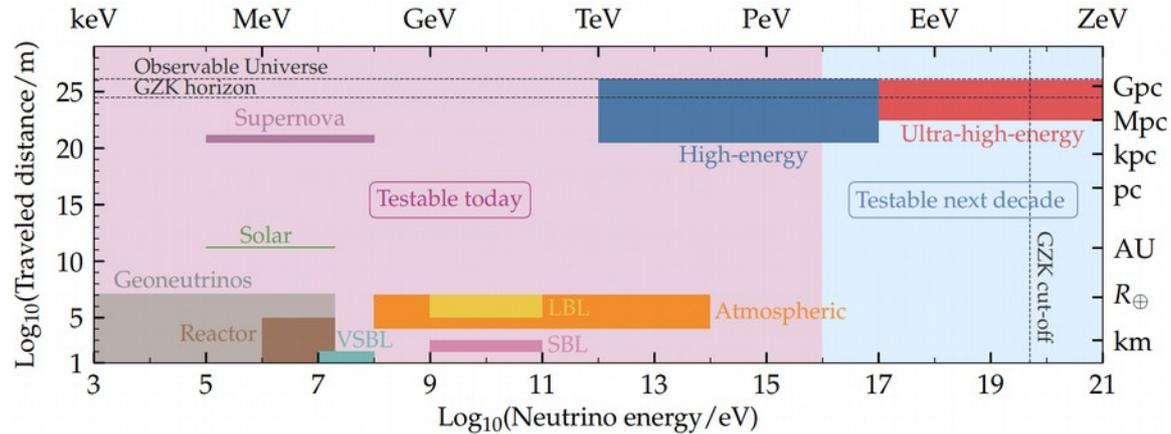
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An exciting decade ahead



Today: TeV–PeV astrophysical ν

$$\kappa_n \sim 4 \cdot 10^{-47} (E/\text{PeV})^{-n} (L/\text{Gpc})^{-1} \text{PeV}^{1-n}$$

Next decade: EeV cosmogenic ν

$$\kappa_n \sim 4 \cdot 10^{-50} (E/\text{EeV})^{-n} (L/\text{Gpc})^{-1} \text{EeV}^{1-n}$$

IceCube + ANTARES + Baikal

+ Growing statistics

+ Improved systematics

IceCube upgrade

IceCube-Gen2

KM3NeT

ANITA

ARA

ARIANNA

Baikal-GVD

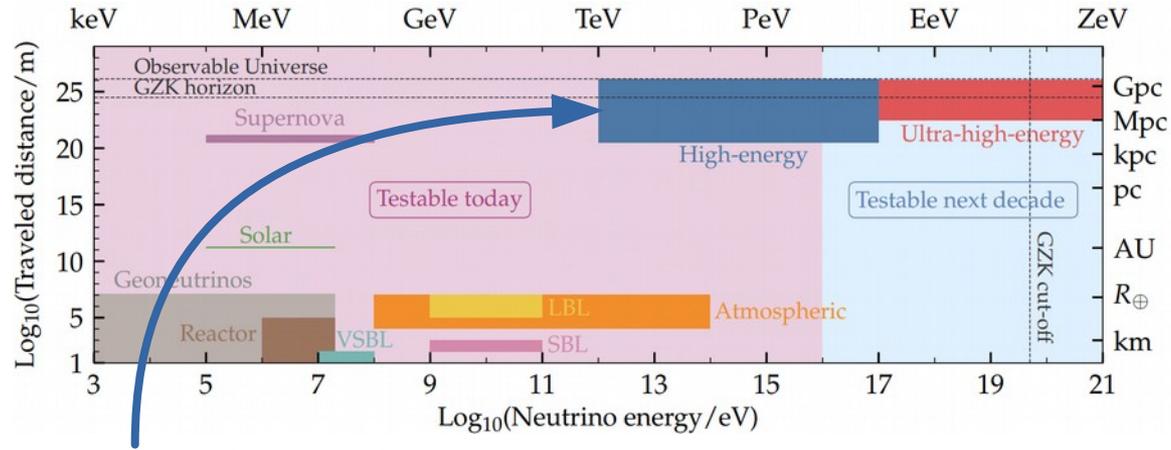
BEACON

GRAND

POEMMA

TRINITY

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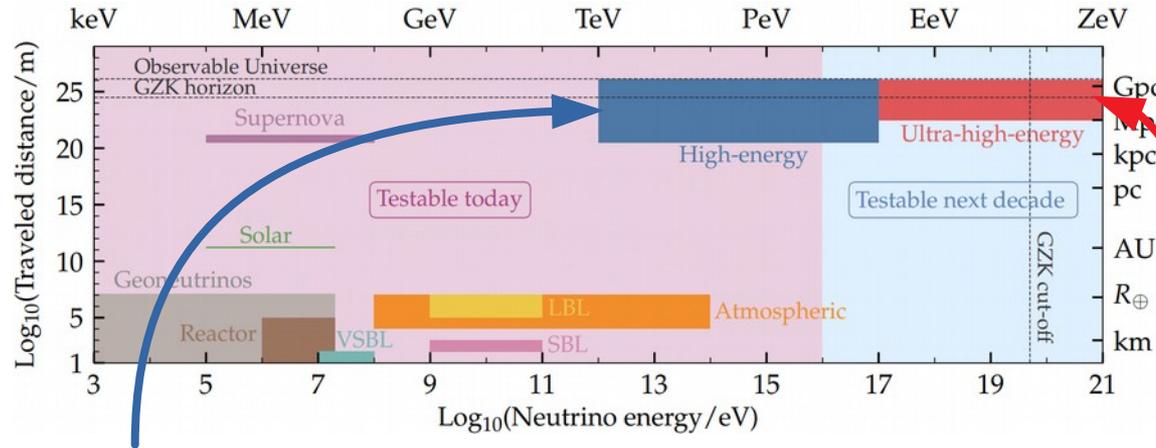
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What are you taking home?

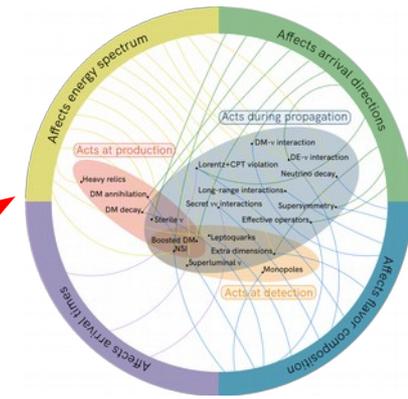
- ▶ Cosmic neutrinos are incisive probes of TeV–PeV physics
- ▶ We can do this *already today*, in spite of unknowns
- ▶ New physics comes in many shapes — so we need to be thorough
- ▶ Exciting prospects: larger statistics, better reconstruction, higher energies

More?

- ▶ *Fundamental physics with high-energy cosmic neutrinos today and in the future*, [1907.08690](#)
- ▶ *Astro2020: Fundamental physics with high-energy cosmic neutrinos*, [1903.04333](#)
- ▶ *Astro2020: Astrophysics uniquely enabled by observations of high-energy cosmic neutrinos*, [1903.04334](#)

What are you taking home?

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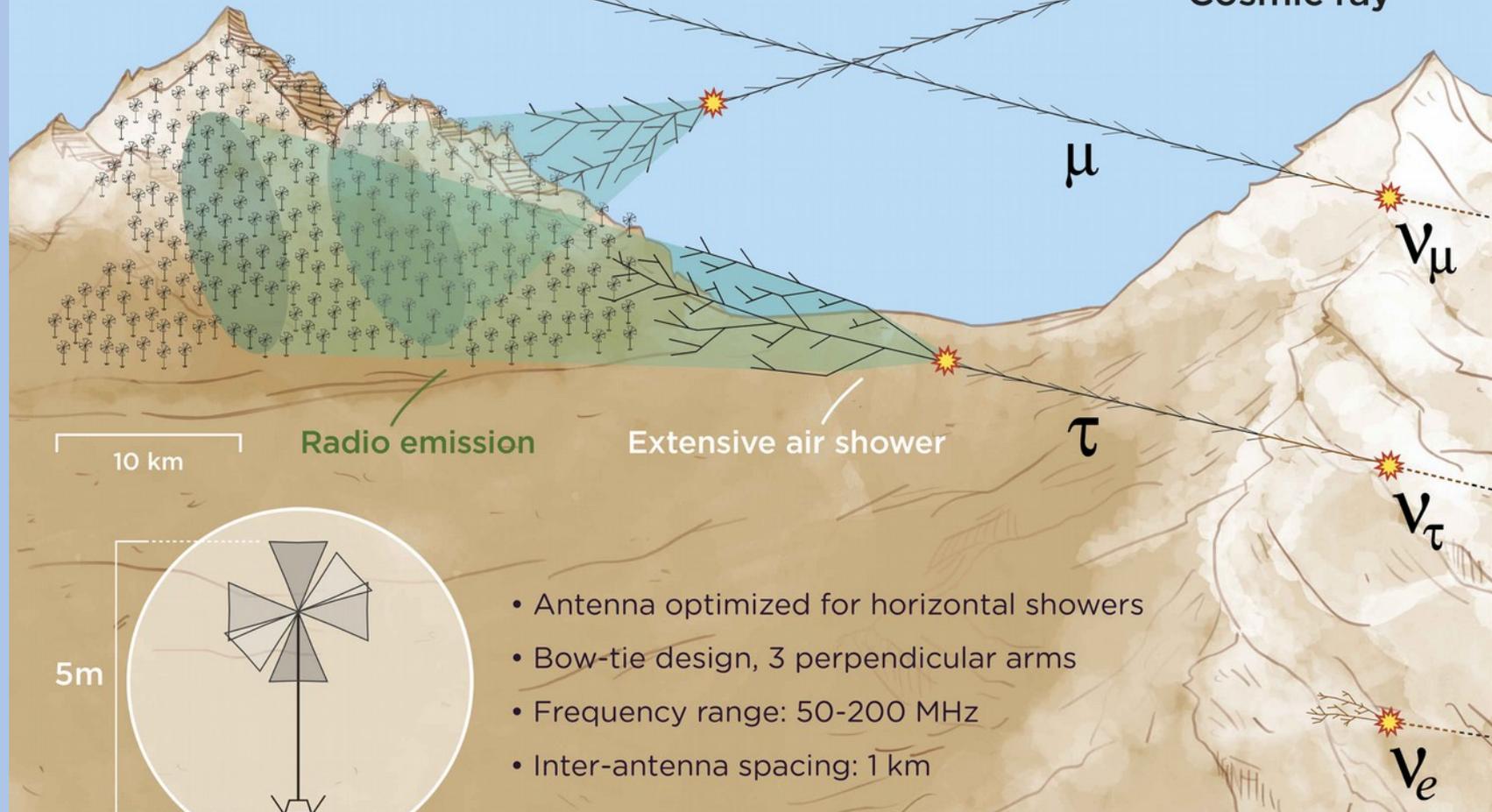
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- ▶ *Astro2020: Fundamental physics with high-energy cosmic neutrinos*, [1903.04333](#)
- ▶ *Astro2020: Astrophysics uniquely enabled by observations of high-energy cosmic neutrinos*, [1903.04334](#)

Backup slides



Giant Radio Array for Neutrino Detection



10 km

Radio emission

Extensive air shower

Cosmic ray

μ

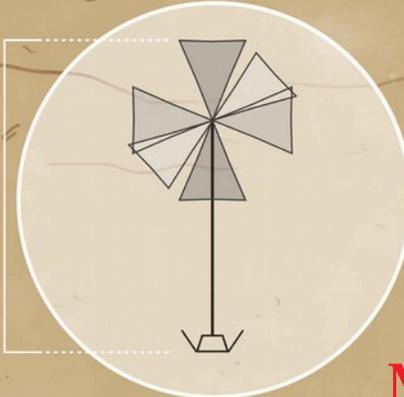
ν_{μ}

τ

ν_{τ}

ν_e

5m



- Antenna optimized for horizontal showers
- Bow-tie design, 3 perpendicular arms
- Frequency range: 50-200 MHz
- Inter-antenna spacing: 1 km

More information about GRAND: grand.cnrs.fr

Flavor-transition probability: the quick and dirty of it

► In matrix form:
$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1}^* & U_{e2}^* & U_{e3}^* \\ U_{\mu 1}^* & U_{\mu 2}^* & U_{\mu 3}^* \\ U_{\tau 1}^* & U_{\tau 2}^* & U_{\tau 3}^* \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

► Pontecorvo-Maki-Nakagawa-Sakata matrix ($c_{ij} = \cos \theta_{ij}$, $s_{ij} = \sin \theta_{ij}$):

$$U = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\text{Atmospheric}} \underbrace{\begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix}}_{\text{Cross mixing}} \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{Solar}} \underbrace{\begin{pmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{Majorana CP phases}}$$

► Probability for $\nu_\alpha \rightarrow \nu_\beta$:
$$P_{\nu_\alpha \rightarrow \nu_\beta} = \delta_{\alpha\beta} - 4 \sum_{i>j} \text{Re}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin^2 \left(\Delta m_{ij}^2 \frac{L}{4E} \right) + 2 \sum_{i>j} \text{Im}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin \left(\Delta m_{ij}^2 \frac{L}{2E} \right)$$

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$\theta_{23} \approx 48^\circ$
 $\theta_{13} \approx 9^\circ$
 $\theta_{12} \approx 34^\circ$
 $\delta \approx 222^\circ$

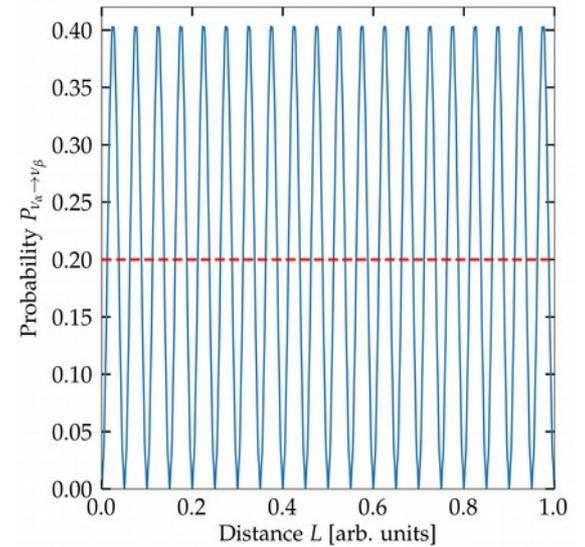
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... But high-energy neutrinos oscillate *fast*

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Oscillation length for 1-TeV ν : $2\pi \times 2E/\Delta m^2 \sim 0.1 \text{ pc}$

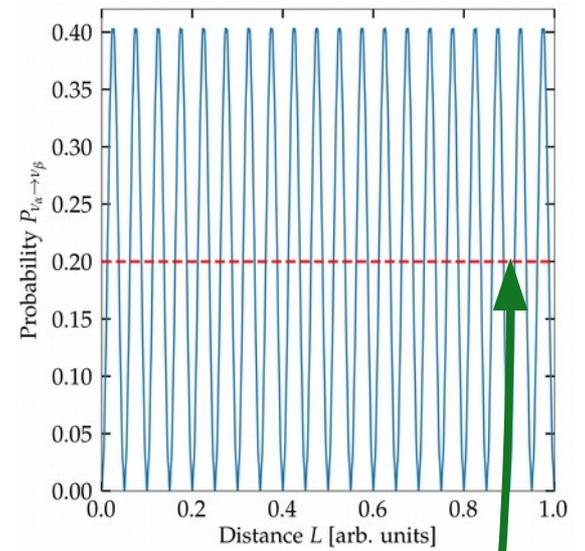
- $\sim 8\%$ of the way to Proxima Centauri
- \ll Distance to Galactic Center (8 kpc)
- \ll Distance to Andromeda (1 Mpc)
- \ll Cosmological distances (few Gpc)

We cannot resolve oscillations, so we use instead the average probability:

$$\langle P_{\nu_\alpha \rightarrow \nu_\beta} \rangle = \sum_{i=1}^3 |U_{\alpha i}|^2 |U_{\beta i}|^2$$

... But high-energy neutrinos oscillate *fast*

$$P_{\nu_\alpha \rightarrow \nu_\beta} = \delta_{\alpha\beta} - 4 \sum_{i>j} \text{Re}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin^2 \left(\Delta m_{ij}^2 \frac{L}{4E} \right) + 2 \sum_{i>j} \text{Im}(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin \left(\Delta m_{ij}^2 \frac{L}{2E} \right)$$



Oscillation length for 1-TeV ν : $2\pi \times 2E/\Delta m^2 \sim 0.1$ pc

$\sim 8\%$ of the way to Proxima Centauri
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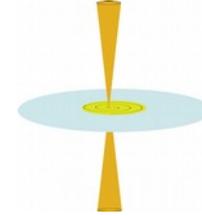
$$\langle P_{\nu_\alpha \rightarrow \nu_\beta} \rangle = \sum_{i=1}^3 |U_{\alpha i}|^2 |U_{\beta i}|^2$$

Inferring the flavor composition at the sources

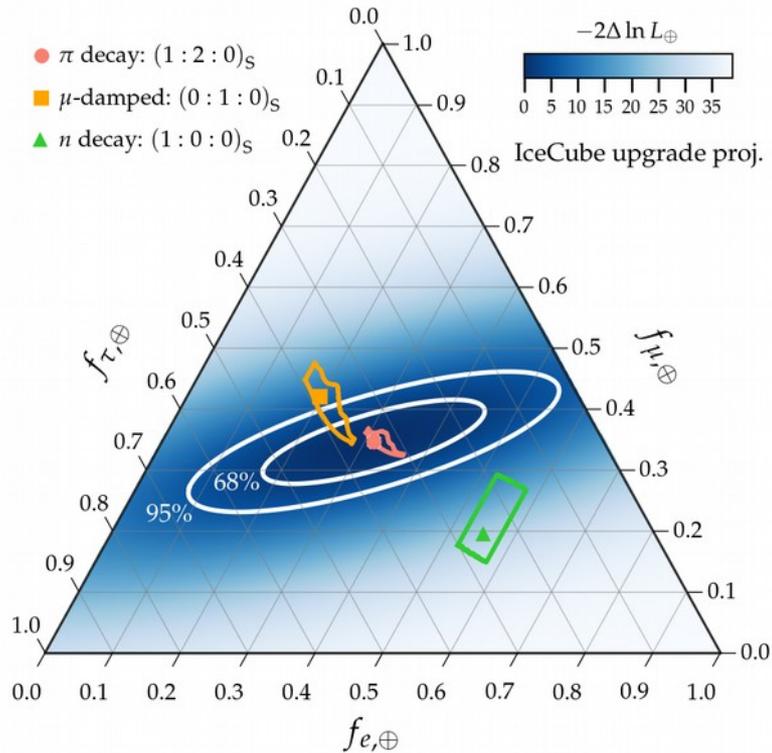
Measured:
Flavor ratios at Earth



Invert flavor oscillations



Inferred:
Flavor ratios at
astrophysical sources

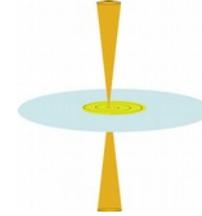


Inferring the flavor composition at the sources

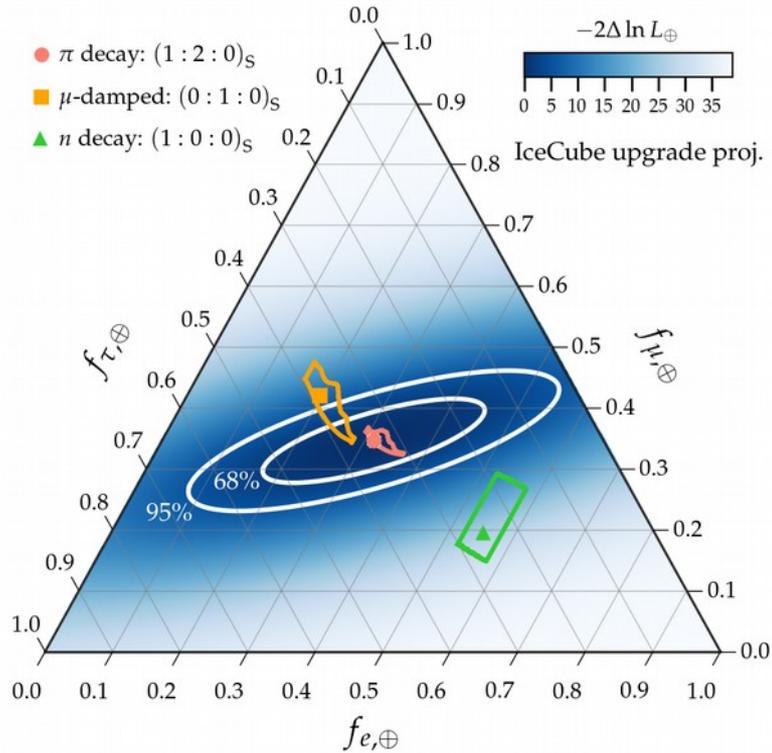
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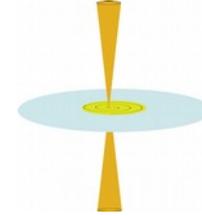


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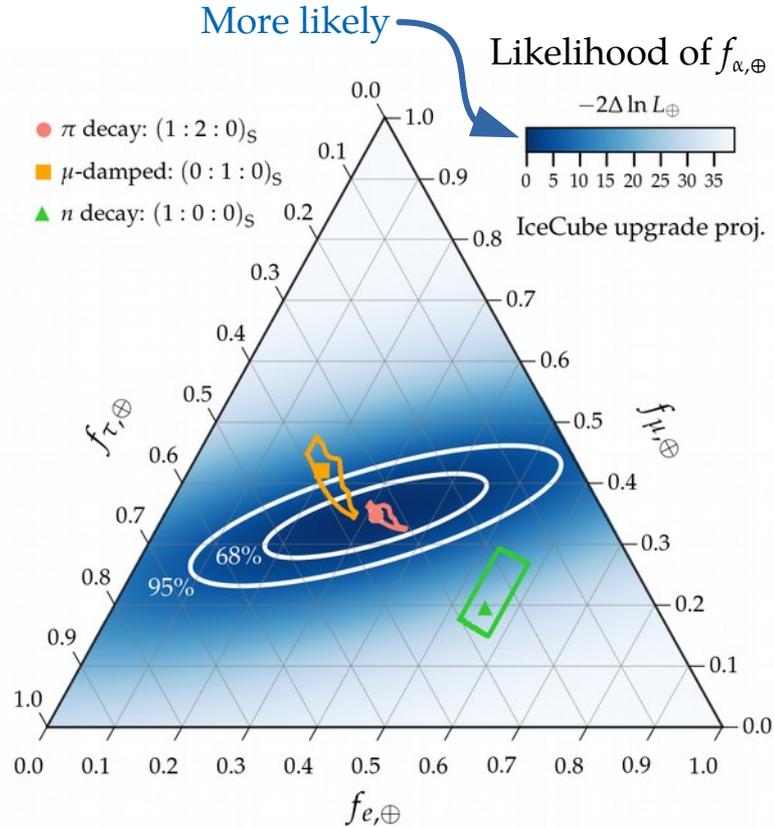
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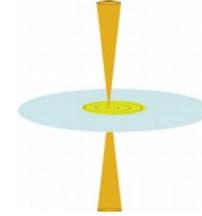


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Flavor ratios at Earth



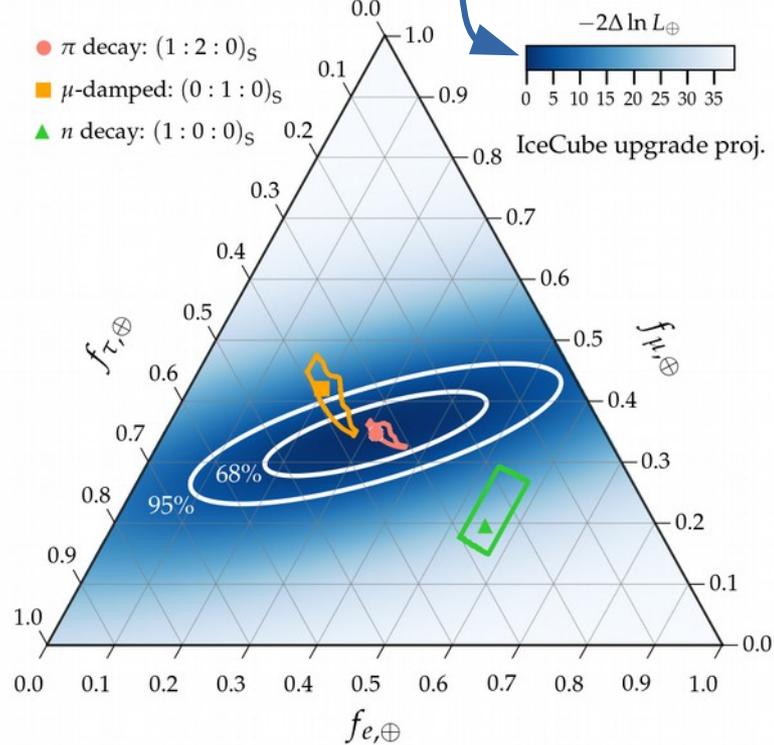
Invert flavor oscillations



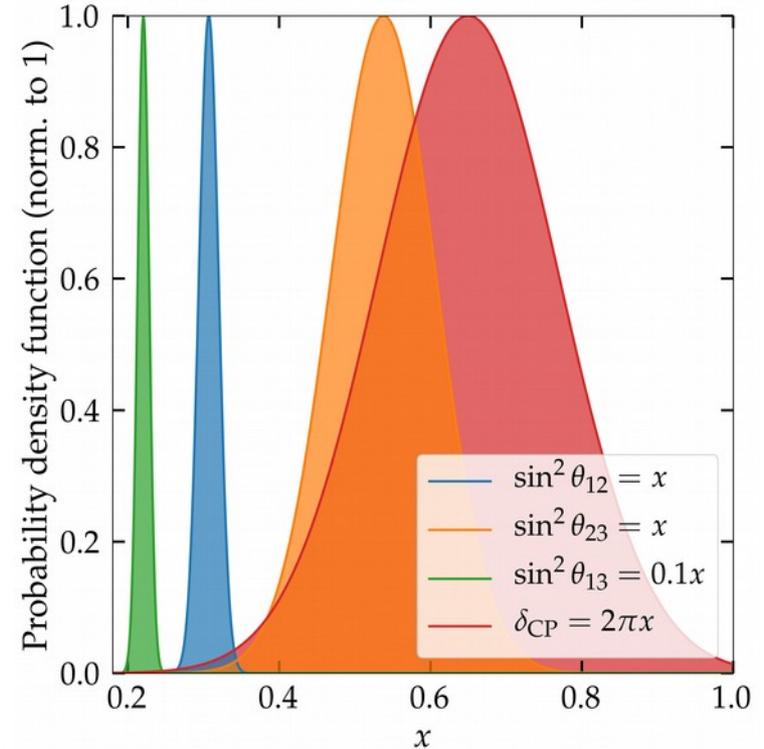
Inferred:
Flavor ratios at
astrophysical sources

More likely

Likelihood of $f_{\alpha,\oplus}$



PDFs of mixing parameters

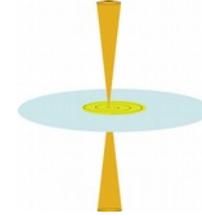


Inferring the flavor composition at the sources

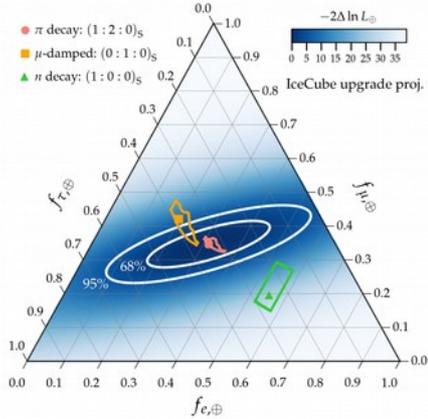
Measured:
Flavor ratios at Earth



Invert flavor oscillations



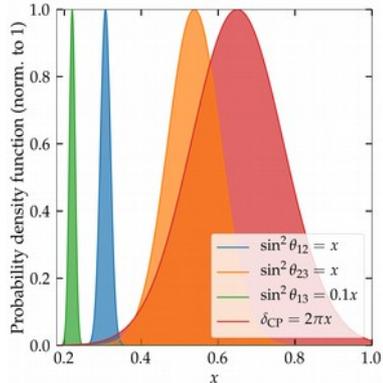
Inferred:
Flavor ratios at
astrophysical sources



Posterior probability density of $f_{\alpha,S}$ being the flavor ratios at the sources:

$$\mathcal{P}(f_{\alpha,S}) \equiv \int d\boldsymbol{\theta} \frac{\mathcal{P}(\boldsymbol{\theta})}{\mathcal{N}(\boldsymbol{\theta})} \mathcal{L}_{\oplus} [f_{e,\oplus}(f_{\alpha,S}, \boldsymbol{\theta}), f_{\mu,\oplus}(f_{\alpha,S}, \boldsymbol{\theta})]$$

$$\boldsymbol{\theta} \equiv (\theta_{12}, \theta_{23}, \theta_{13}, \delta_{CP})$$



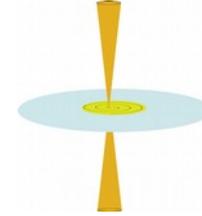
$$\left[\text{Normalization: } \mathcal{N}(\boldsymbol{\theta}) \equiv \int_0^1 df_{e,S} \int_0^{1-f_{e,S}} df_{\mu,S} \mathcal{L}_{\oplus} [f_{e,\oplus}(f_{\alpha,S}, \boldsymbol{\theta}), f_{\mu,\oplus}(f_{\alpha,S}, \boldsymbol{\theta})] \right]$$

Inferring the flavor composition at the sources

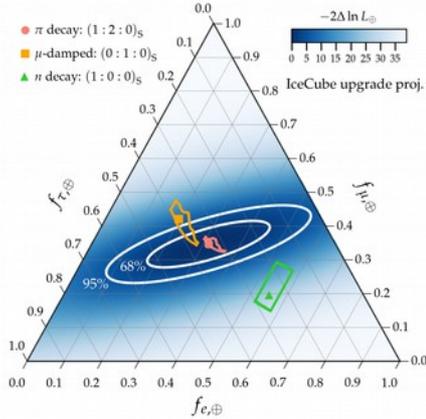
Measured:
Flavor ratios at Earth



Invert flavor oscillations



Inferred:
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astrophysical sources

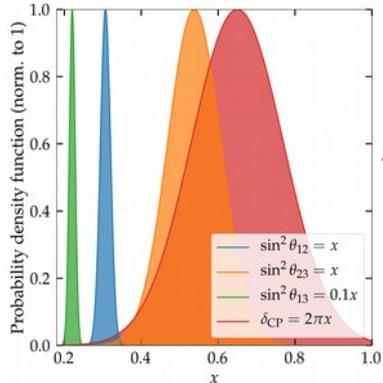


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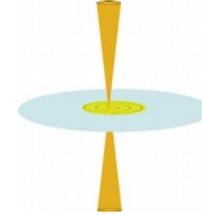


Inferring the flavor composition at the sources

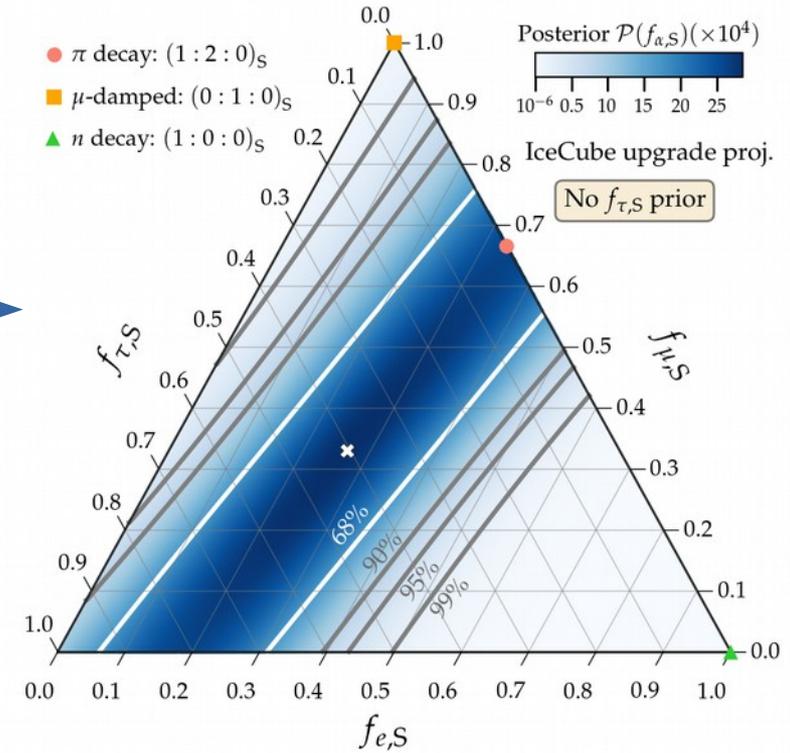
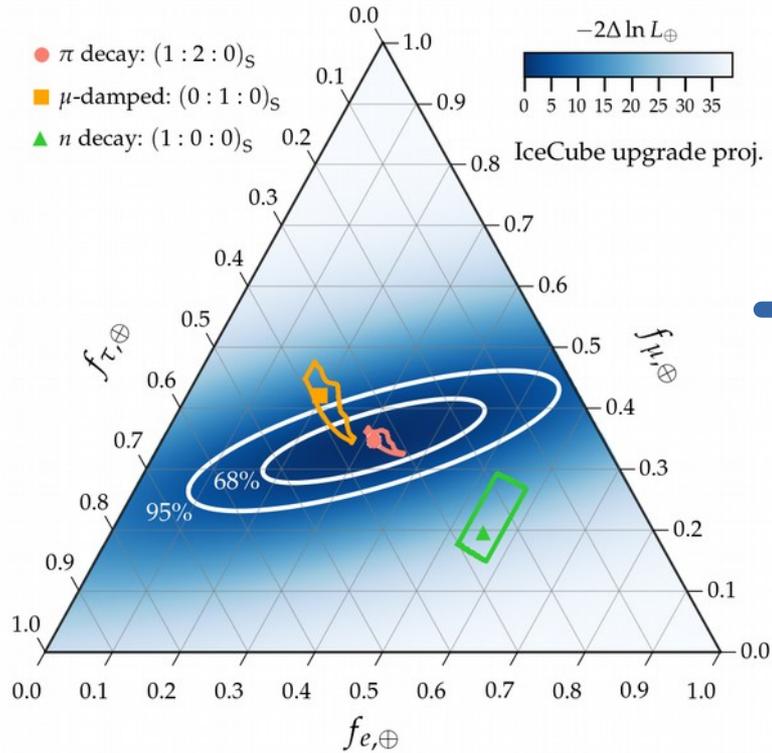
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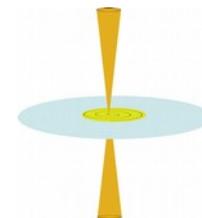


Inferring the flavor composition at the sources

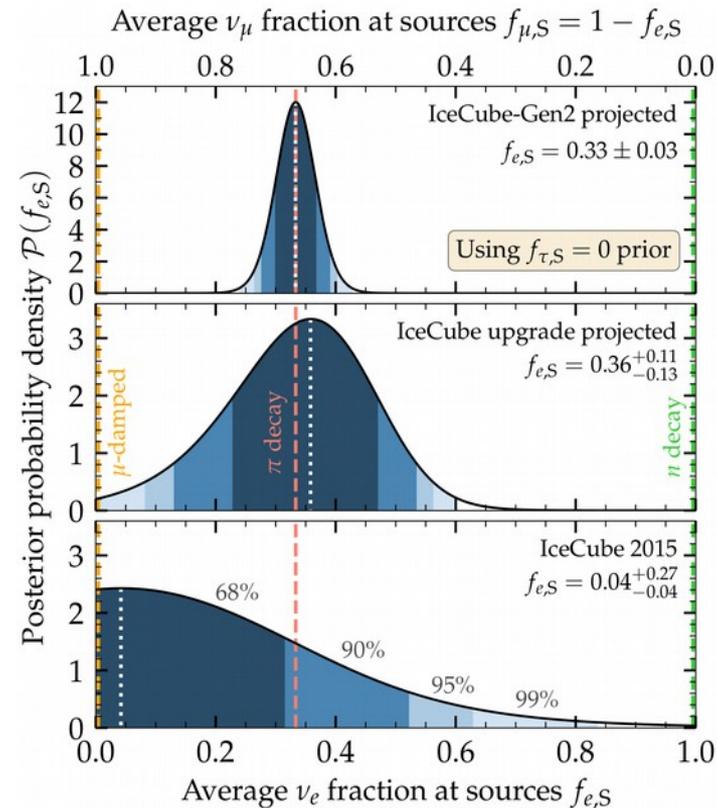
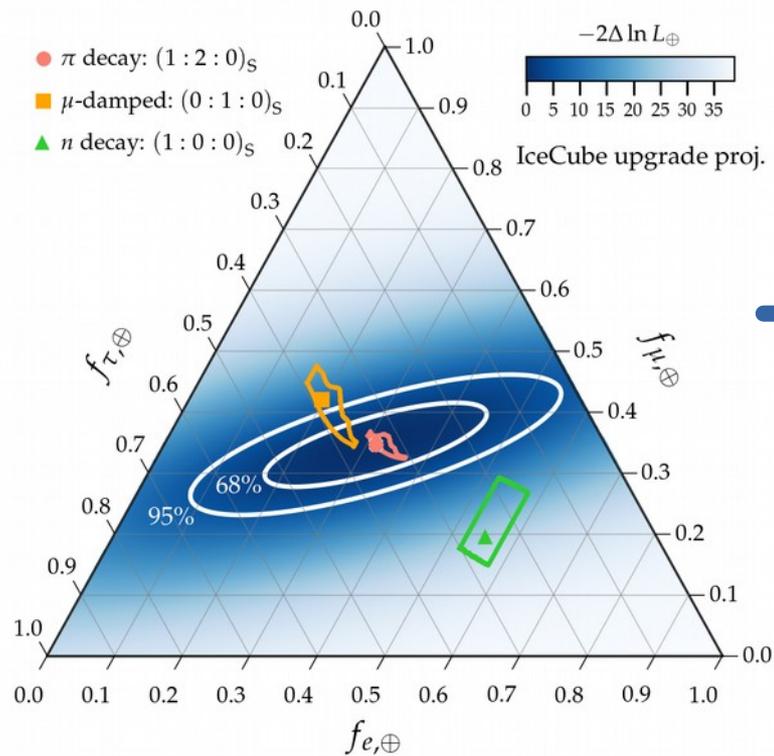
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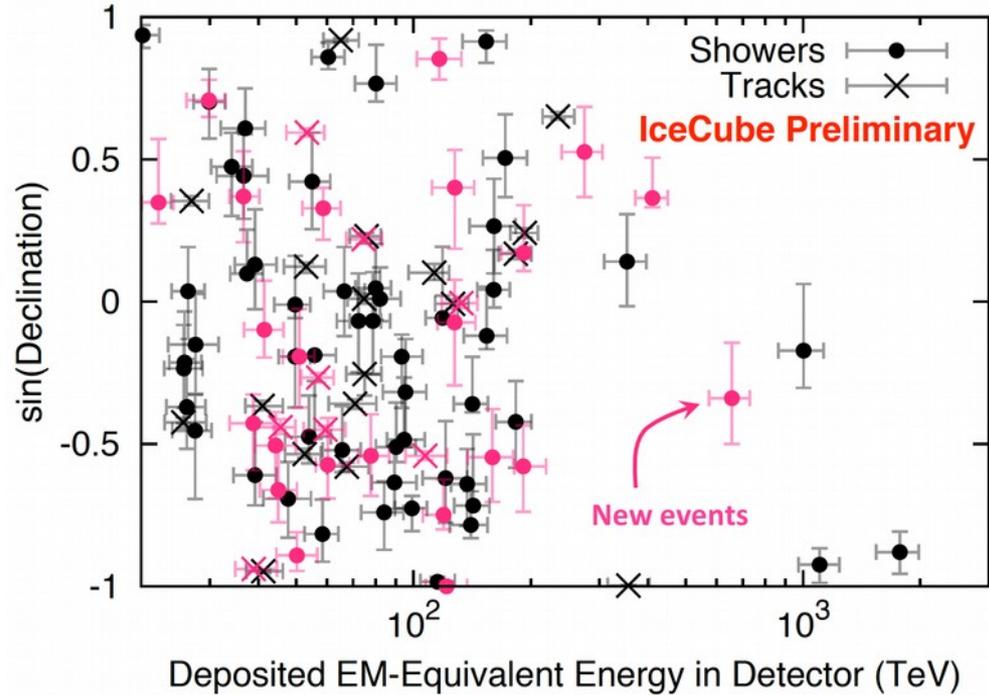


What has IceCube found so far (7.5 years)?

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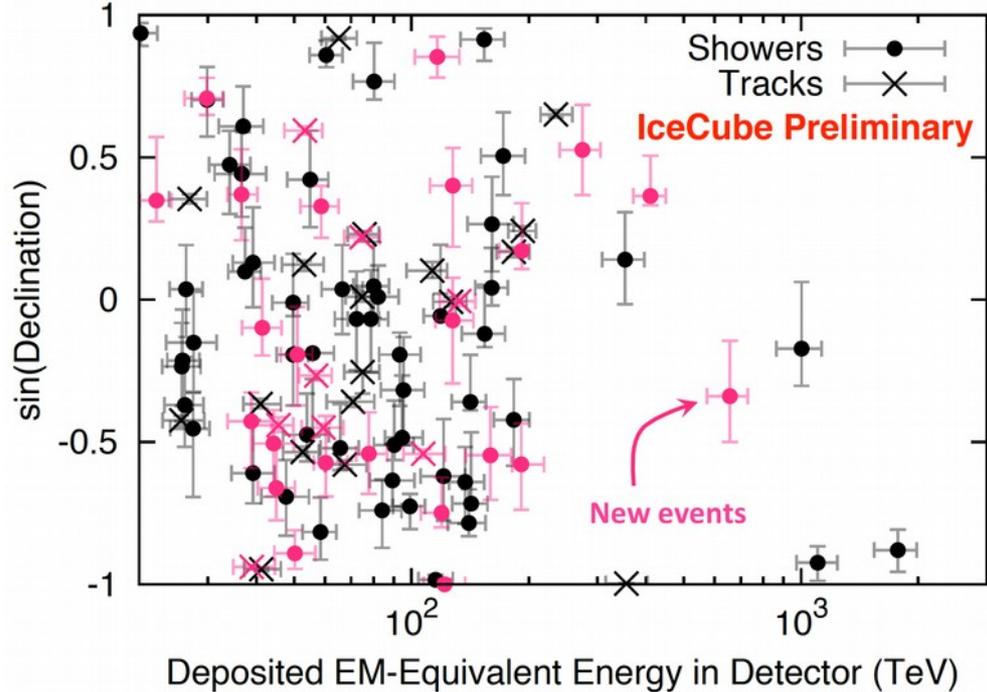
What has IceCube found so far (7.5 years)?

103 contained events between 15 TeV – 2 PeV

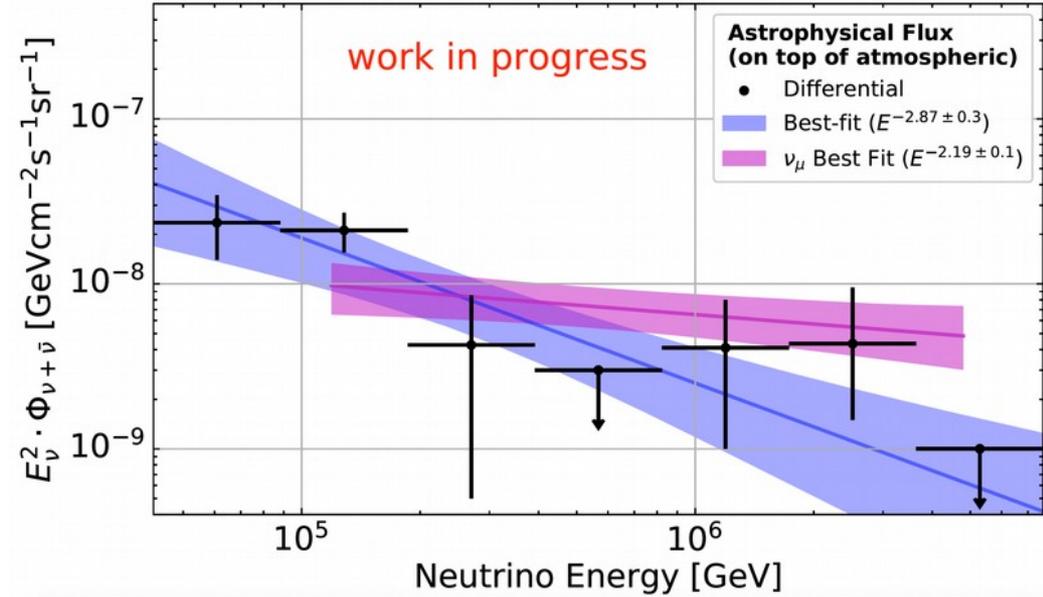


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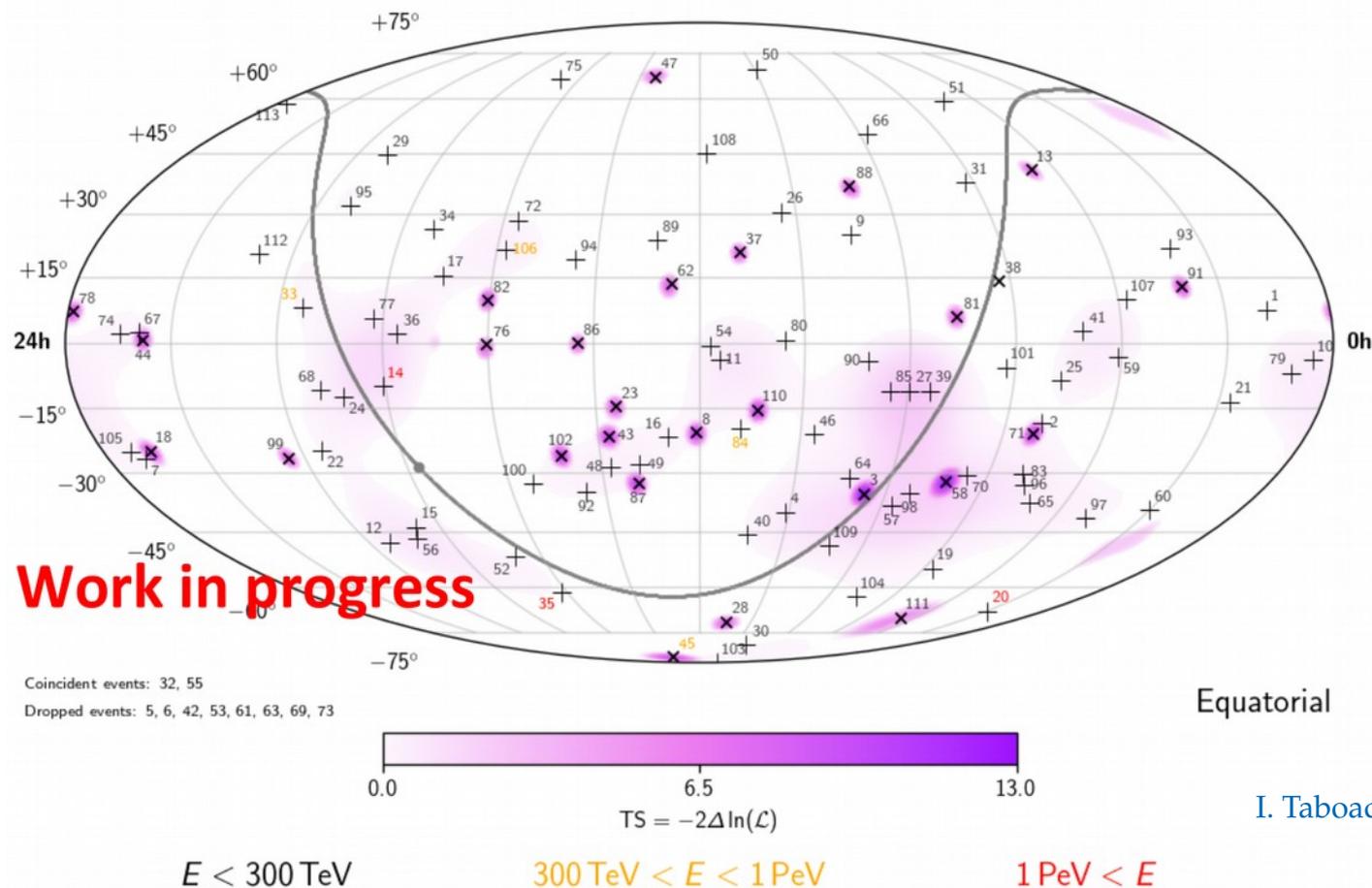


Astrophysical ν flux detected at $> 7\sigma$
(Normalization ok, but steep spectrum)



What has IceCube found so far (7.5 years)?

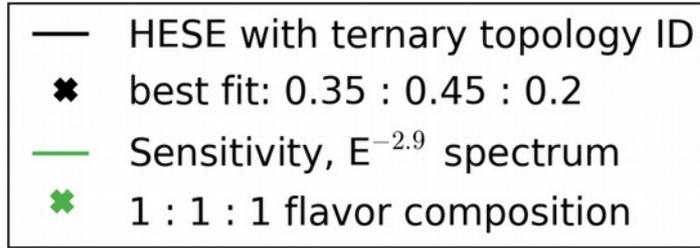
Arrival directions compatible with isotropy



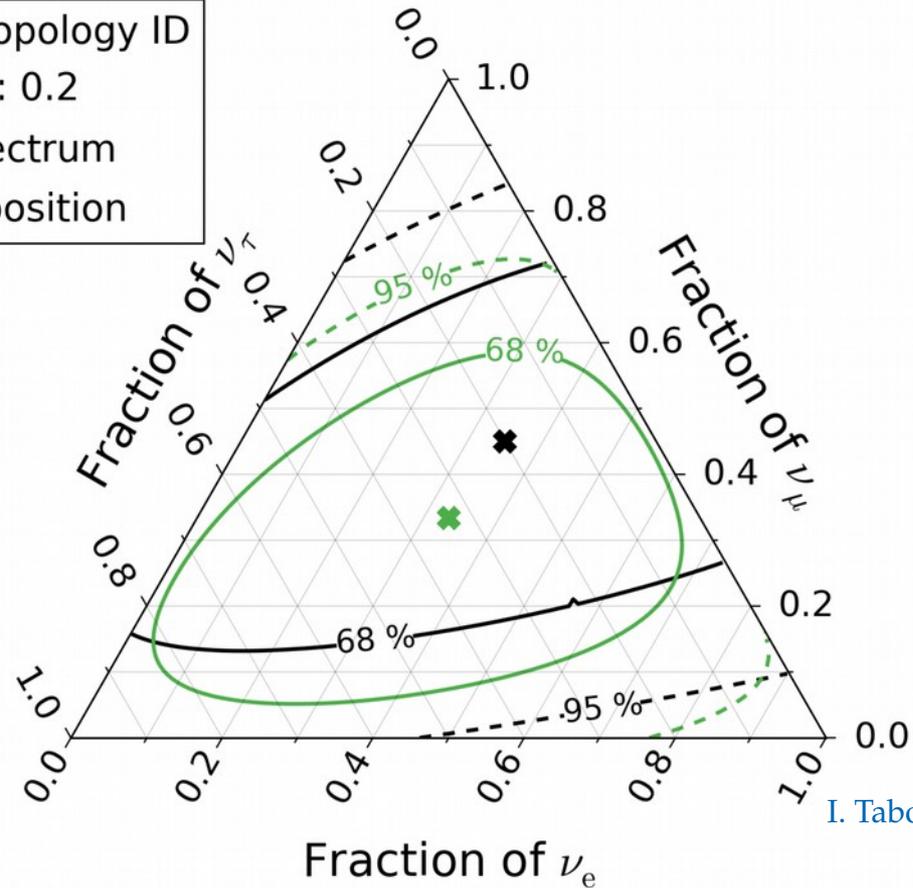
I. Taboada, Neutrino 2018

What has IceCube found so far (7.5 years)?

Flavor composition compatible with equal proportion of each flavor

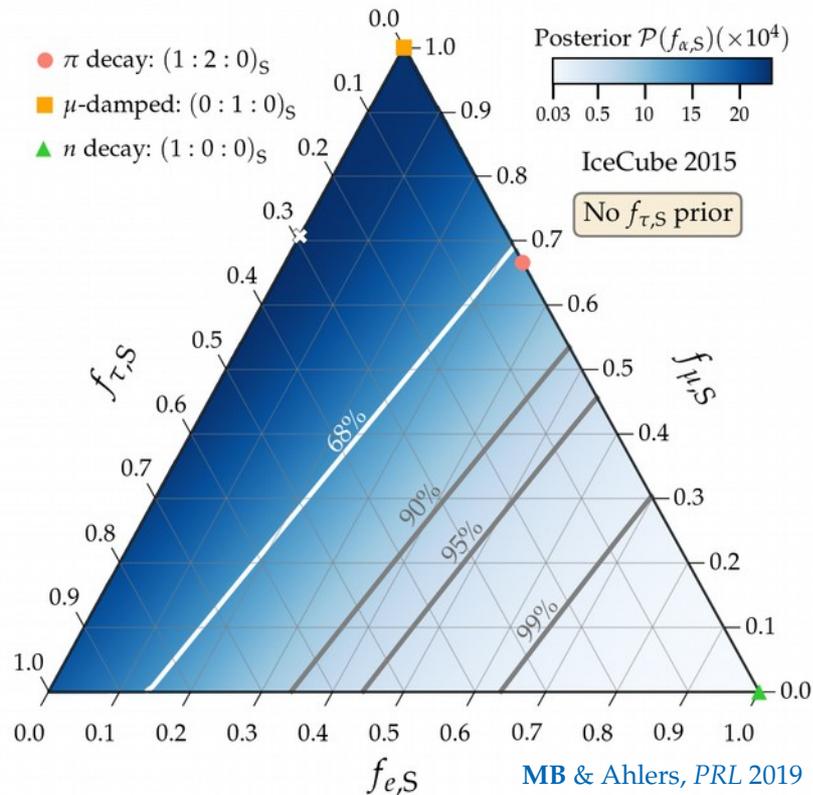


WORK IN PROGRESS

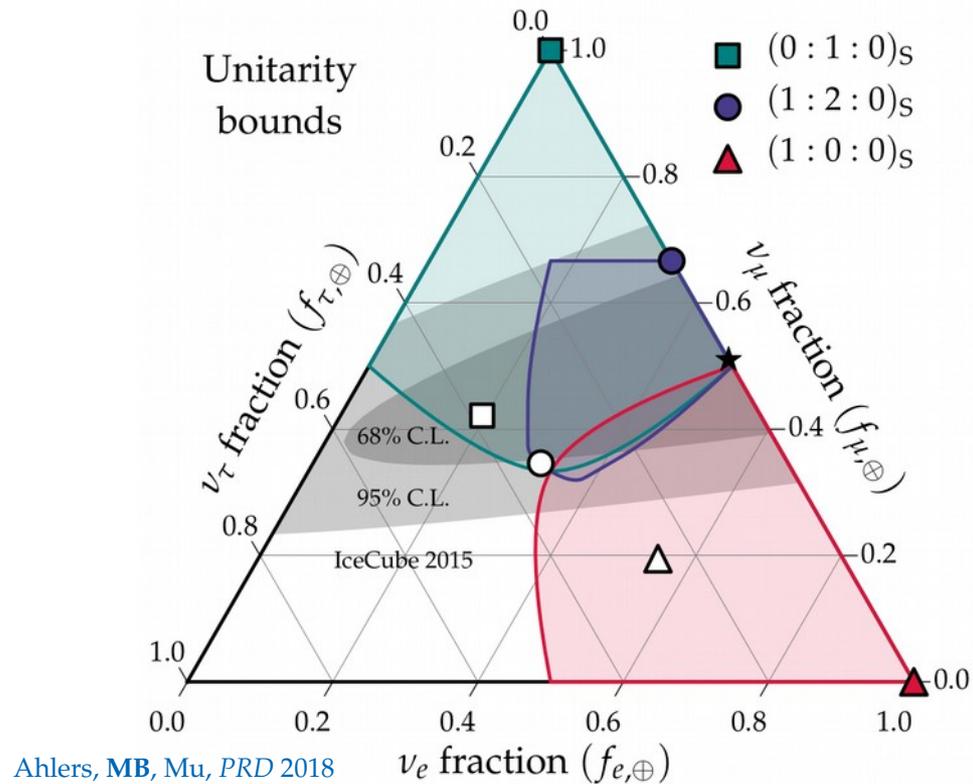


Flavor – What is it good for?

Trusting **particle physics**
and learning about **astrophysics**



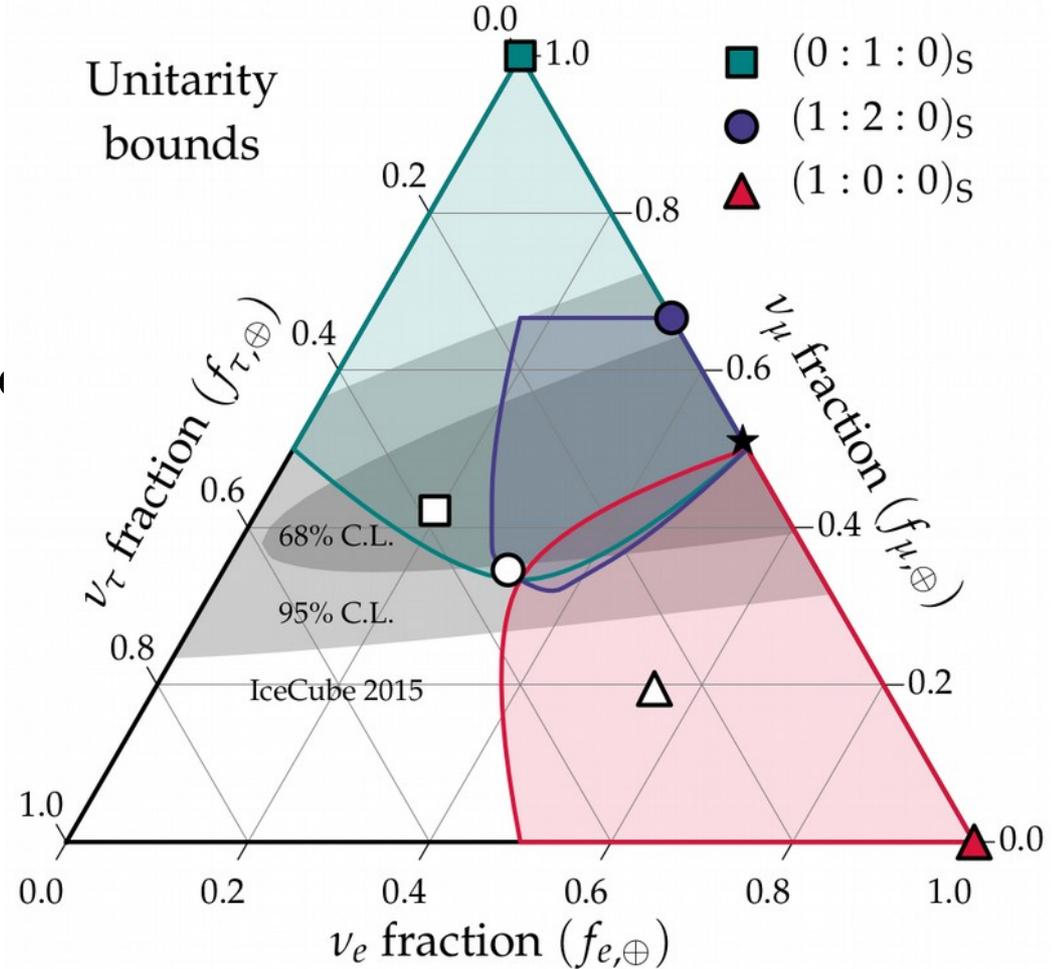
Trusting **astrophysics**
and learning about **particle physics**



Using unitarity to constrain new physics

$$H_{\text{tot}} = H_{\text{std}} + H_{\text{NP}}$$


- ▶ New mixing angles unconstrained
- ▶ Use unitarity ($U_{\text{NP}}U_{\text{NP}}^\dagger = 1$) to bound all possible flavor ratios at Earth
- ▶ Can be used as prior in new-physics searches in IceCube



What lies beyond? *Take your pick*

- ▶ High-energy effective field theories
 - ▶ Violation of Lorentz and CPT invariance
[Barenboim & Quigg, *PRD* 2003; MB, Gago, Peña-Garay, *JHEP* 2010; Kostelecky & Mewes 2004]
 - ▶ Violation of equivalence principle
[Gasperini, *PRD* 1989; Glashow *et al.*, *PRD* 1997]
 - ▶ Coupling to a gravitational torsion field
[De Sabbata & Gasperini, *Nuovo Cim.* 1981]
 - ▶ Renormalization-group-running of mixing parameters
[MB, Gago, Jones, *JHEP* 2011]
 - ▶ General non-unitary propagation
[Ahlers, MB, Mu, *PRD* 2018]
- ▶ Active-sterile mixing
[Aeikens *et al.*, *JCAP* 2015; Brdar, *JCAP* 2017]
- ▶ Flavor-violating physics
 - ▶ New neutrino-electron interactions
[MB & Agarwalla, *PRL* 2019]
 - ▶ New $\nu\nu$ interactions
[Ng & Beacom, *PRD* 2014; Cherry, Friedland, Shoemaker, 1411.1071; Blum, Hook, Murase, 1408.3799]
- ▶ ...



Toho Company Ltd.

Ultra-long-range flavorful interactions

- ▶ **Simple extension of the SM:** Promote the global lepton-number symmetries L_e-L_μ , L_e-L_τ to local symmetries
- ▶ They introduce new interaction between electrons and ν_e and ν_μ or ν_τ mediated by a new neutral vector boson (Z'):
 - ▶ Affects oscillations
 - ▶ If the Z' is *very* light, *many* electrons can contribute

The new potential sourced by an electron

Under the L_e-L_μ or L_e-L_τ symmetry, an electron sources a Yukawa potential —

$$V \sim \frac{g'_{e\beta}{}^2}{r} e^{-m'_{e\beta} r}$$

A neutrino “feels” all the electrons within the interaction range $\sim(1/m')$

The new potential sourced by an electron

Under the L_e-L_μ or L_e-L_τ symmetry, an electron sources a Yukawa potential —

$$V \sim \frac{g'_{e\beta}{}^2}{r} e^{-m'_{e\beta} r}$$

Z' coupling \rightarrow $g'_{e\beta}{}^2$ \leftarrow *Z'* mass $m'_{e\beta}$

r \leftarrow Distance to neutrino

A neutrino “feels” all the electrons within the interaction range $\sim (1/m')$

Electron-neutrino interactions can kill oscillations

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$$H_{\text{tot}} = H_{\text{vac}}$$


Standard oscillations:
Neutrinos change flavor
because this is non-diagonal

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$$H_{\text{tot}} = H_{\text{vac}}$$


Standard oscillations:
Neutrinos change flavor
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$$P_{\nu_\alpha \rightarrow \nu_\beta}(\theta_{ij}, \delta_{\text{CP}})$$

Electron-neutrino interactions can kill oscillations

$$H_{\text{tot}} = H_{\text{vac}} + \underbrace{V_{e\beta}}_{= \text{diag}(V_{e\mu}, -V_{e\mu}, 0)}$$

New neutrino-electron interaction:
This is diagonal

Electron-neutrino interactions can kill oscillations

$$H_{\text{tot}} = H_{\text{vac}} + \underbrace{V_{e\beta}}_{= \text{diag}(V_{e\mu}, -V_{e\mu}, 0)}$$

New neutrino-electron interaction:
This is diagonal

↓

$$P_{\nu_\alpha \rightarrow \nu_\beta} \left(\theta_{ij}, \delta_{\text{CP}}, \Delta m_{ij}^2, E_\nu, \overbrace{g'_{e\mu}, m'_{e\mu}}^{\text{Z' parameters}} \right)$$

Electron-neutrino interactions can kill oscillations

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New neutrino-electron interaction:
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↓

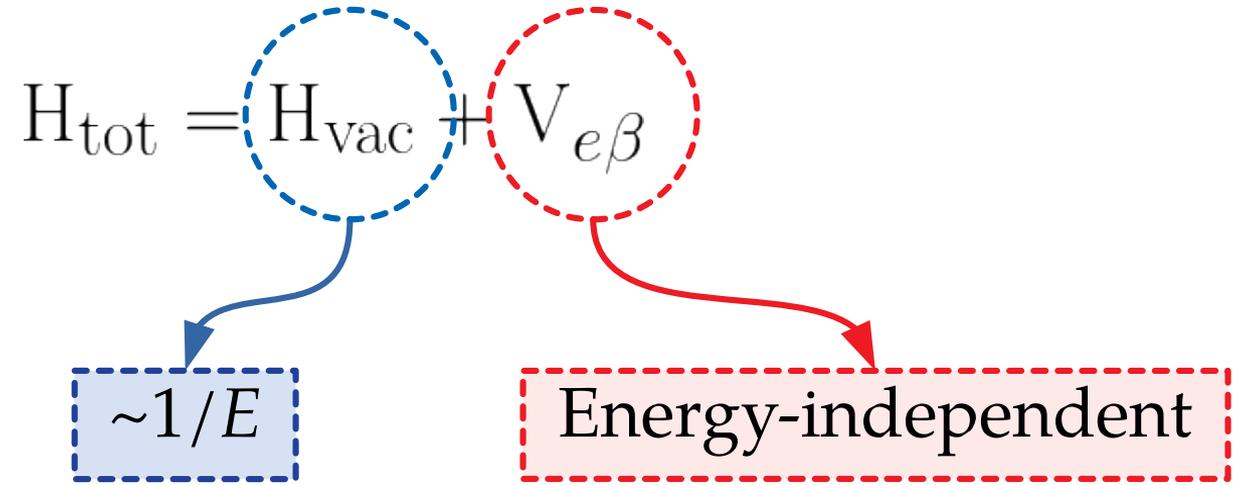
$$P_{\nu_\alpha \rightarrow \nu_\beta} \left(\theta_{ij}, \delta_{\text{CP}}, \Delta m_{ij}^2, E_\nu, \overbrace{g'_{e\mu}, m'_{e\mu}}^{\text{Z' parameters}} \right)$$

If $V_{e\beta}$ dominates ($g' \gg 1, m' \ll 1$), oscillations turn off

Electron-neutrino interactions can kill oscillations

$$H_{\text{tot}} = H_{\text{vac}} + V_{e\beta}$$

Electron-neutrino interactions can kill oscillations



Electron-neutrino interactions can kill oscillations

$$H_{\text{tot}} = H_{\text{vac}} + V_{e\beta}$$

$\sim 1/E$

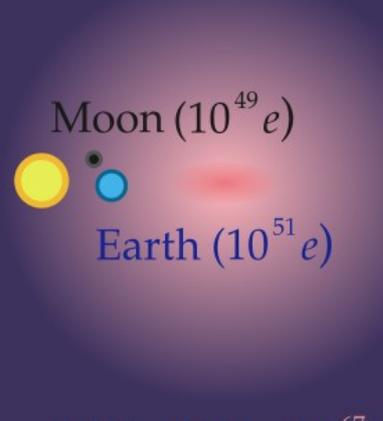
Energy-independent

\therefore We can use high-energy astrophysical neutrinos

The total potential

Cosmological electrons ($10^{79} e$)

Sun ($10^{57} e$)



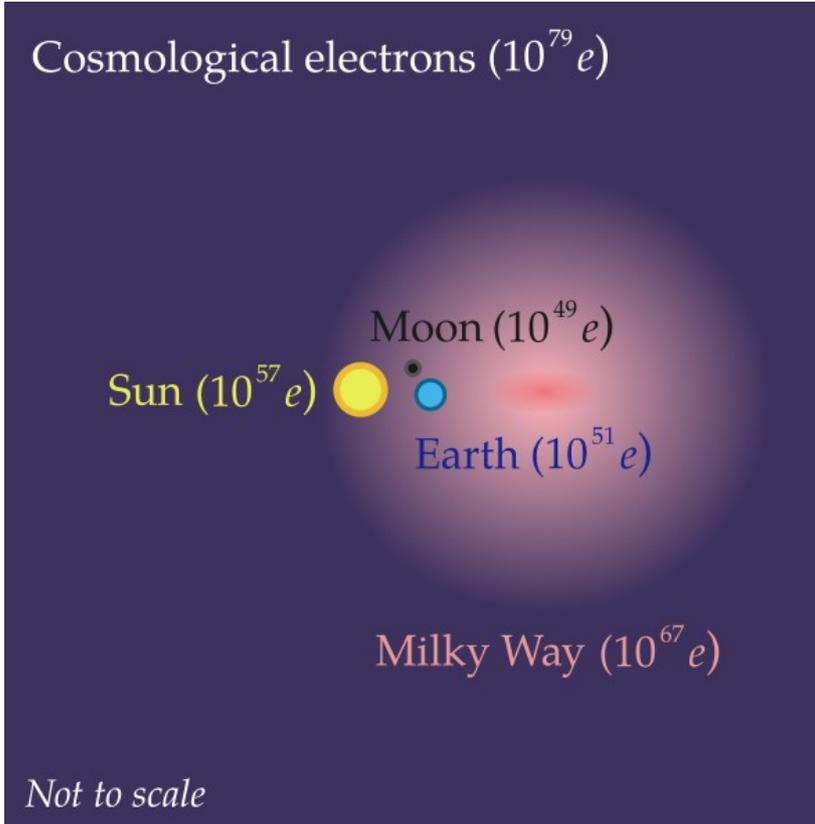
Moon ($10^{49} e$)

Earth ($10^{51} e$)

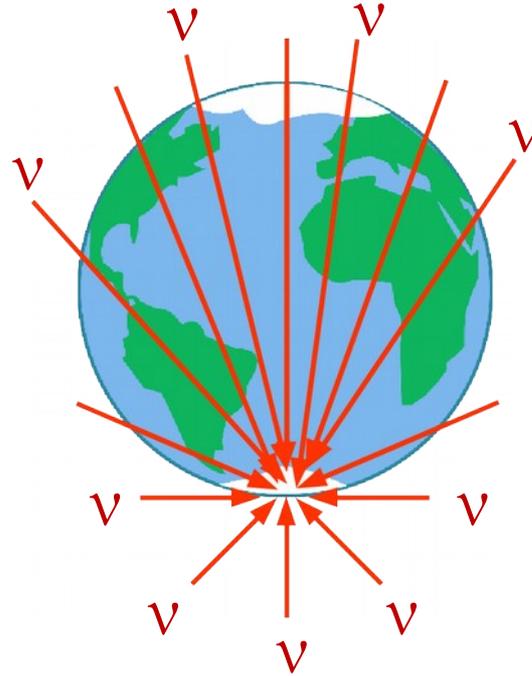
Milky Way ($10^{67} e$)

Not to scale

The total potential

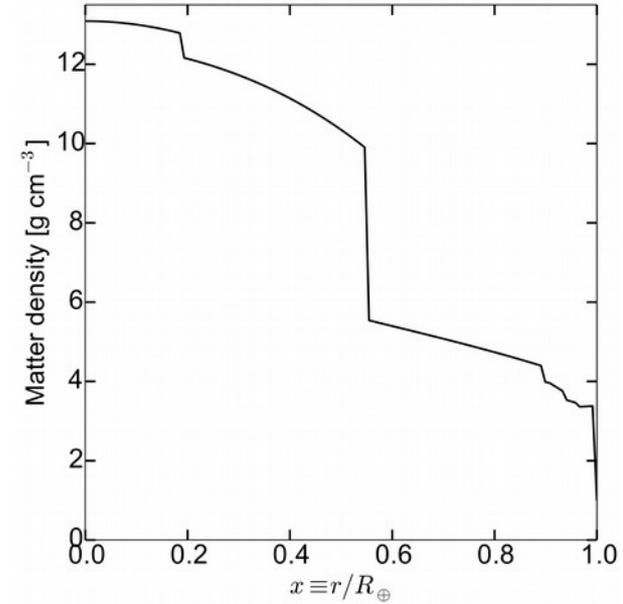


Earth:



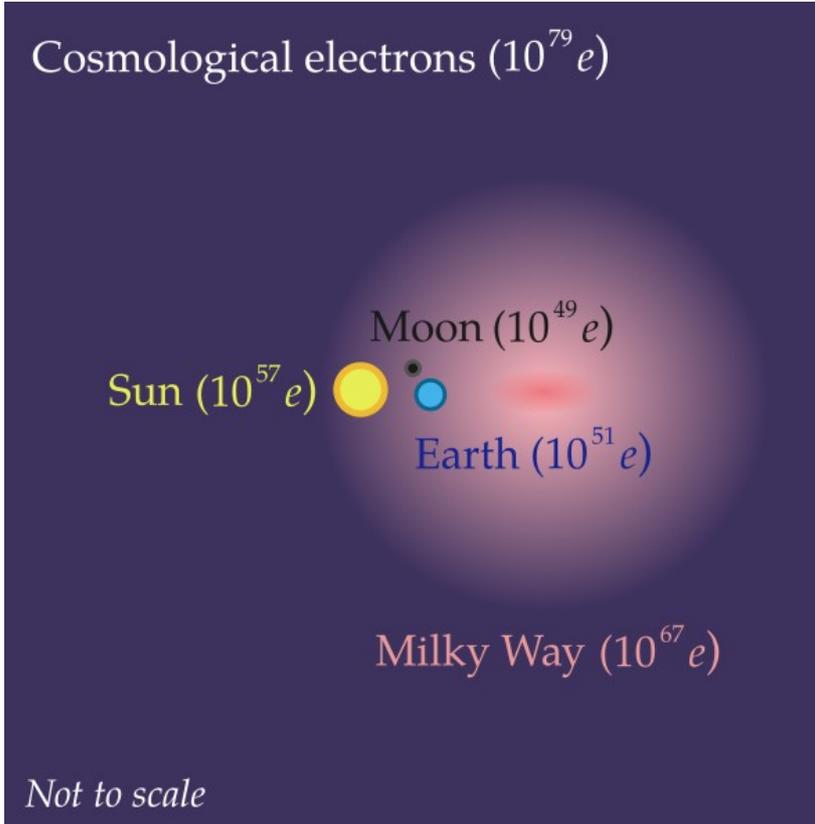
Neutrinos traverse different electron column depths

Preliminary Reference Earth Model
Dziewonski & Anderson 1981

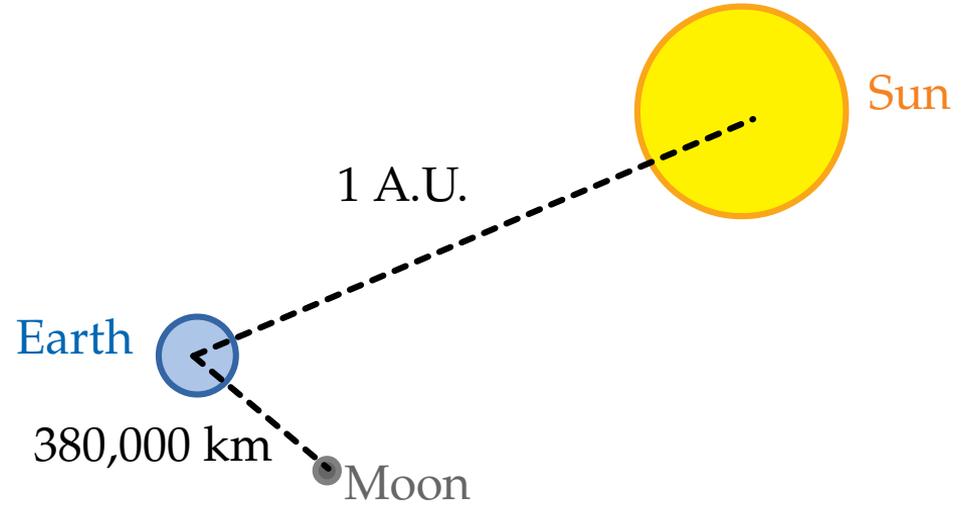


$$V_{e\beta} = V_{e\beta}^{\oplus}$$

The total potential



Moon and Sun:



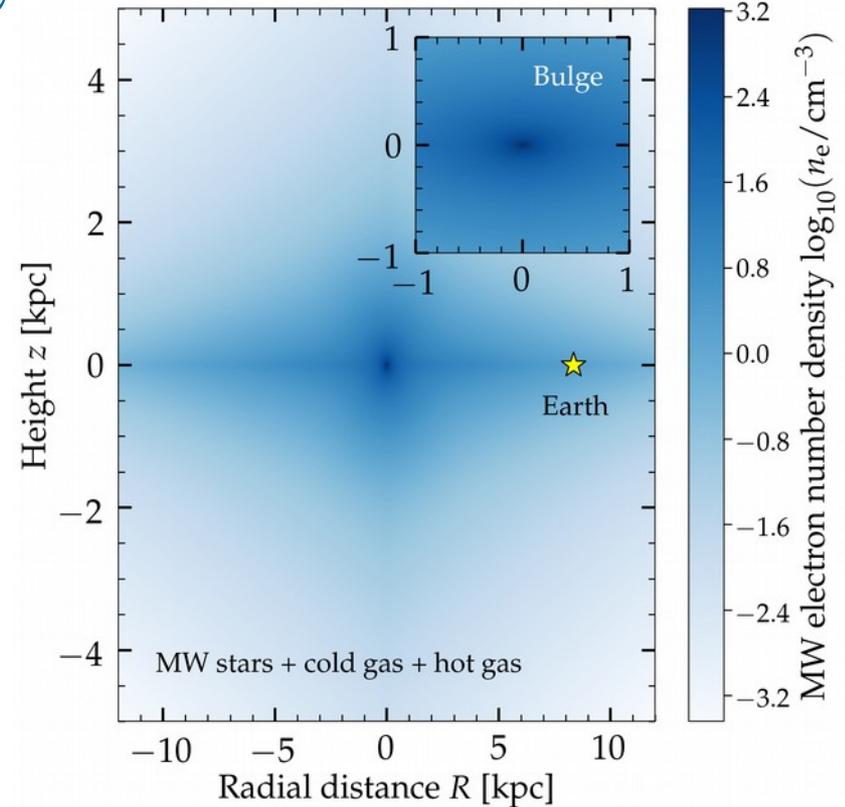
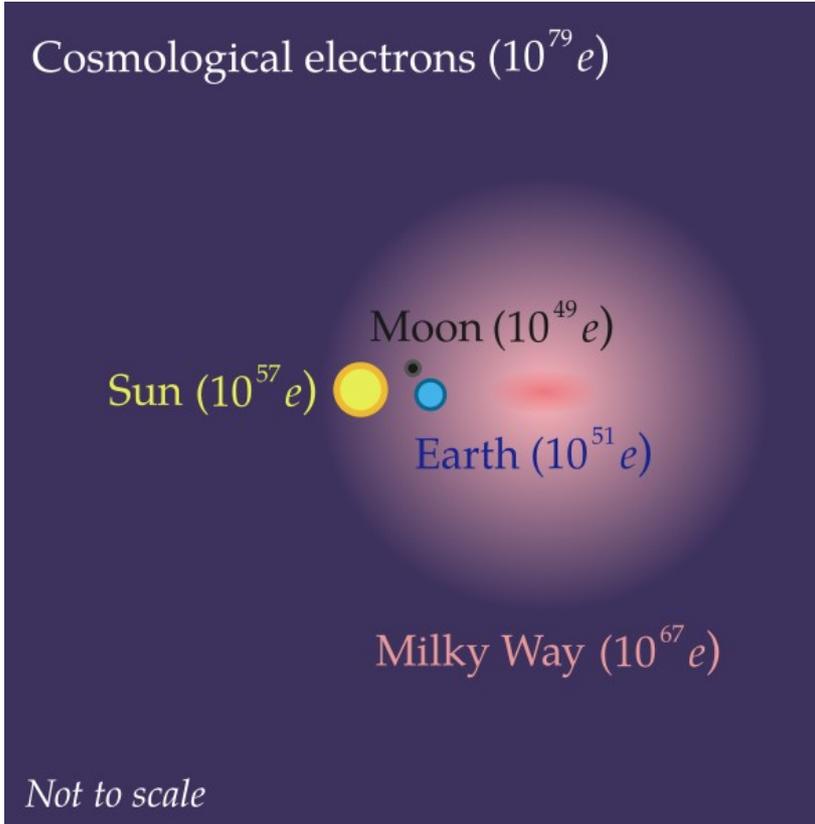
Treated as point sources of electrons

$$V_{e\beta} = V_{e\beta}^{\oplus} + V_{e\beta}^{\text{Moon}} + V_{e\beta}^{\odot}$$

The total potential

Milky Way:

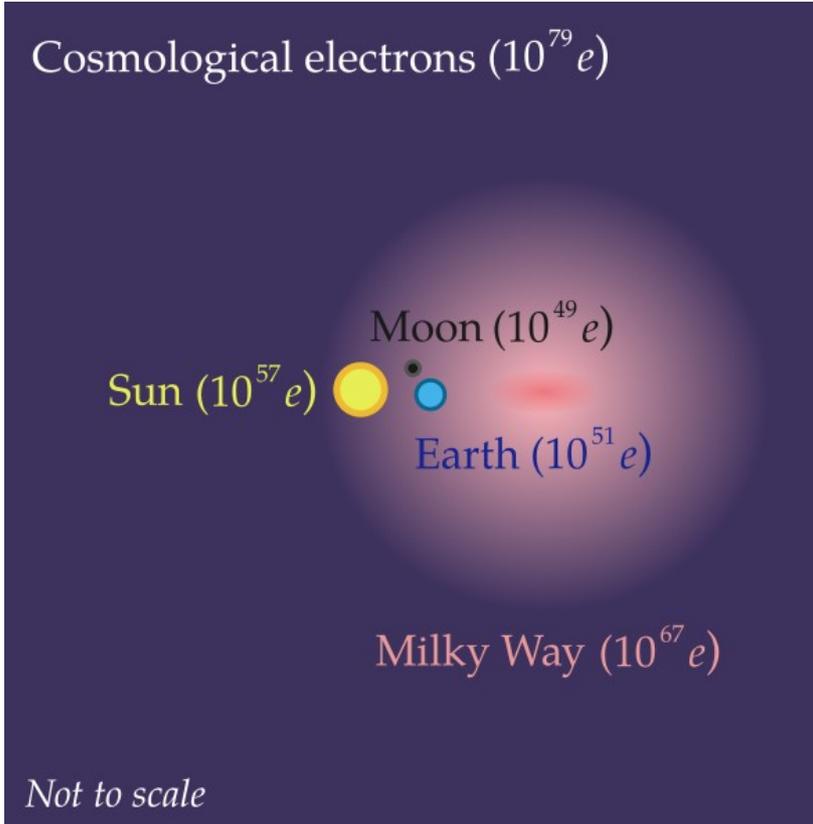
P. McMillan 2011
M.J. Miller & J.N. Bregman 2013



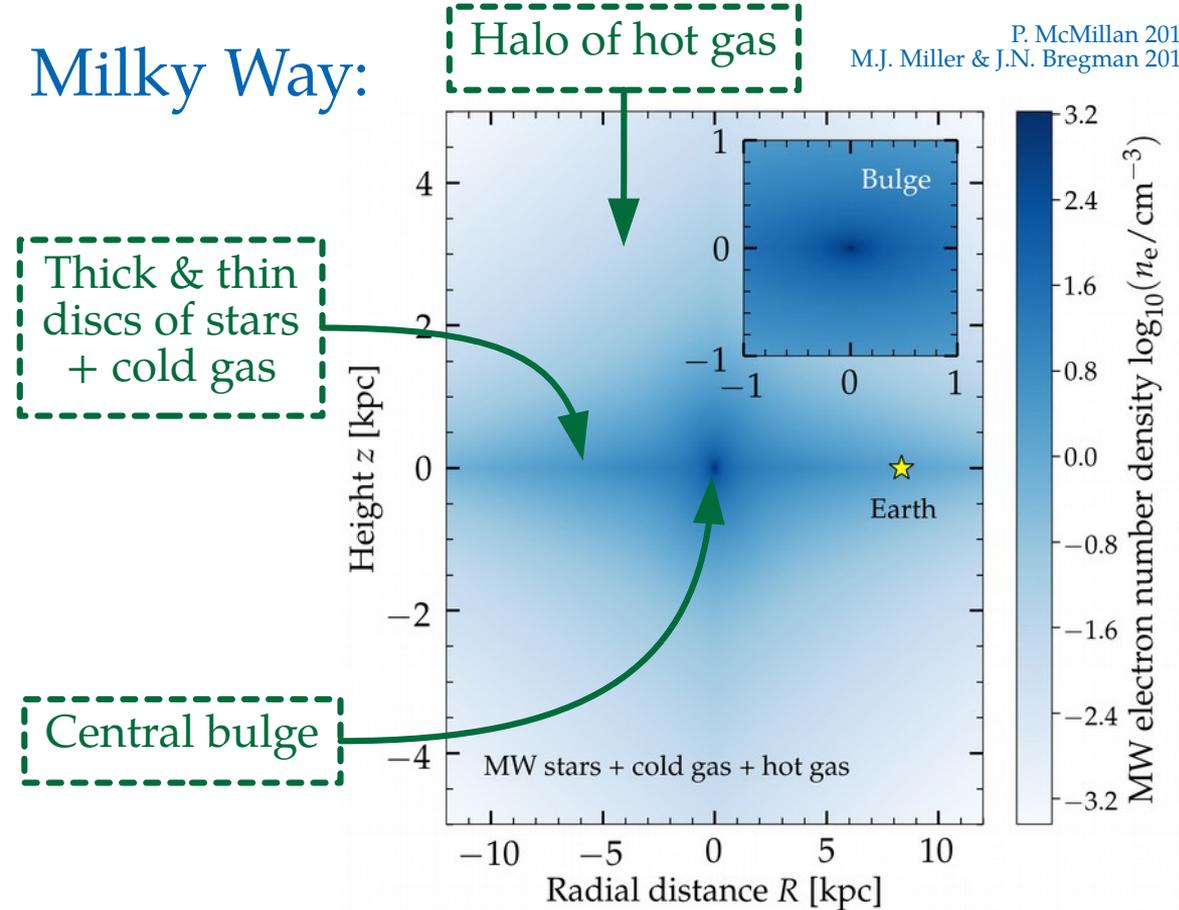
$$V_{e\beta} = V_{e\beta}^{\oplus} + V_{e\beta}^{\text{Moon}} + V_{e\beta}^{\odot} + V_{e\beta}^{\text{MW}}$$

The total potential

P. McMillan 2011
M.J. Miller & J.N. Bregman 2013

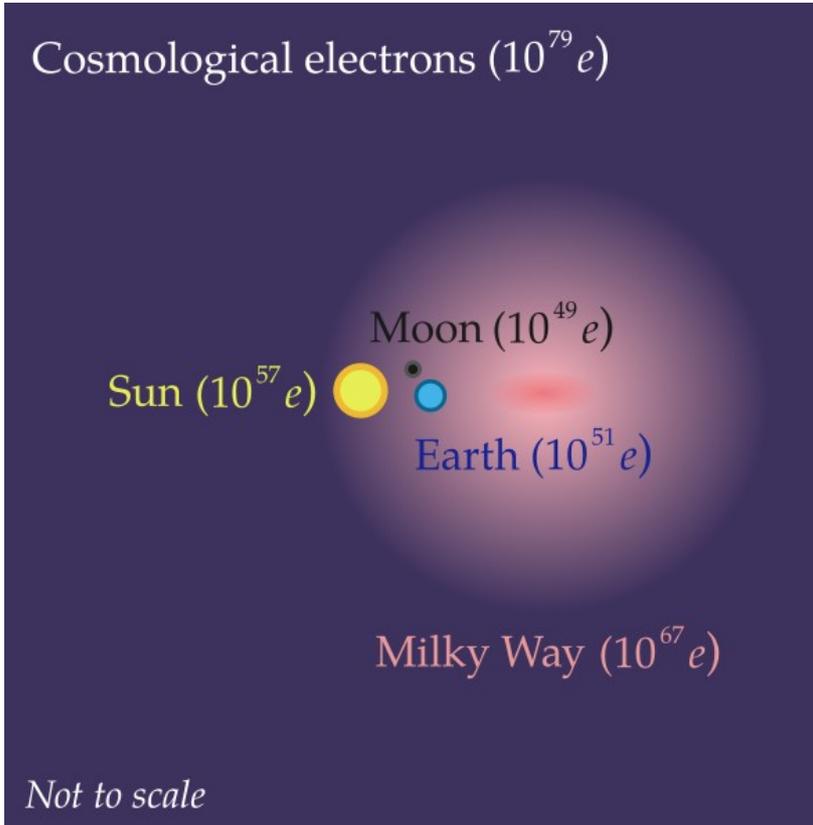


Milky Way:



$$V_{e\beta} = V_{e\beta}^{\oplus} + V_{e\beta}^{\text{Moon}} + V_{e\beta}^{\ominus} + V_{e\beta}^{\text{MW}}$$

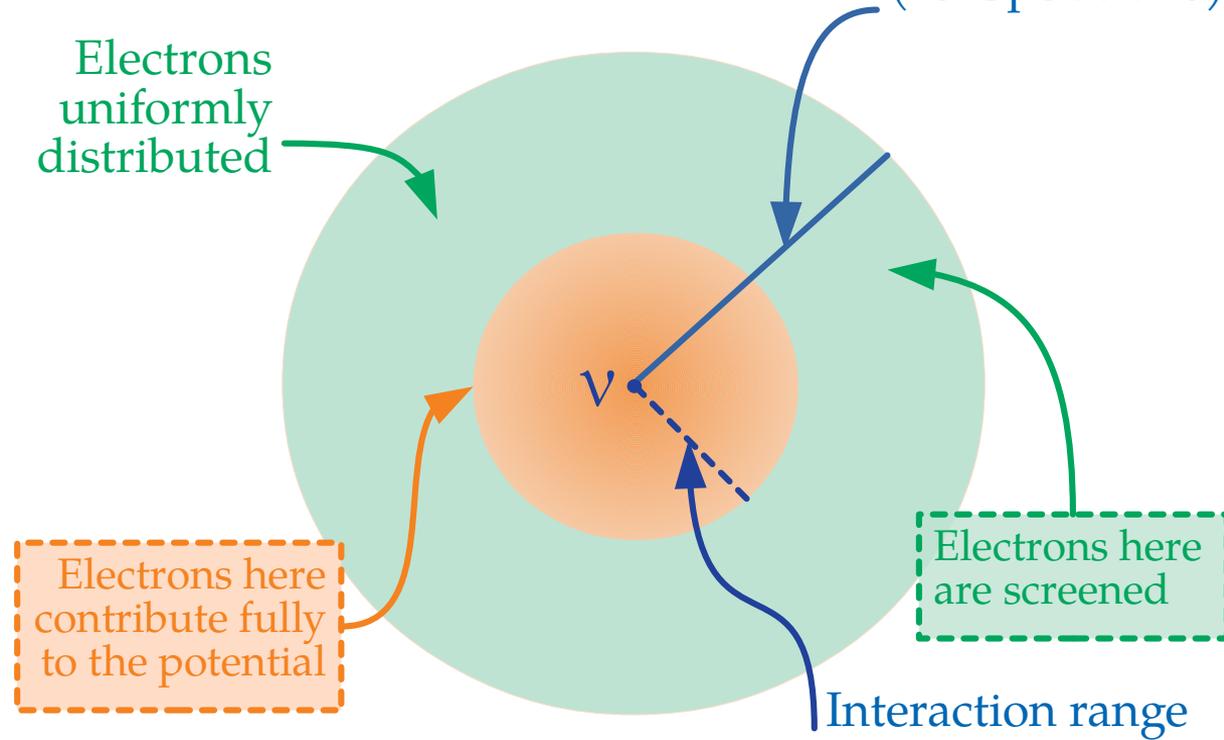
The total potential



Cosmological electrons:

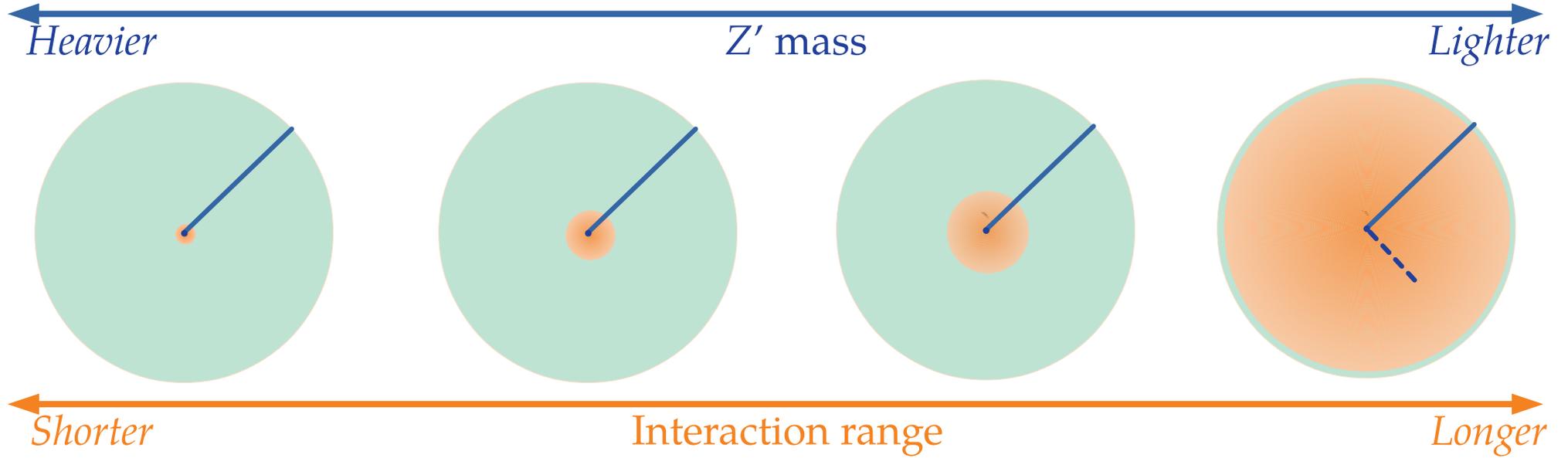
Electrons uniformly distributed

Causal horizon (15 Gpc at $z=0$)



$$V_{e\beta} = V_{e\beta}^{\oplus} + V_{e\beta}^{\text{Moon}} + V_{e\beta}^{\ominus} + V_{e\beta}^{\text{MW}} + V_{e\beta}^{\text{COS}}$$

The total potential



$$V_{e\beta} = V_{e\beta}^{\oplus} + V_{e\beta}^{\text{Moon}} + V_{e\beta}^{\ominus} + V_{e\beta}^{\text{MW}} + V_{e\beta}^{\text{cos}}$$

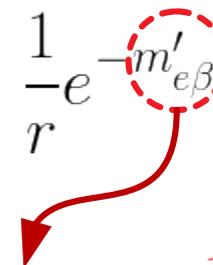
Electrons in the local and distant Universe

Potential:

$$V_{e\beta} \propto \frac{1}{r} e^{-m'_{e\beta} r}$$

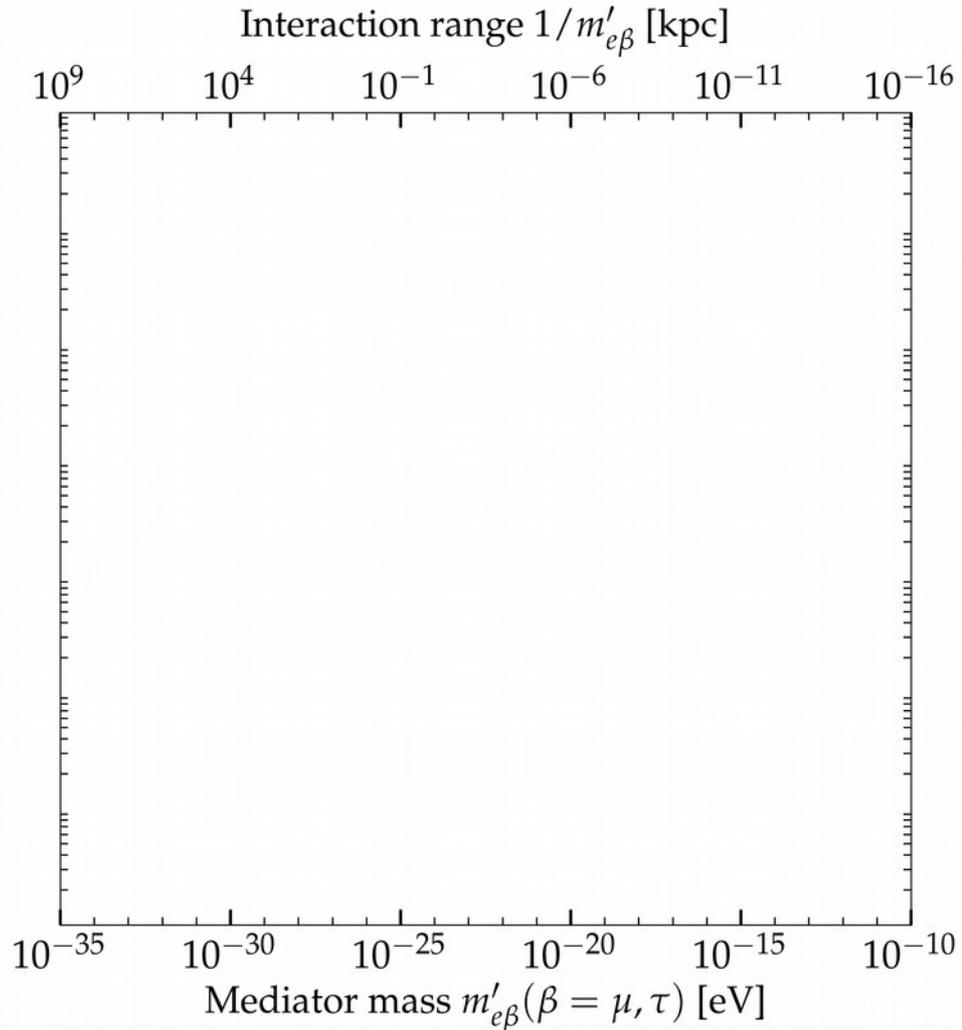
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Interaction range: $\frac{1}{m'_{e\beta}}$

Electrons in the local and distant Universe

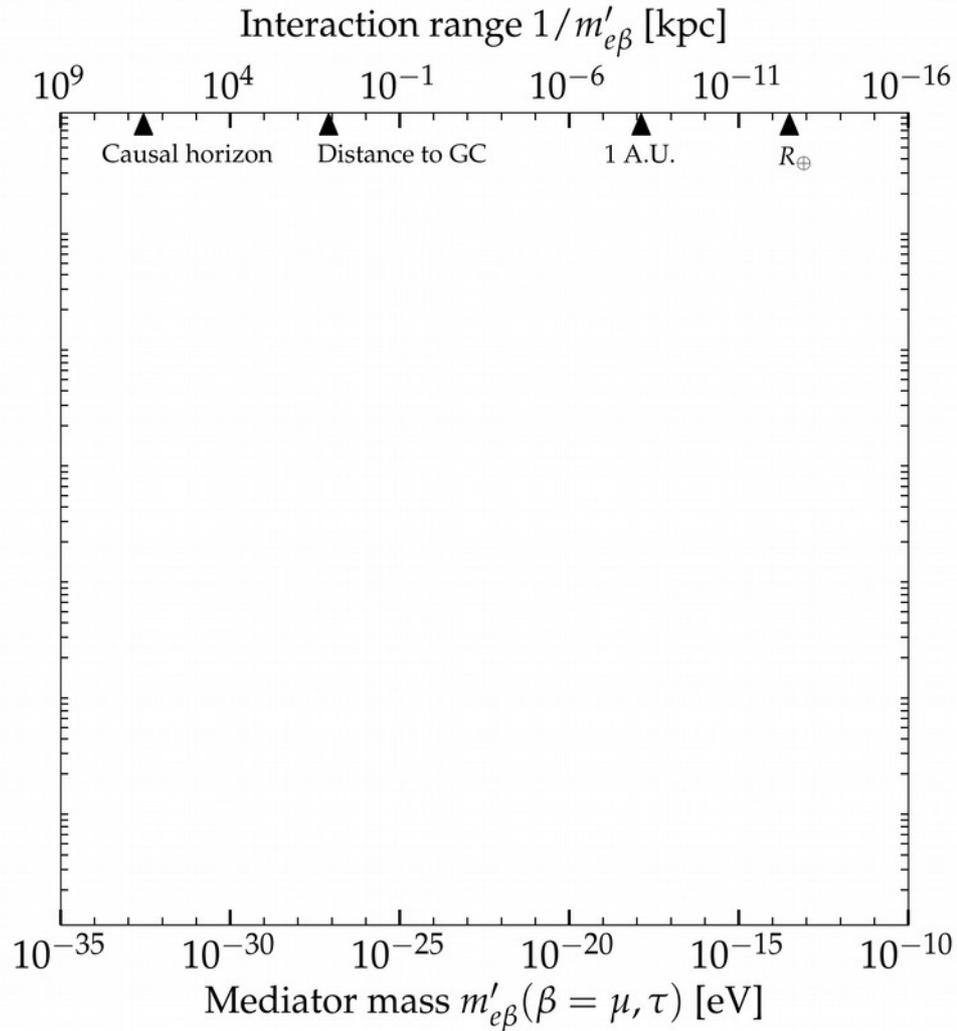


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Electrons in the local and distant Universe

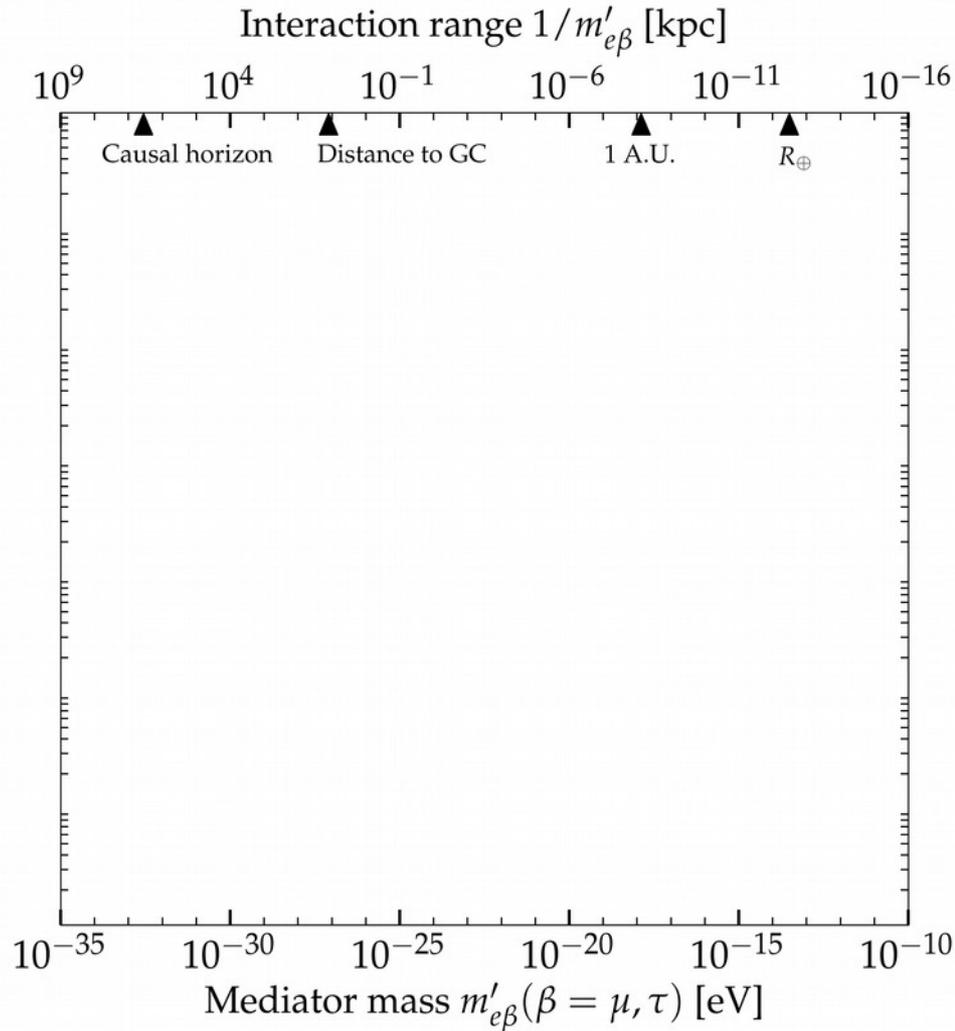


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Electrons in the local and distant Universe



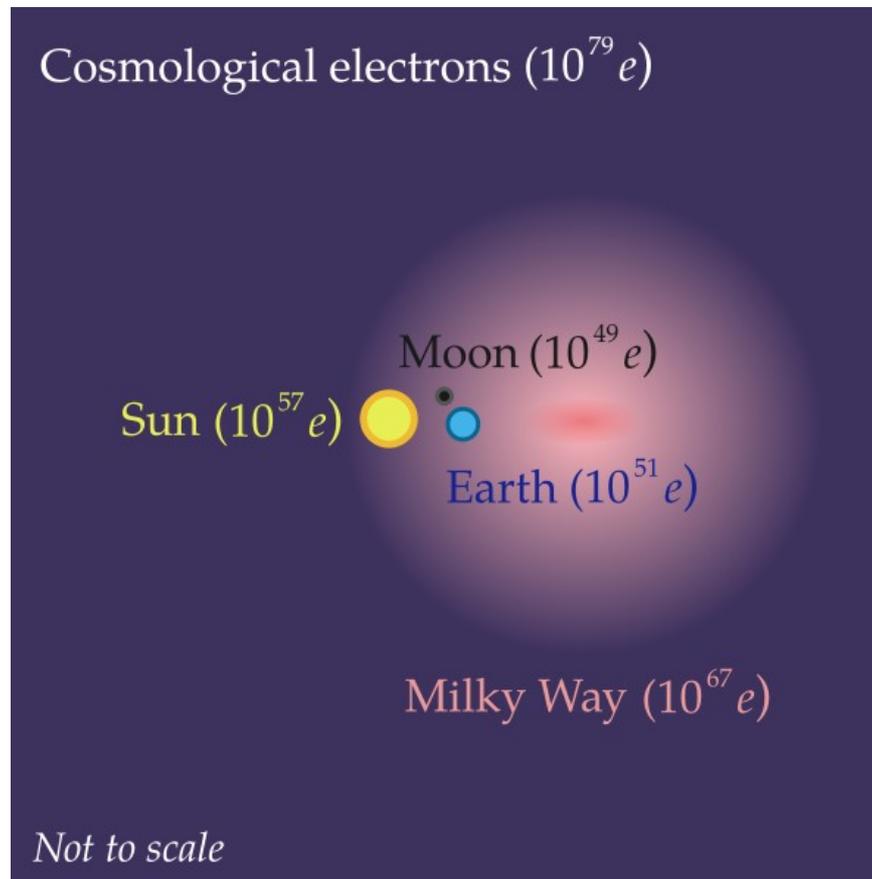
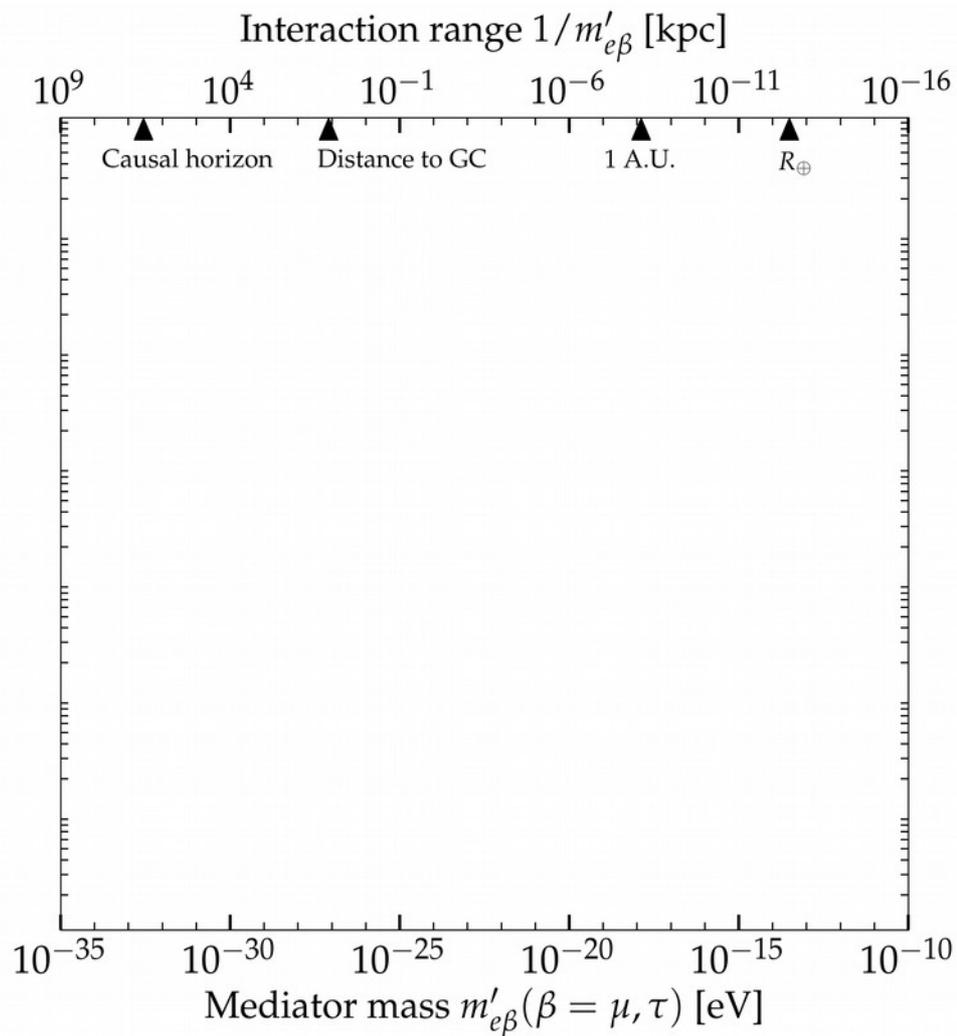
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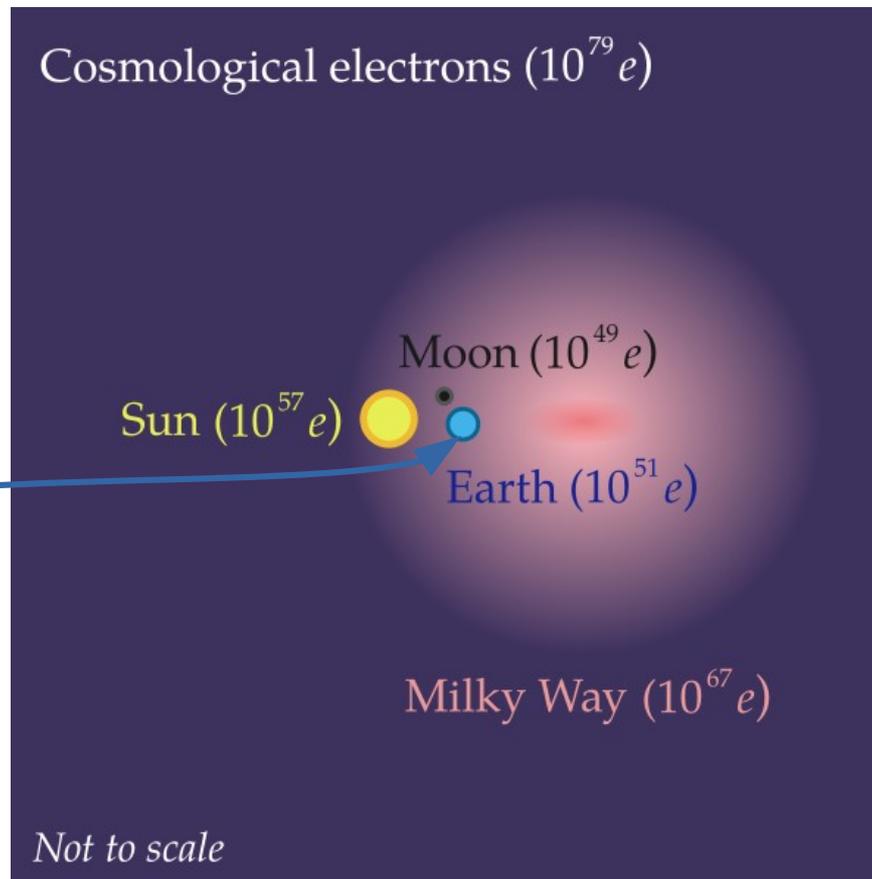
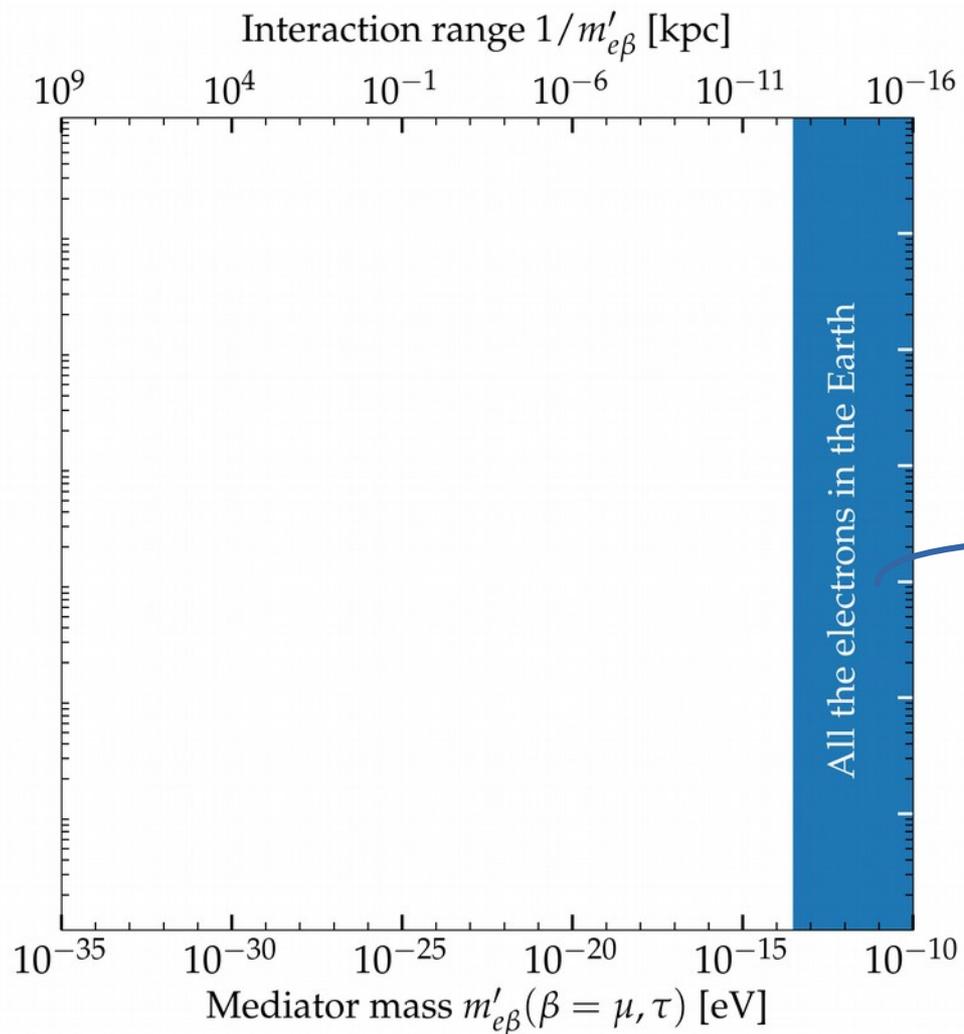
Interaction range: $\frac{1}{m'_{e\beta}}$

Light mediators
 \Rightarrow Long interaction ranges

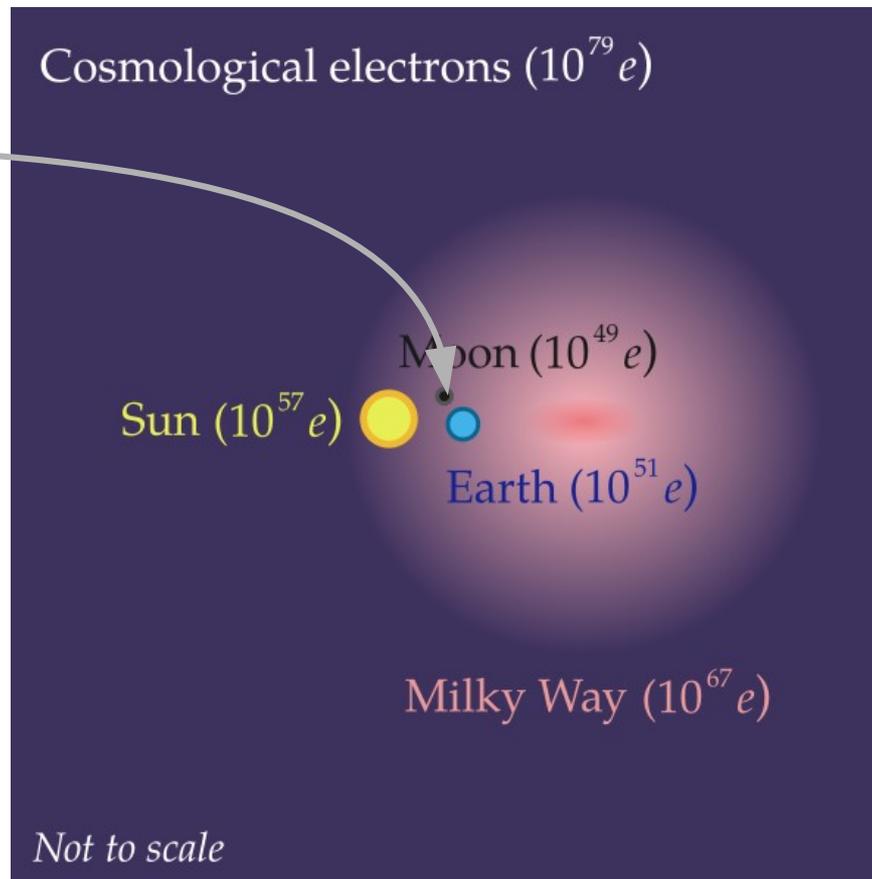
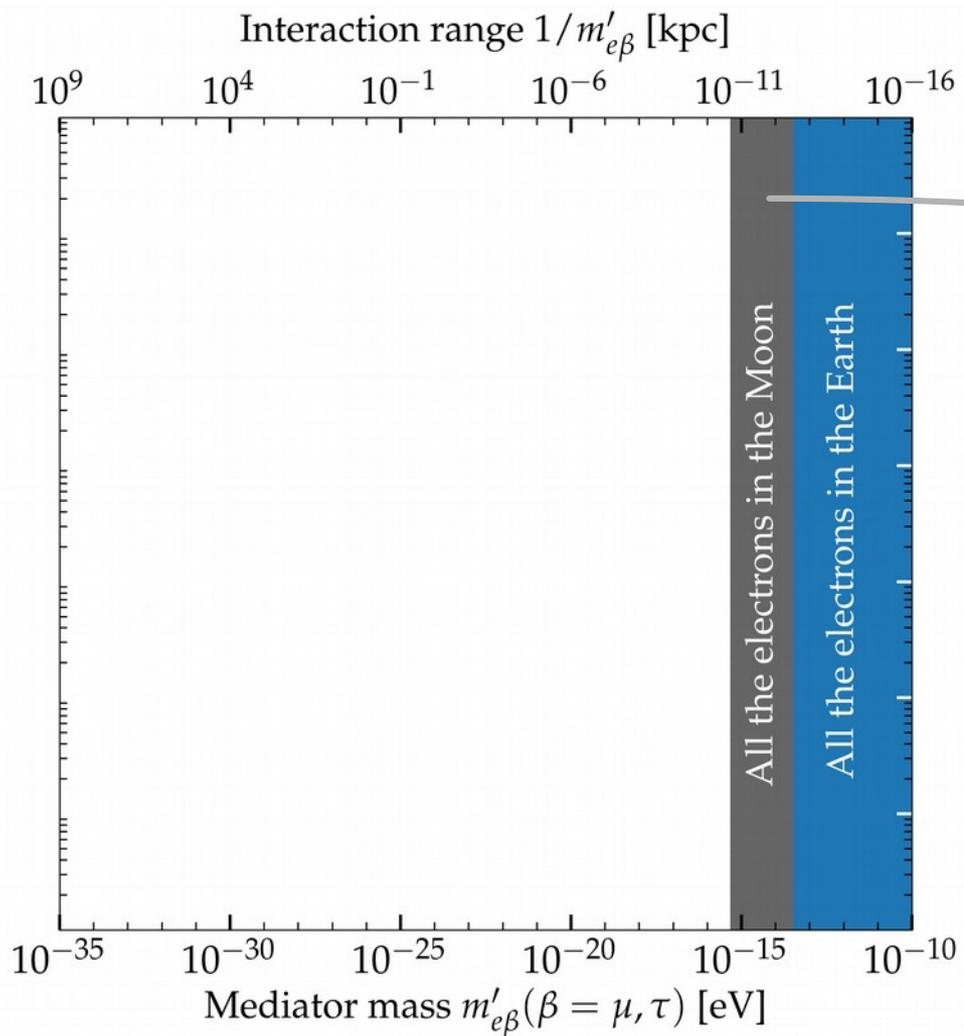
Electrons in the local and distant Universe



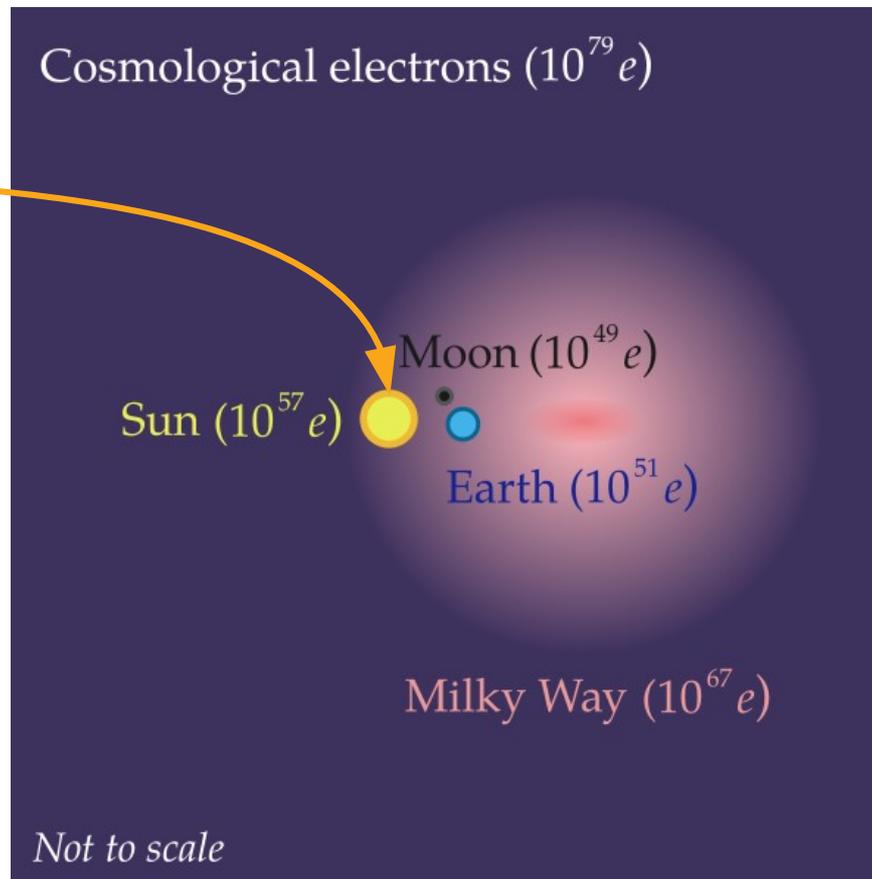
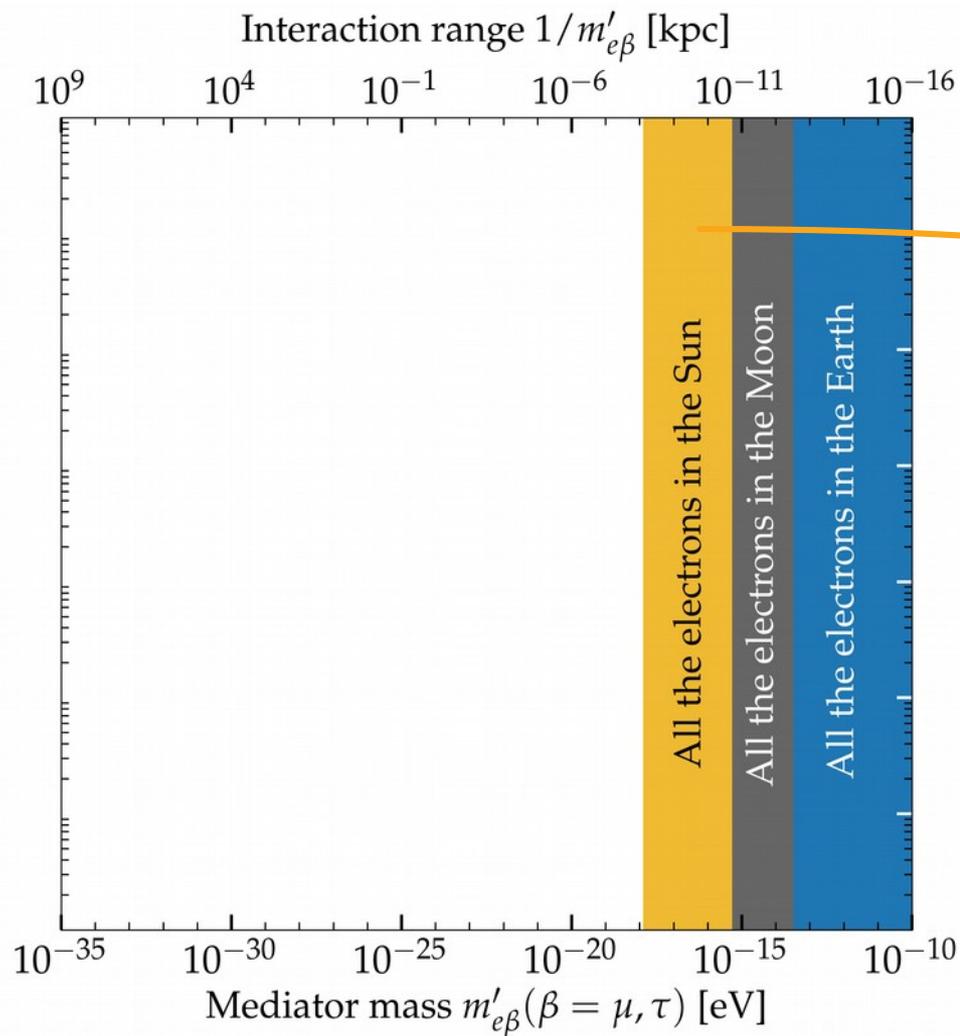
Electrons in the local and distant Universe



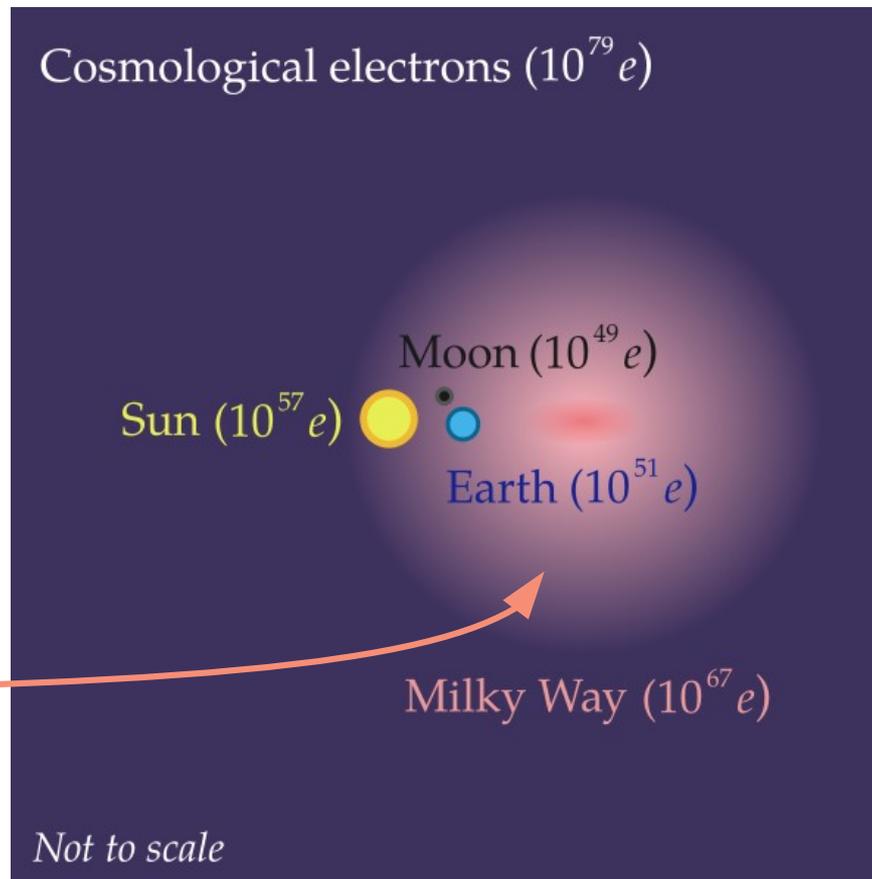
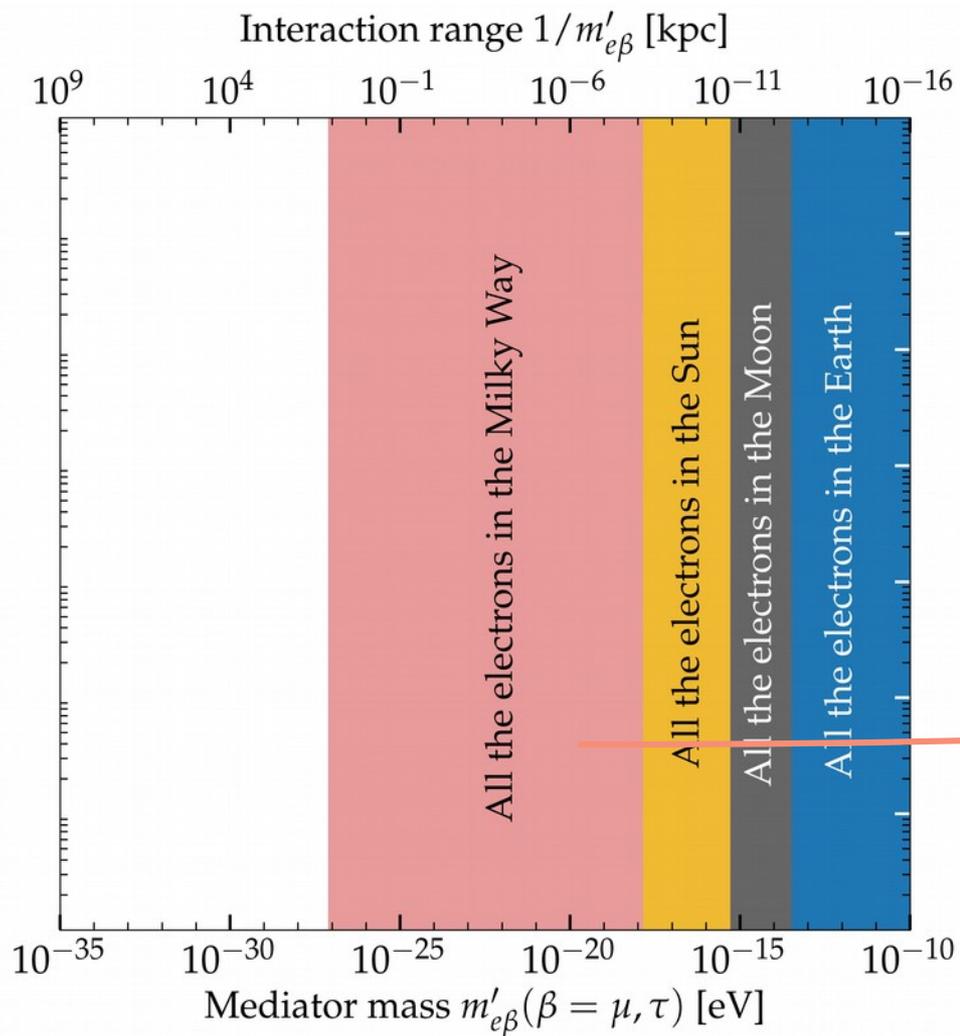
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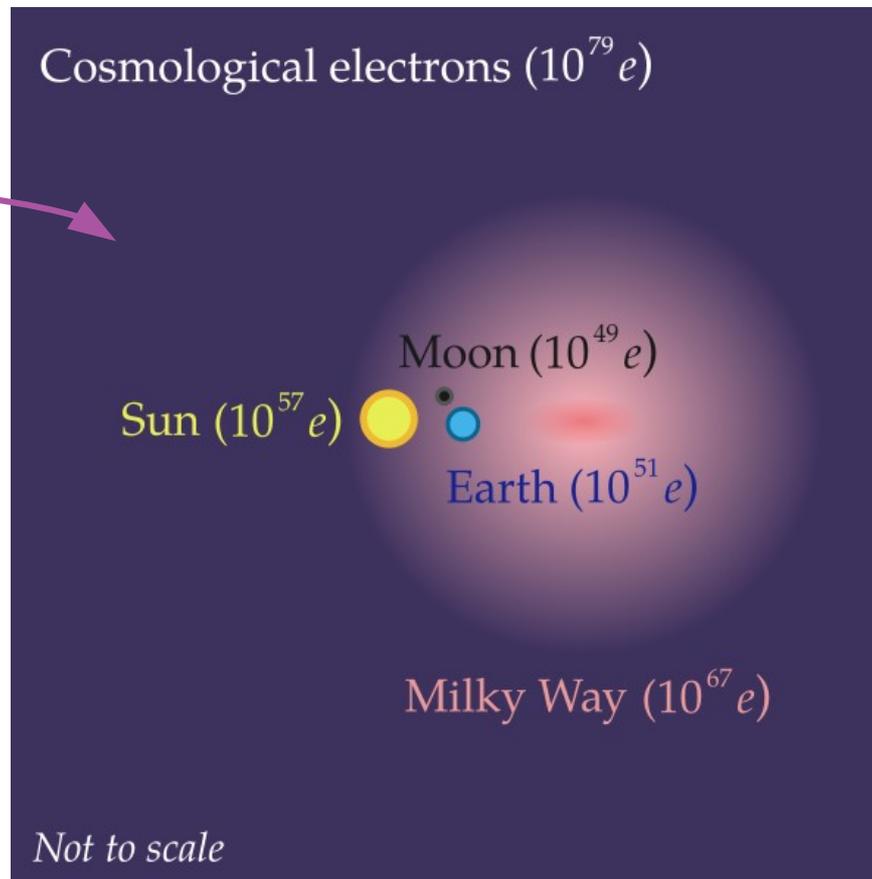
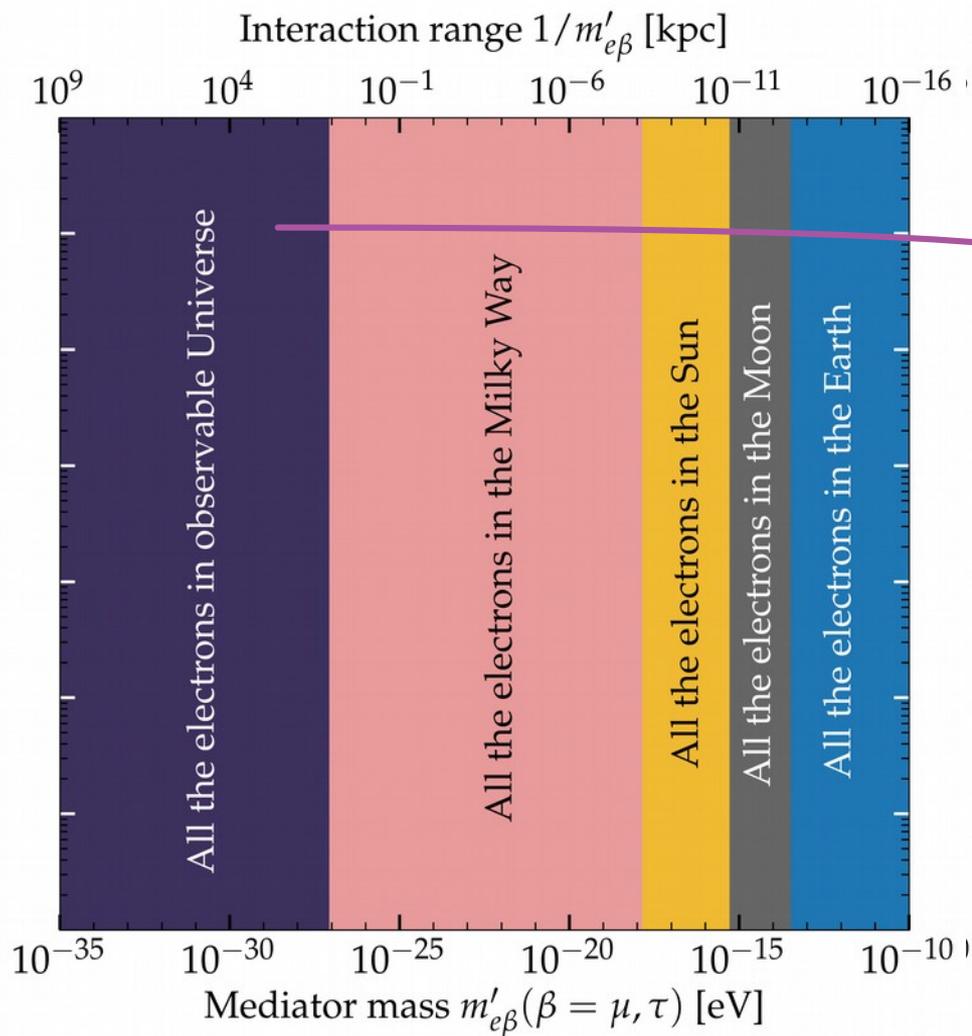
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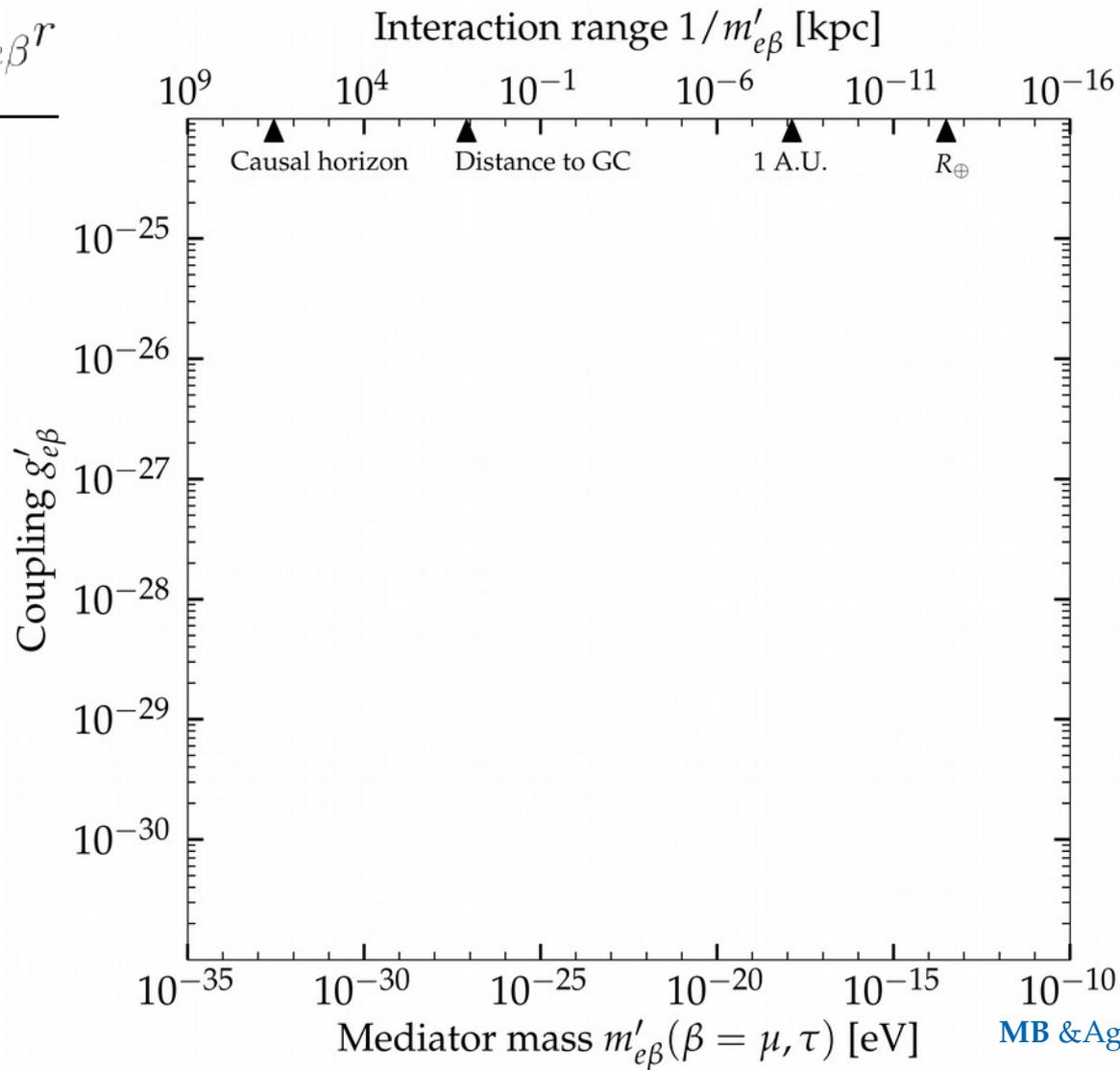
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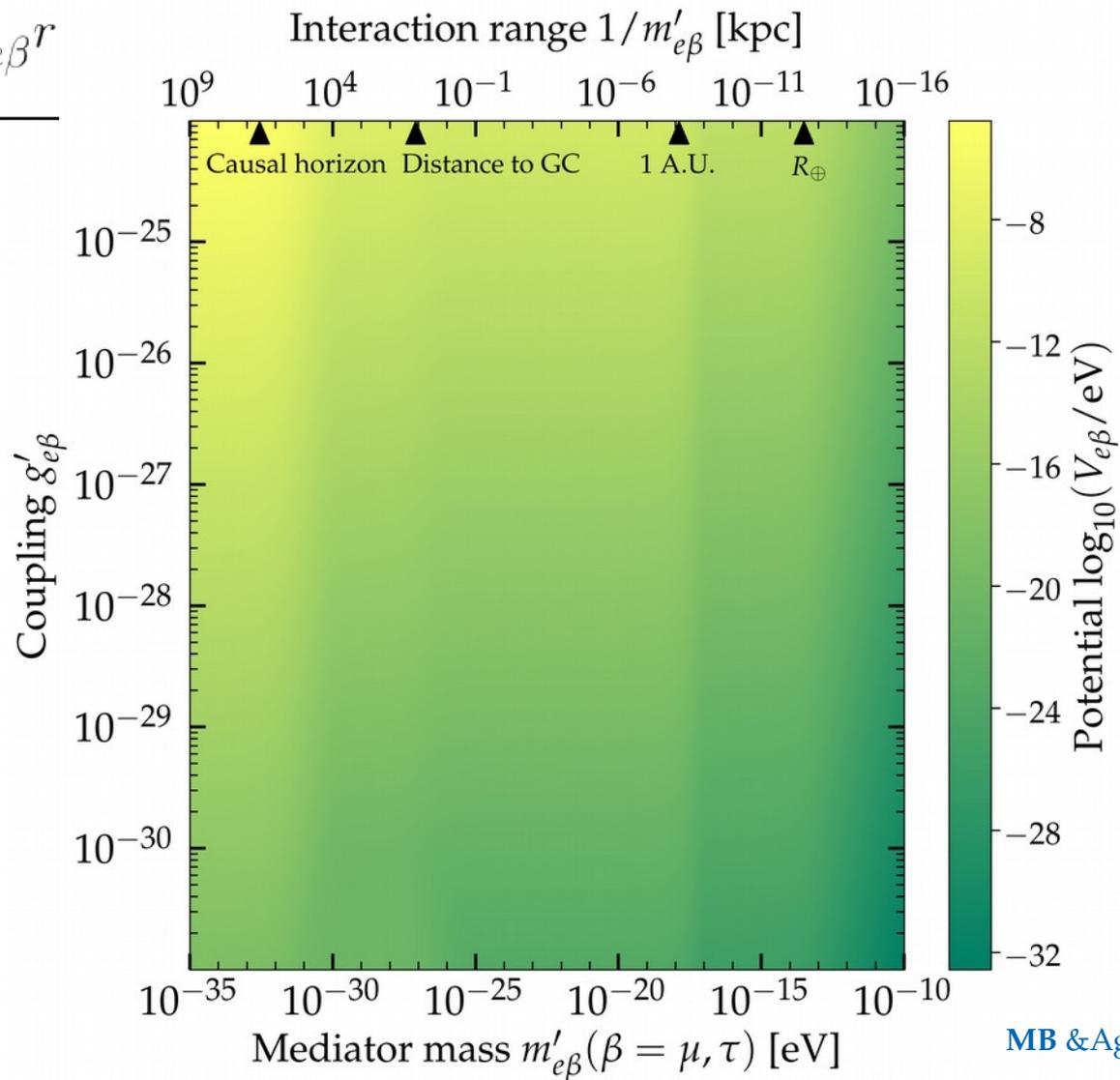
Electrons in the local and distant Universe



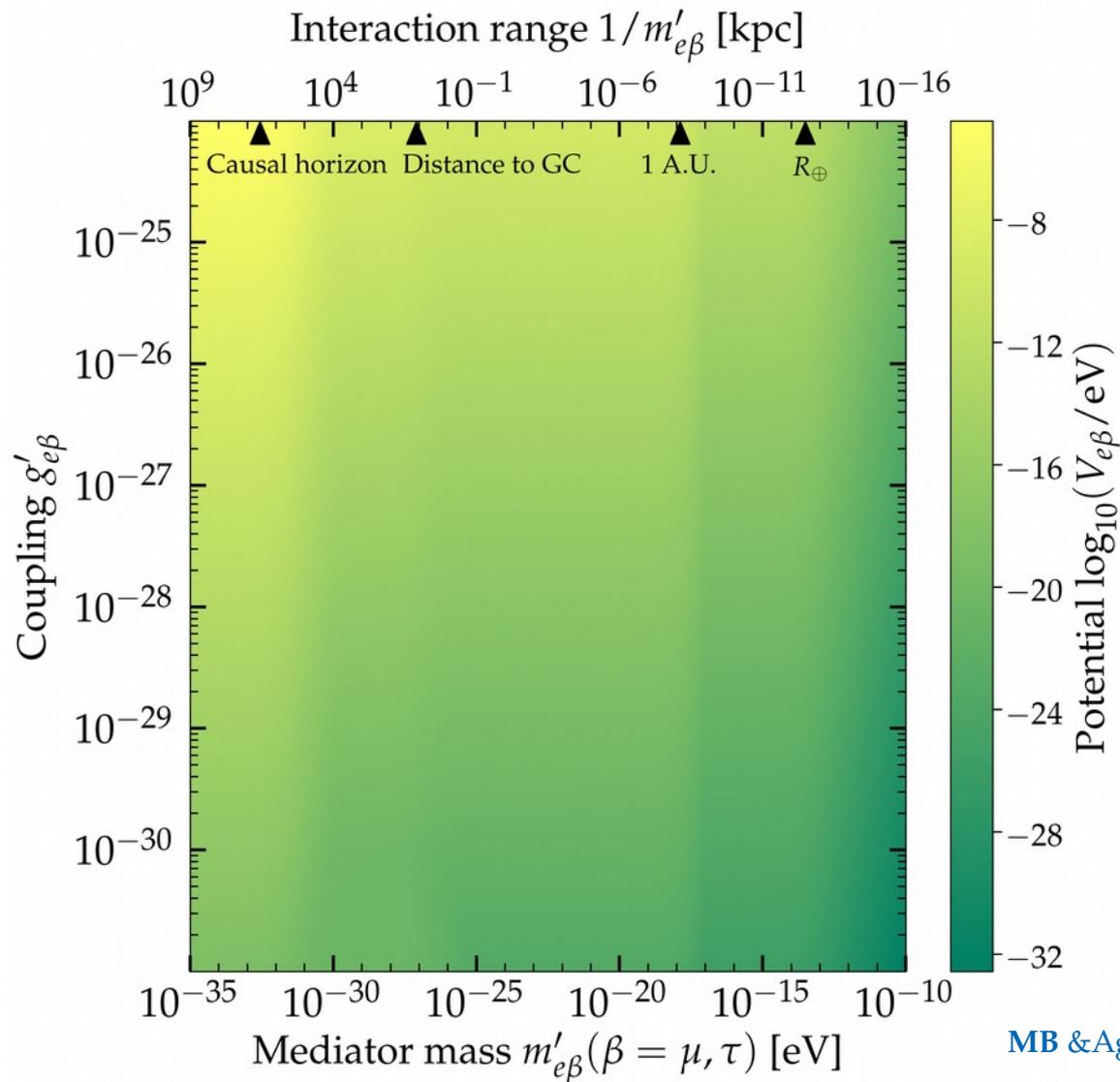
$$V_{e\beta} = \frac{g'_{e\beta}{}^2}{4\pi} \frac{e^{-m'_{e\beta}r}}{r}$$



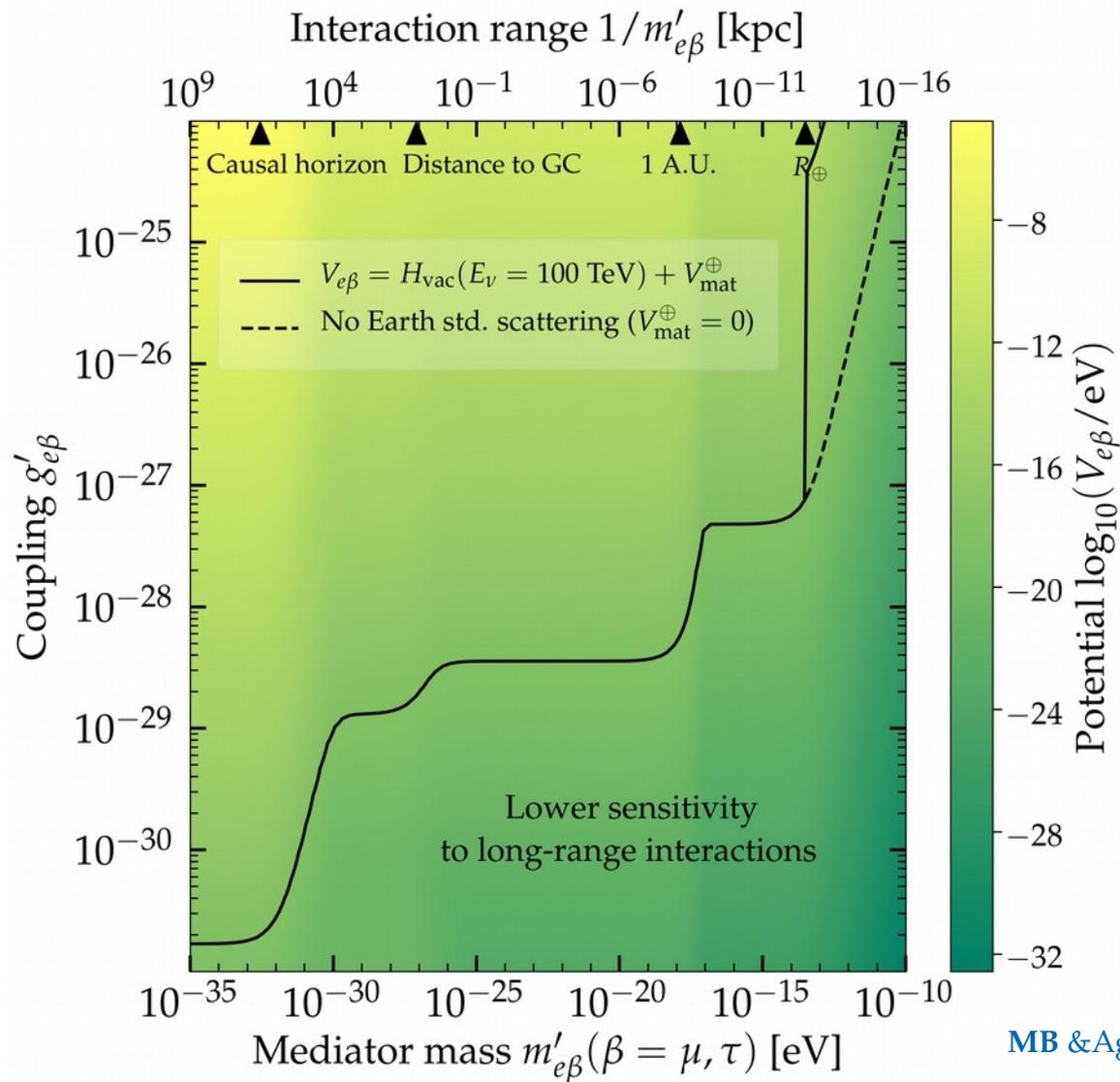
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$g_{\text{strong}} \sim 13.5$
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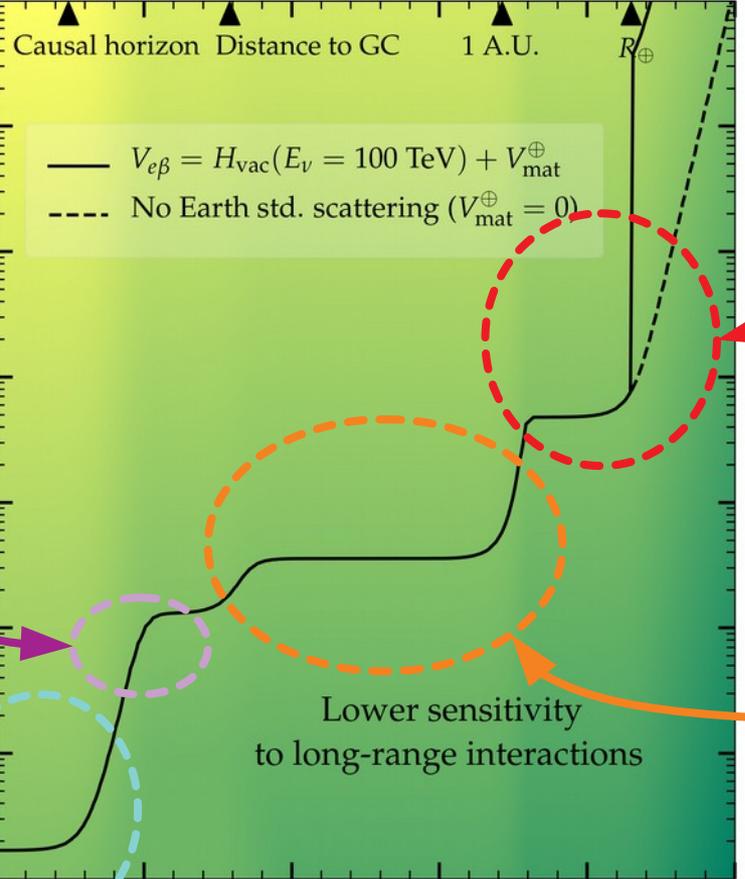



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Interaction range $1/m'_{e\beta}$ [kpc]

10^9 10^4 10^{-1} 10^{-6} 10^{-11} 10^{-16}



Dominated by Milky-Way e

Dominated by cosmological e

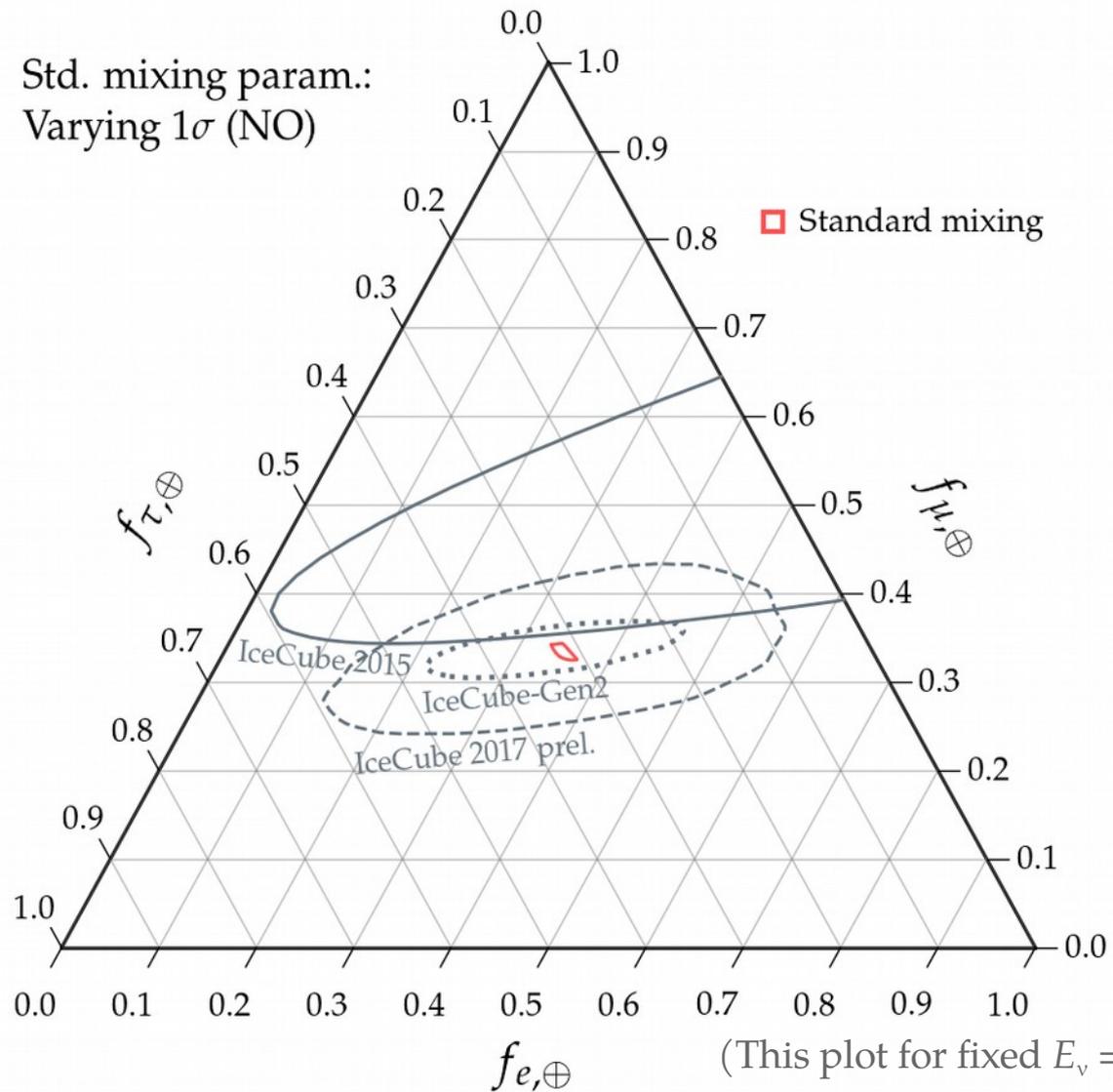
Dominated by electrons in the Earth + Moon

Dominated by solar electrons (+ Milky-Way e)

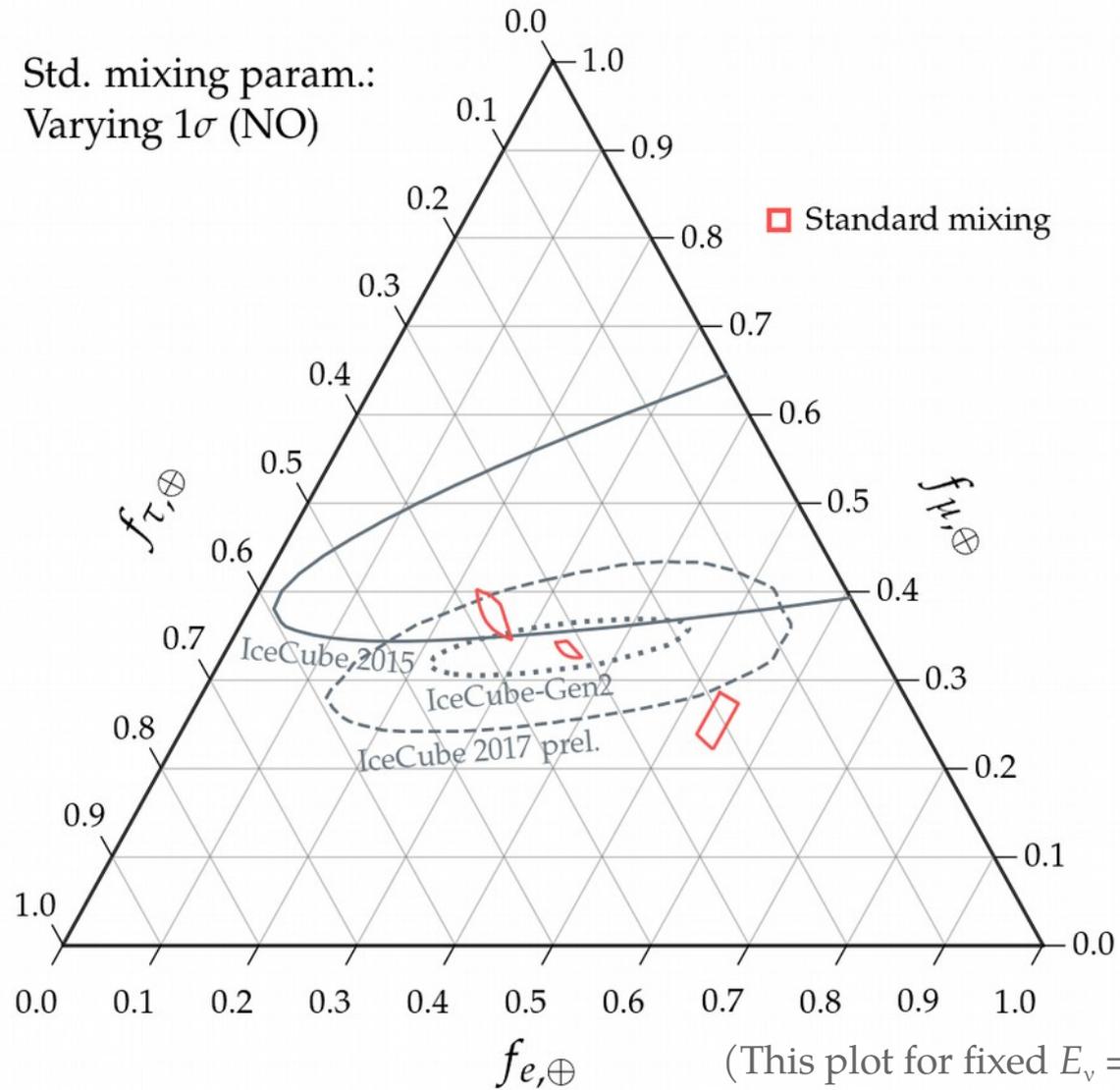
Lower sensitivity to long-range interactions

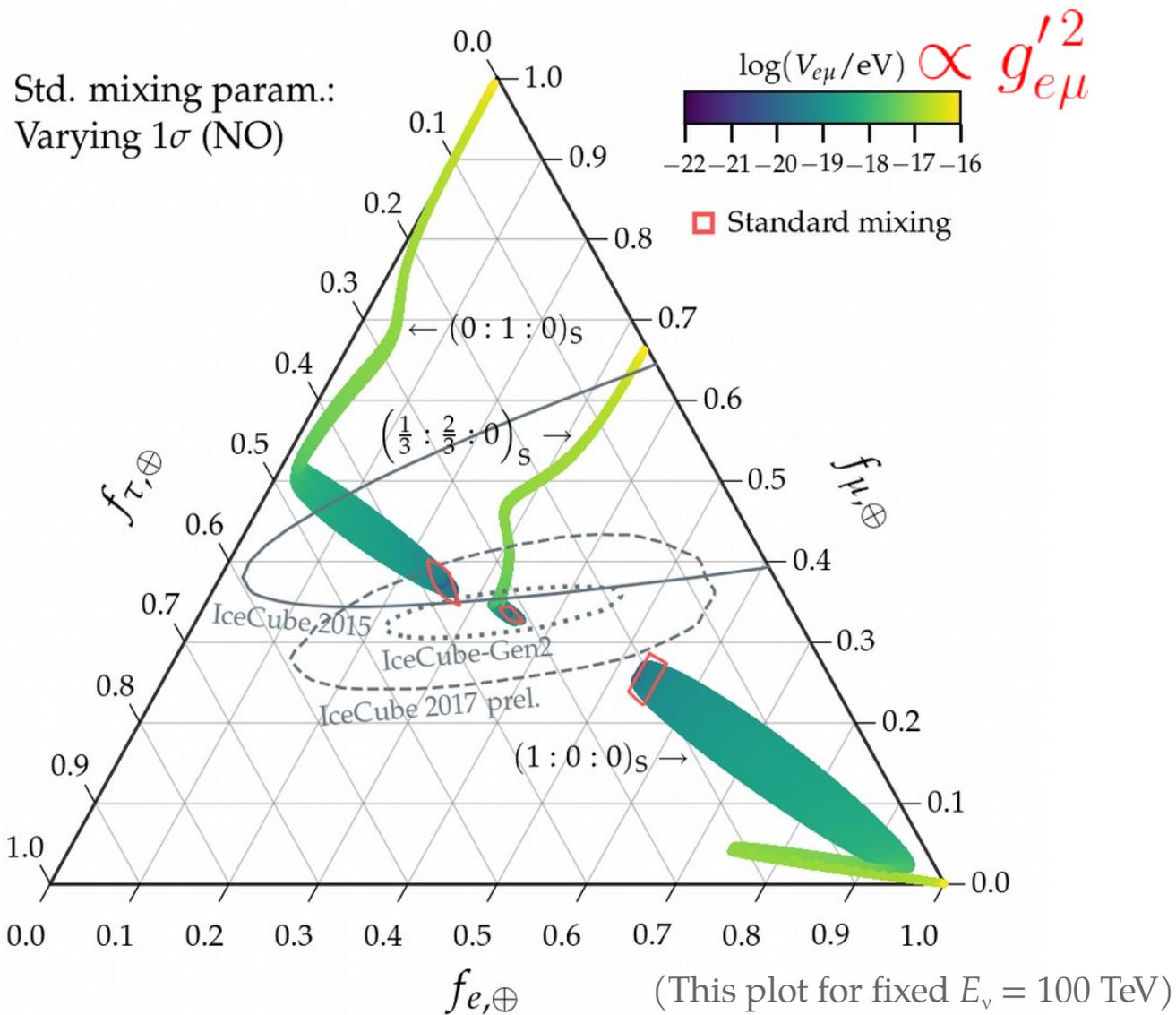
Mediator mass $m'_{e\beta}$ ($\beta = \mu, \tau$) [eV]

Std. mixing param.:
Varying 1σ (NO)

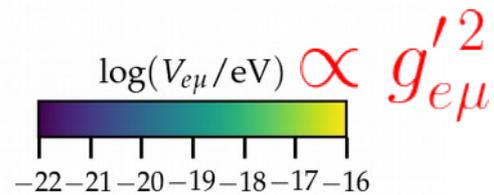


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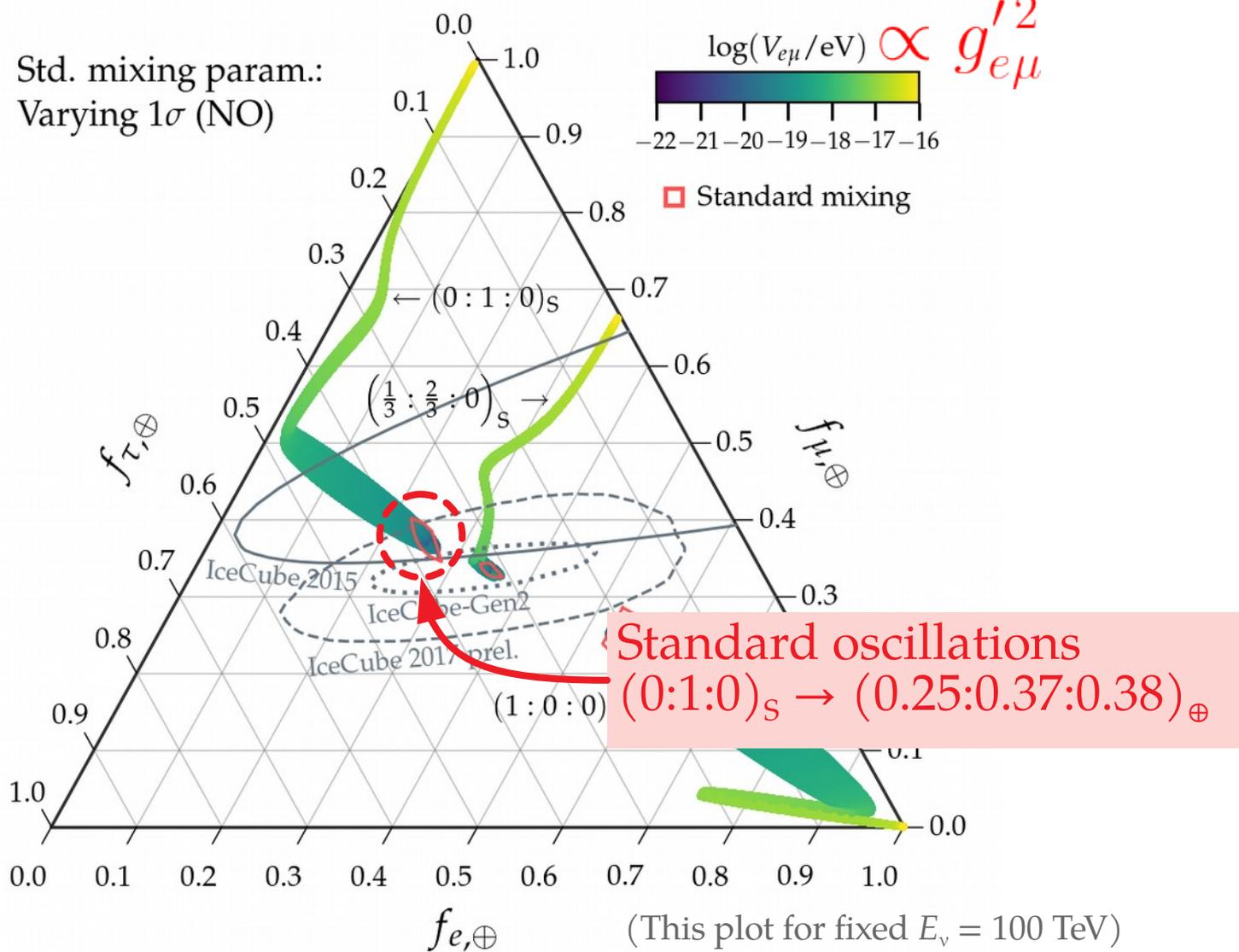




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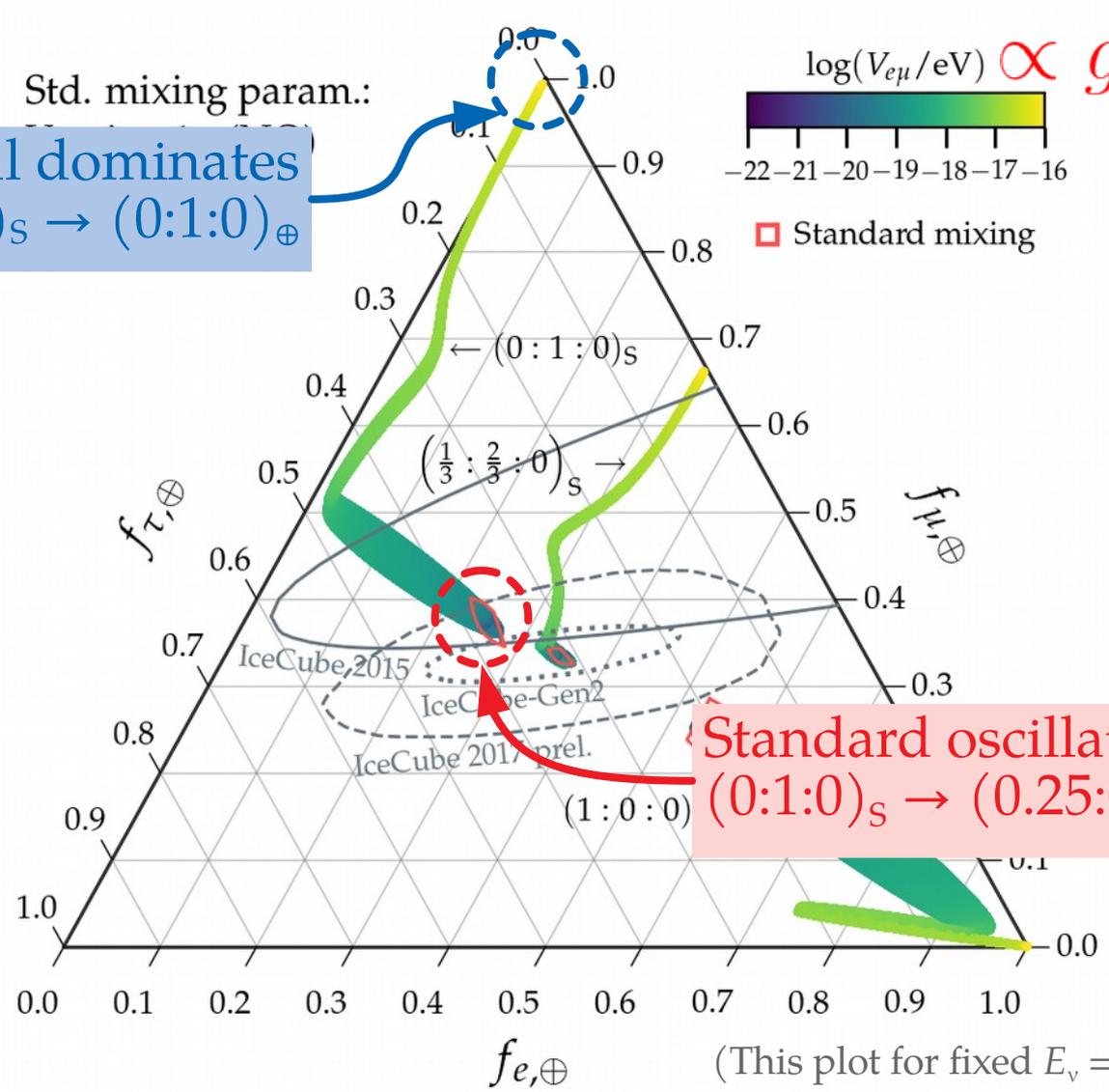
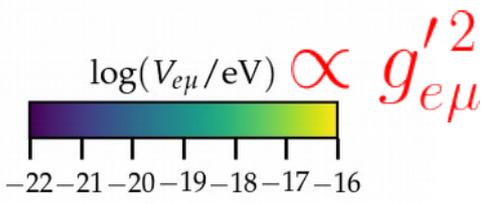


□ Standard mixing



New potential dominates
 $(0:1:0)_s \rightarrow (0:1:0)_\oplus$

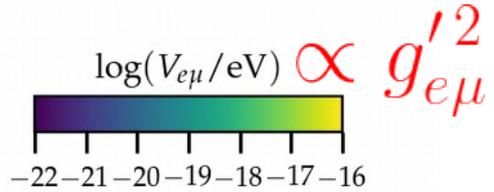
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Standard oscillations
 $(0:1:0)_s \rightarrow (0.25:0.37:0.38)_\oplus$

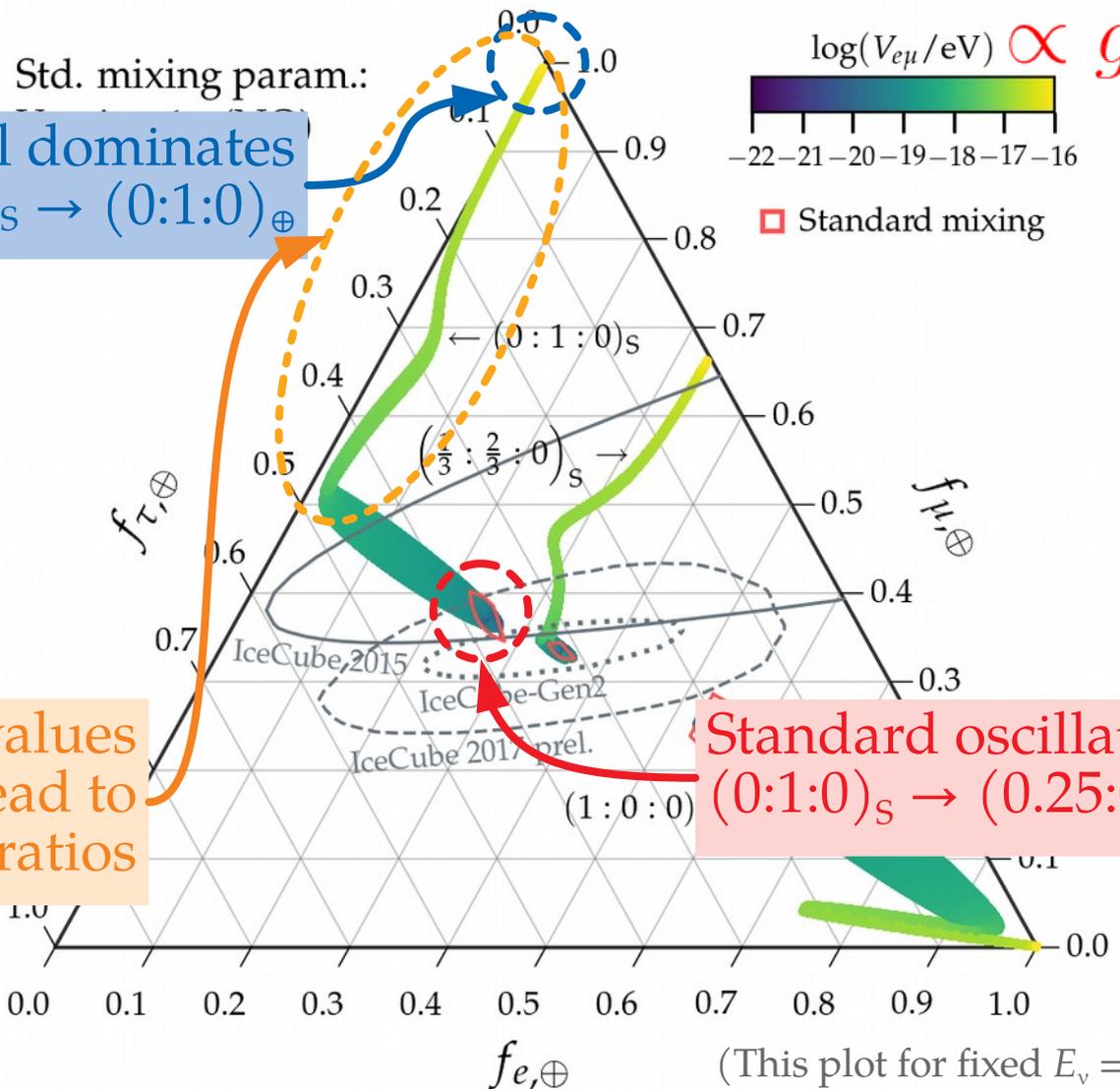
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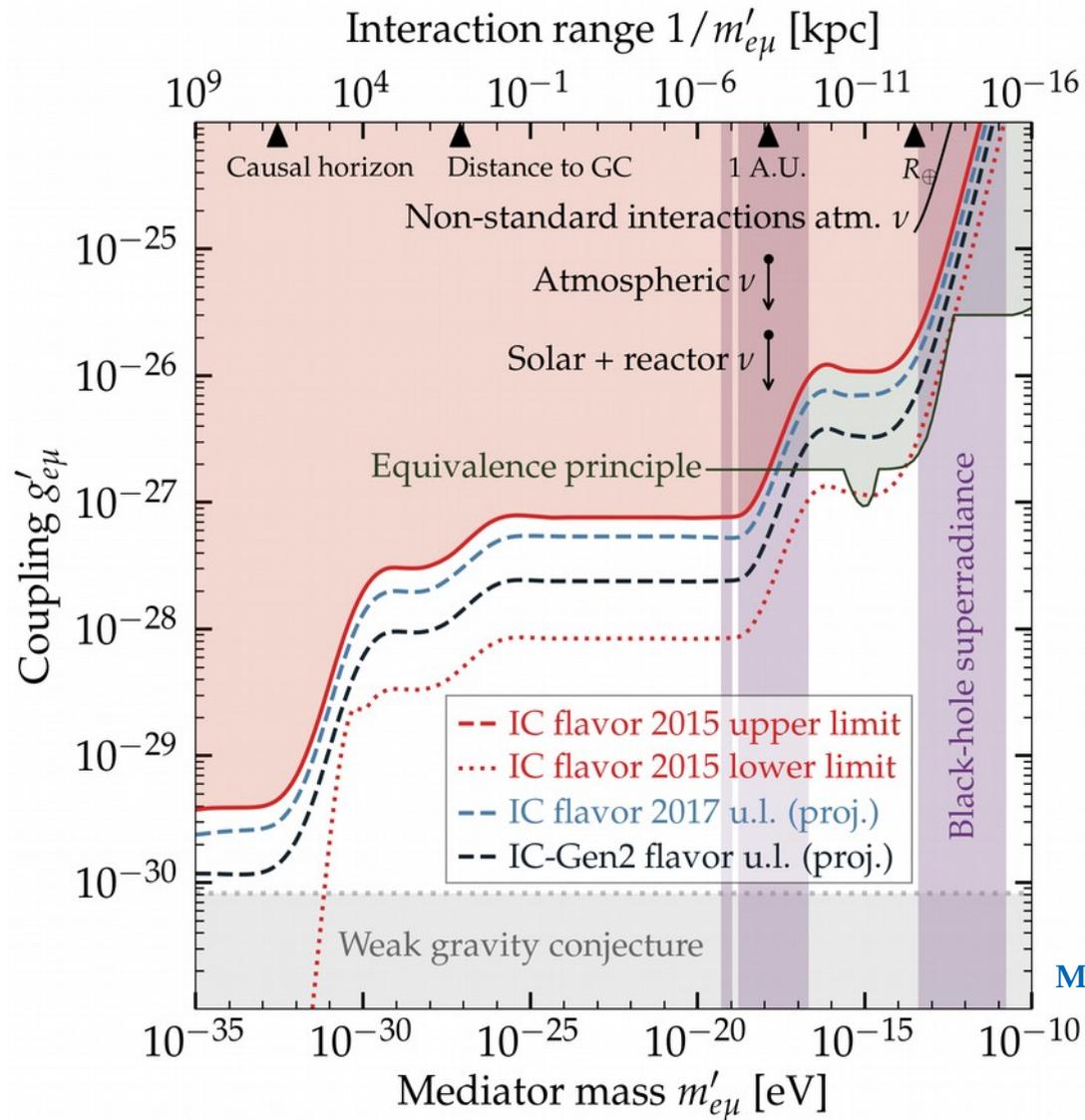
Std. mixing param.:



We can disfavor all values of m' and g' that lead to these flavor ratios

Standard oscillations
 $(0:1:0)_s \rightarrow (0.25:0.37:0.38)_\oplus$





MB & Agarwalla, PRL 2019

Bonus: Measuring the inelasticity $\langle y \rangle$

- ▶ Inelasticity in CC ν_μ interaction $\nu_\mu + N \rightarrow \mu + X$:

$$E_X = y E_\nu \quad \text{and} \quad E_\mu = (1-y) E_\nu \quad \Rightarrow \quad y = (1 + E_\mu/E_X)^{-1}$$

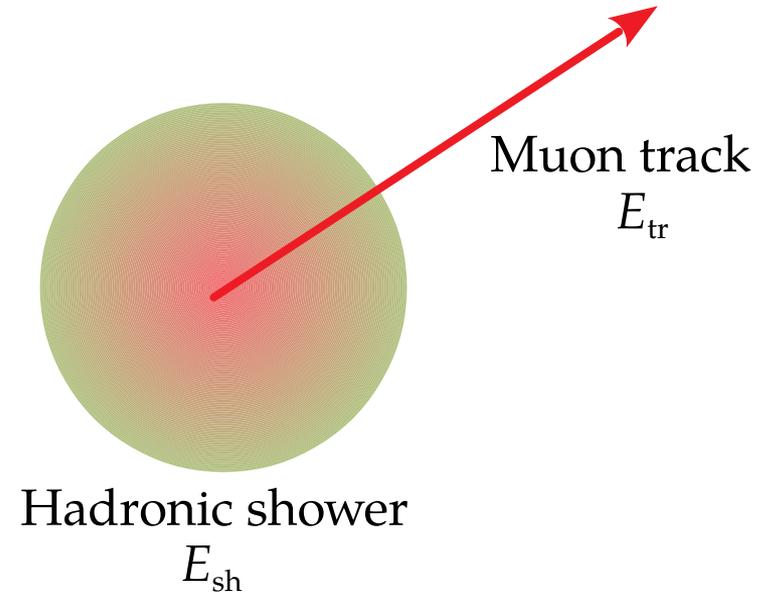
- ▶ The value of y follows a distribution $d\sigma/dy$

- ▶ In a HESE starting track:

$$\left. \begin{array}{l} E_X = E_{\text{sh}} \text{ (energy of shower)} \\ E_\mu = E_{\text{tr}} \text{ (energy of track)} \end{array} \right\} y = (1 + E_{\text{tr}}/E_{\text{sh}})^{-1}$$

- ▶ New IceCube analysis:

- ▶ 5 years of starting-track data (2650 tracks)
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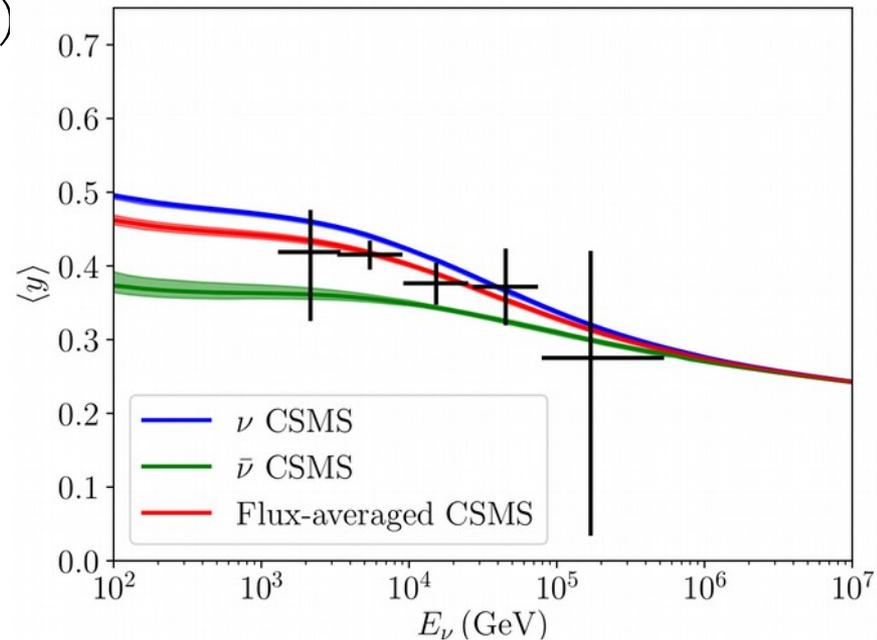
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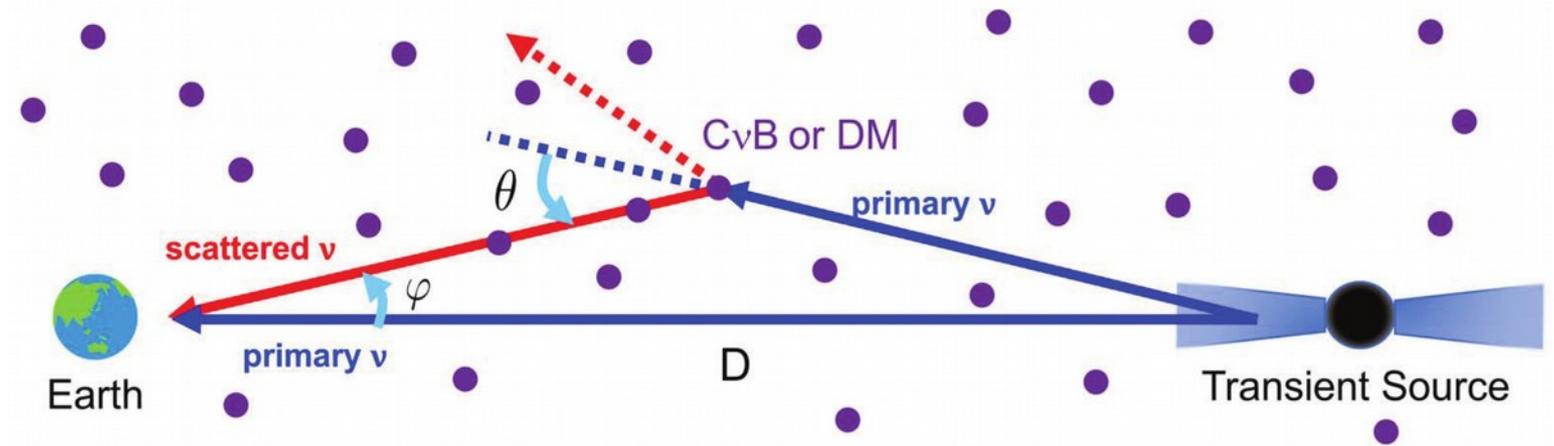
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IceCube, PRD 2019

New physics in timing — TeV–PeV

Multiple secret $\nu\nu$ scatterings may delay the arrival of neutrinos from a transient



Shoemaker & Murase, 1903.08607

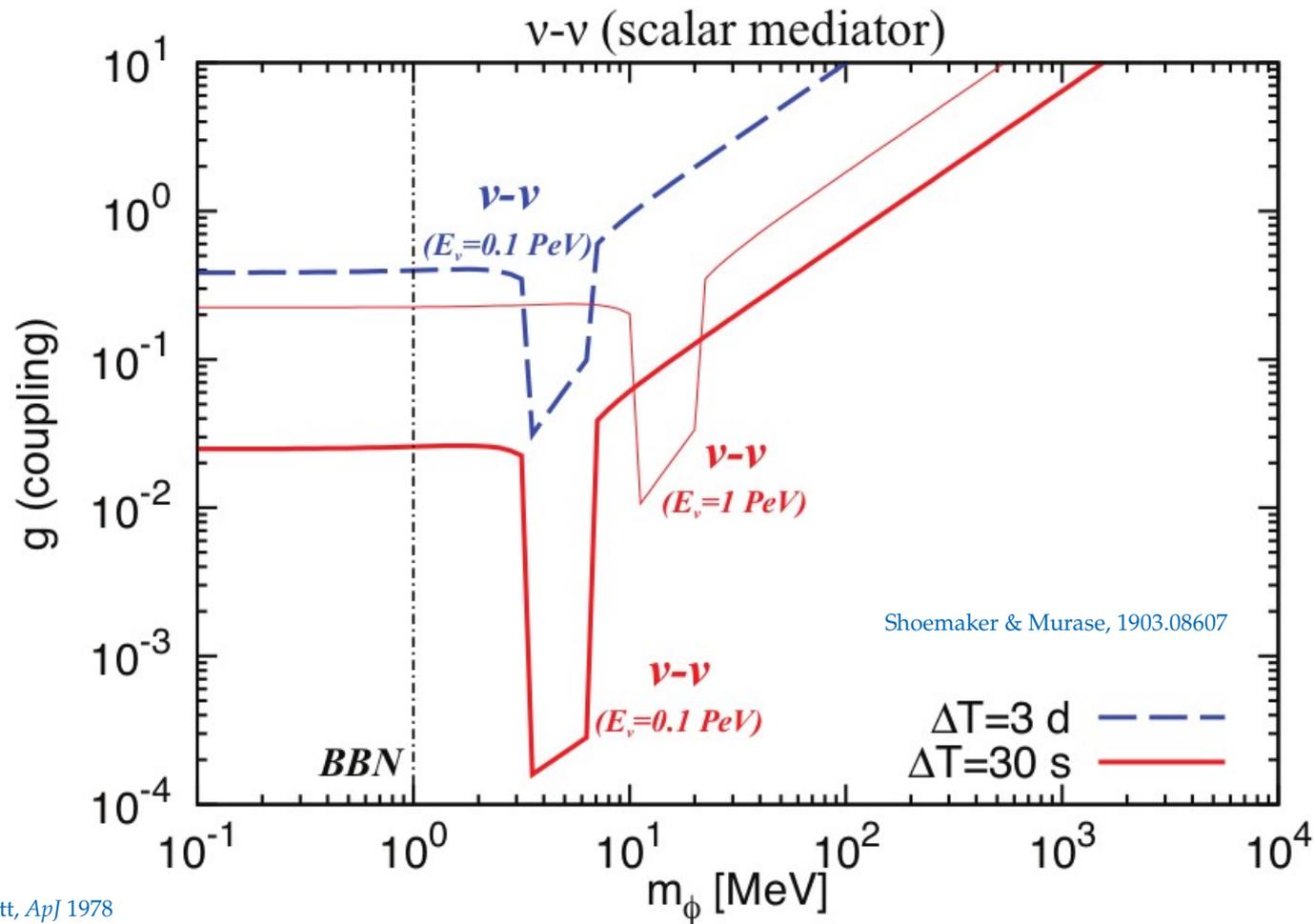
Characteristic time delay —

$$\text{Optical depth to } \nu\nu: \tau_{\nu\nu} = n_\nu \sigma_{\nu\nu} D$$

$$\Delta t \approx 1500 \text{ s} \left(\frac{\tau_{\nu\nu}}{30} \right) \left(\frac{D}{3 \text{ Gpc}} \right) \left(\frac{m_\nu}{0.1 \text{ eV}} \right) \left(\frac{0.1 \text{ PeV}}{E_\nu} \right)$$

See also: Alcock & Hatchett, *ApJ* 1978

New physics in timing — TeV–PeV



See also: Alcock & Hatchett, *ApJ* 1978

Neutrino zenith angle distribution

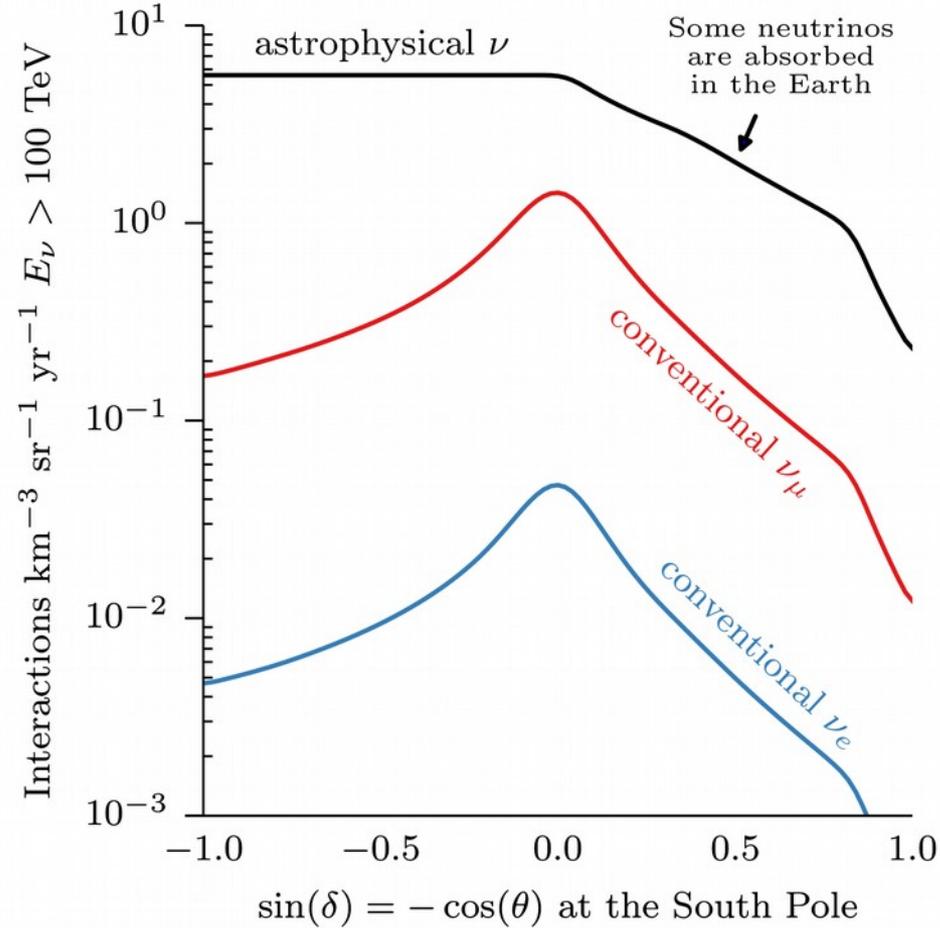
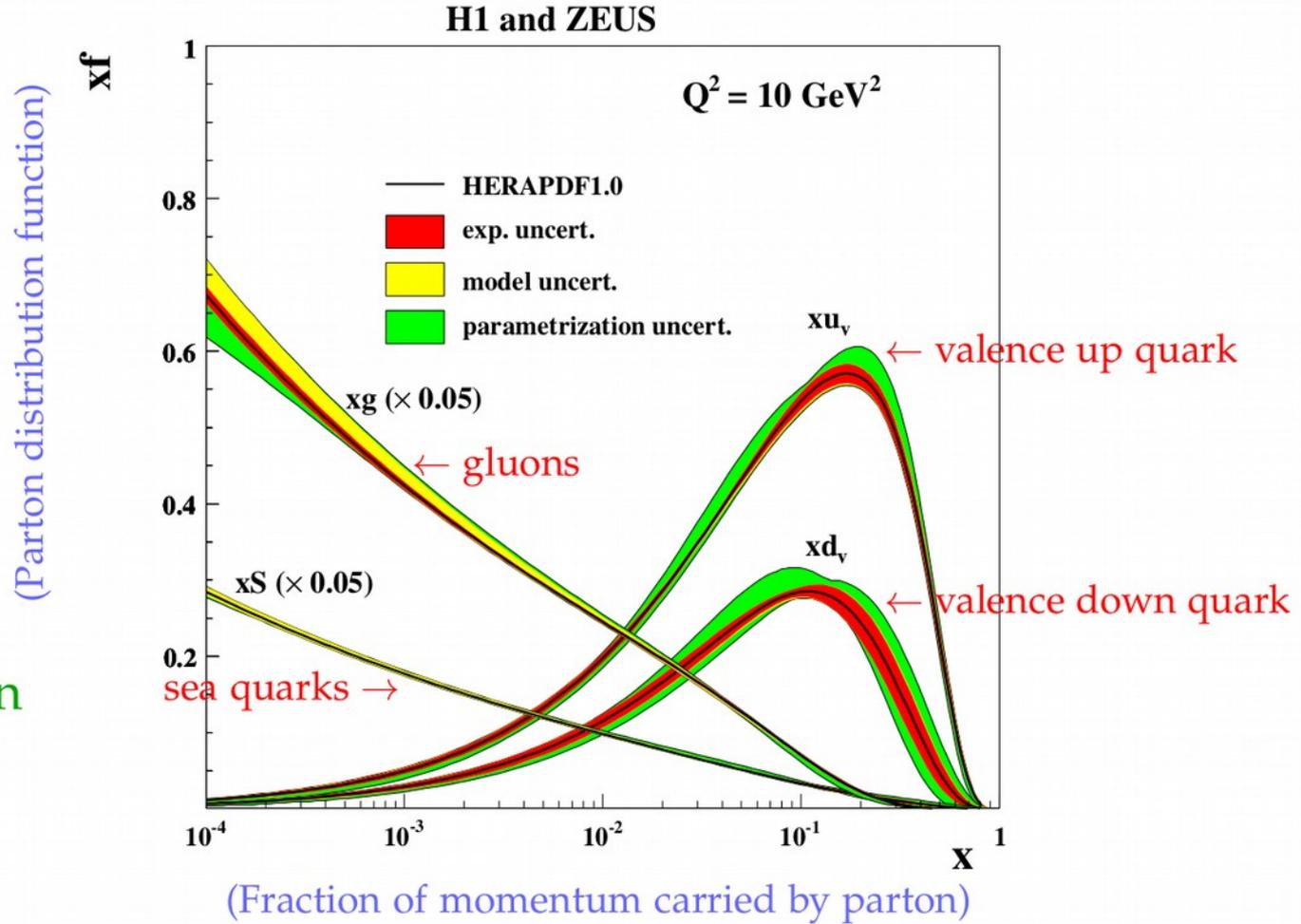


Figure by
Jakob Van Santen
ICRC 2017

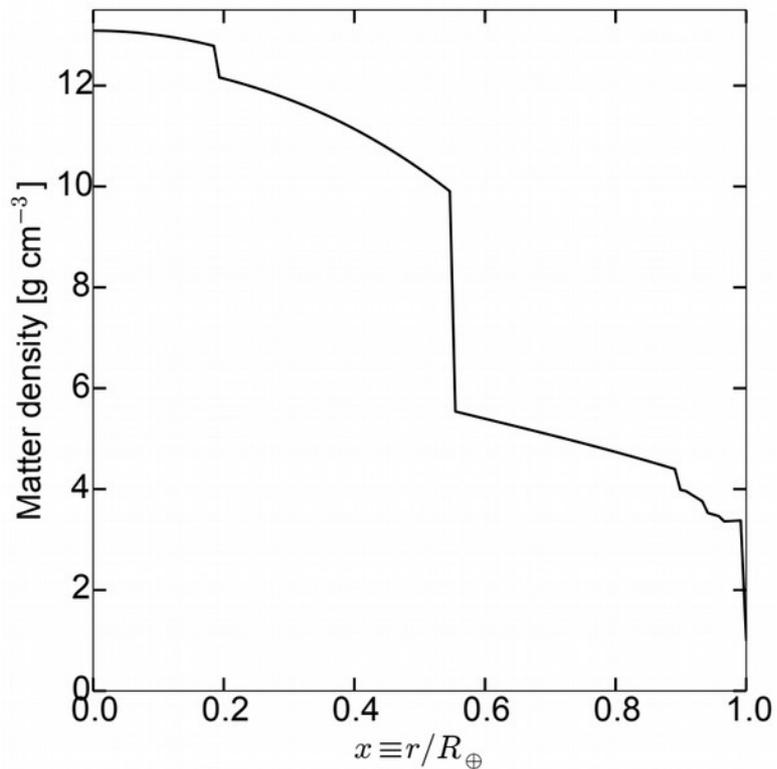
Peeking inside a proton



A feel for the in-Earth attenuation

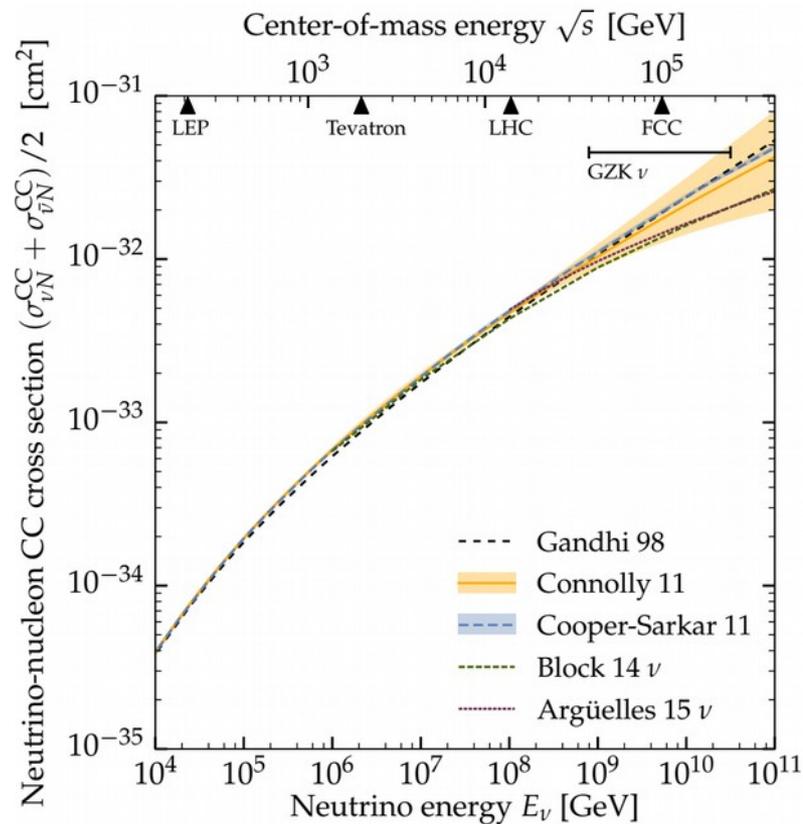
Earth matter density

(Preliminary Reference Earth Model)

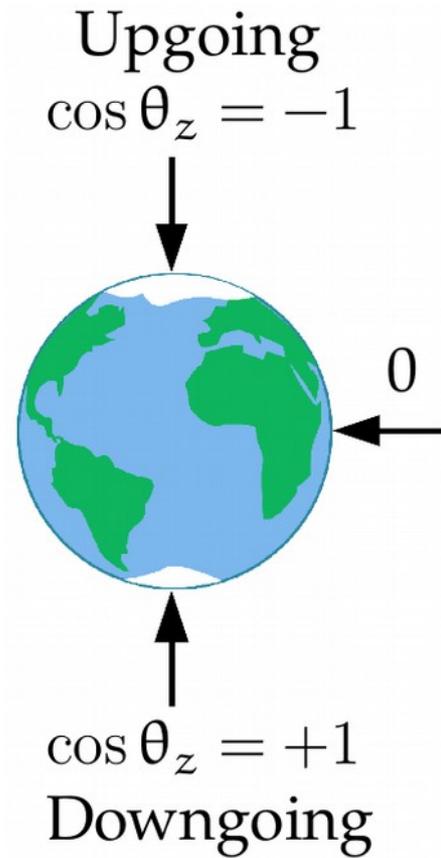
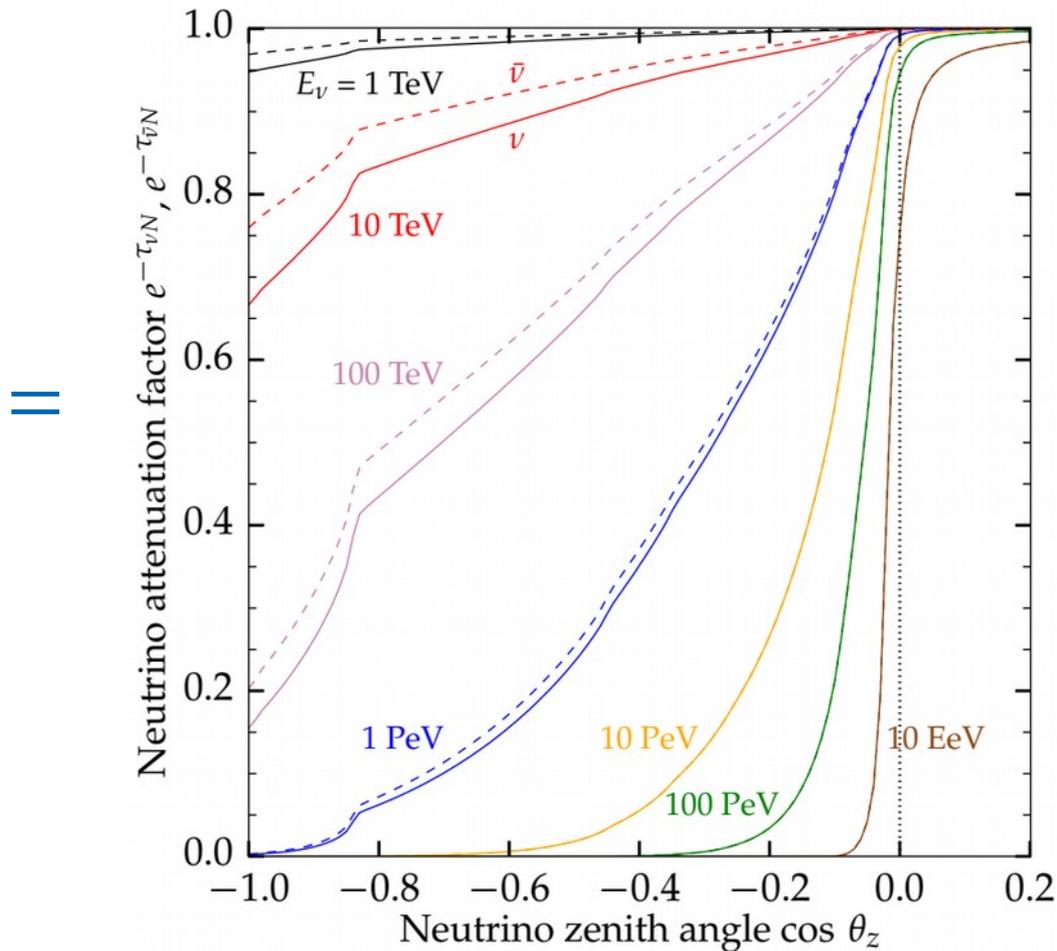


+

Neutrino-nucleon cross section



A feel for the in-Earth attenuation



What goes into the (likelihood) mix?

- ▶ Inside each energy bin, we freely vary
 - ▶ N_{ast} (showers from astrophysical neutrinos)
 - ▶ N_{atm} (showers from atmospheric neutrinos)
 - ▶ γ (astrophysical spectral index)
 - ▶ σ_{CC} (neutrino-nucleon charged-current cross section)
- ▶ For each combination, we generate the angular and energy shower spectrum...
- ▶ ... and compare it to the observed HESE spectrum via a likelihood
- ▶ Maximum likelihood yields σ_{CC} (marginalized over nuisance parameters)
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Including detector resolution
(10% in energy, 15° in direction)

Marginalized cross section in each bin

TABLE I. Neutrino-nucleon charged-current inclusive cross sections, averaged between neutrinos ($\sigma_{\nu N}^{\text{CC}}$) and anti-neutrinos ($\sigma_{\bar{\nu} N}^{\text{CC}}$), extracted from 6 years of IceCube HESE showers. To obtain these results, we fixed $\sigma_{\bar{\nu} N}^{\text{CC}} = \langle \sigma_{\bar{\nu} N}^{\text{CC}} / \sigma_{\nu N}^{\text{CC}} \rangle \cdot \sigma_{\nu N}^{\text{CC}}$ — where $\langle \sigma_{\bar{\nu} N}^{\text{CC}} / \sigma_{\nu N}^{\text{CC}} \rangle$ is the average ratio of $\bar{\nu}$ to ν cross sections calculated using the standard prediction from Ref. [60](#) — and $\sigma_{\nu N}^{\text{NC}} = \sigma_{\nu N}^{\text{CC}}/3$, $\sigma_{\bar{\nu} N}^{\text{NC}} = \sigma_{\bar{\nu} N}^{\text{CC}}/3$. Uncertainties are statistical plus systematic, added in quadrature.

E_ν [TeV]	$\langle E_\nu \rangle$ [TeV]	$\langle \sigma_{\bar{\nu} N}^{\text{CC}} / \sigma_{\nu N}^{\text{CC}} \rangle$	$\log_{10}[\frac{1}{2}(\sigma_{\nu N}^{\text{CC}} + \sigma_{\bar{\nu} N}^{\text{CC}})/\text{cm}^2]$
18–50	32	0.752	-34.35 ± 0.53
50–100	75	0.825	-33.80 ± 0.67
100–400	250	0.888	-33.84 ± 0.67
400–2004	1202	0.957	$> -33.21 (1\sigma)$

Energy and angular shower spectra

Rate from all flavors, CC + NC:

$$\frac{d^2 N_{\text{sh}}}{dE_{\text{sh}} d \cos \theta_z} = \frac{d^2 N_{\text{sh},e}^{\text{CC}}}{dE_{\text{sh}} d \cos \theta_z} + \text{Br}_{\tau \rightarrow \text{sh}} \frac{d^2 N_{\text{sh},\tau}^{\text{CC}}}{dE_{\text{sh}} d \cos \theta_z} + \sum_{l=e,\mu,\tau} \frac{d^2 N_{\text{sh},l}^{\text{NC}}}{dE_{\text{sh}} d \cos \theta_z}$$

$\text{Br}_{\tau \rightarrow \text{sh}} = 0.83$

Contribution from one flavor CC:

$$\frac{d^2 N_{\text{sh},l}^{\text{CC}}}{dE_{\text{sh}} d \cos \theta_z}(E_{\text{sh}}, \cos \theta_z) \simeq -2\pi\rho_{\text{ice}} N_A V T \left\{ \Phi_l(E_\nu) \sigma_{\nu N}^{\text{CC}}(E_\nu) e^{-\tau_\nu N(E_\nu, \theta_z)} + \Phi_{\bar{l}}(E_\nu) \sigma_{\bar{\nu} N}^{\text{CC}}(E_\nu) e^{-\tau_{\bar{\nu}} N(E_\nu, \theta_z)} \right\} \Big|_{E_\nu = E_{\text{sh}}/f_{l,\text{CC}}}$$

Conversion between shower energy and neutrino energy:

$$f_{l,t} \equiv \frac{E_{\text{sh}}}{E_\nu} \simeq \begin{cases} 1 & \text{for } l = e \text{ and } t = \text{CC} \\ [\langle y \rangle + 0.7(1 - \langle y \rangle)] \simeq 0.8 & \text{for } l = \tau \text{ and } t = \text{CC} \\ \langle y \rangle \simeq 0.25 & \text{for } l = e, \mu, \tau \text{ and } t = \text{NC} \end{cases}$$

Detector resolution

Number of contained showers:

$$\frac{d^2 N_{\text{sh}}}{dE_{\text{dep}} d \cos \theta_z} = \int dE_{\text{sh}} \int d \cos \theta'_z \frac{d^2 N_{\text{sh}}}{dE_{\text{sh}} d \cos \theta'_z} R_E(E_{\text{sh}}, E_{\text{dep}}, \sigma_E(E_{\text{sh}})) R_\theta(\cos \theta'_z, \cos \theta_z, \sigma_{\cos \theta_z})$$

Energy resolution: [Palomares-Ruiz, Vincent, Mena *PRD* 2015; Vincent, Palomares-Ruiz, Mena *PRD* 2016; MB, Beacom, Murase, *PRD* 2016]

$$R_E(E_{\text{sh}}, E_{\text{dep}}, \sigma_E(E_{\text{sh}})) = \frac{1}{\sqrt{2\pi\sigma_E^2(E_{\text{sh}})}} \exp \left[-\frac{(E_{\text{sh}} - E_{\text{dep}})^2}{2\sigma_E^2(E_{\text{sh}})} \right] \quad \text{with } \sigma_E(E_{\text{sh}}) = 0.1E_{\text{sh}}$$

IceCube, JINST 2014

Angular resolution:

$$R_\theta(\cos \theta'_z, \cos \theta_z, \sigma_{\cos \theta_z}) = \frac{1}{\sqrt{2\pi\sigma_{\cos \theta_z}^2}} \exp \left[-\frac{(\cos \theta'_z - \cos \theta_z)^2}{2\sigma_{\cos \theta_z}^2} \right]$$

with $\sigma_{\cos \theta_z} \equiv \frac{1}{2} [|\cos(\theta_z + \sigma_{\theta_z}) - \cos \theta_z| + |\cos(\theta_z - \sigma_{\theta_z}) - \cos \theta_z|]$ and $\sigma_{\theta_z} = 15^\circ$

Likelihood

In an energy bin containing $N_{\text{sh}}^{\text{obs}}$ observed showers, the likelihood is

Each energy bin is independent

$$\mathcal{L} = \frac{e^{-(N_{\text{sh}}^{\text{atm}} + N_{\text{sh}}^{\text{ast}})} N_{\text{sh}}^{\text{obs}}}{N_{\text{sh}}^{\text{obs}}!} \prod_{i=1}^{N_{\text{sh}}^{\text{obs}}} \mathcal{L}_i$$

Partial likelihood, *i.e.*, relative probability of the i -th shower being from an atmospheric neutrino or an astrophysical neutrino:

Depends on $\sigma_{\nu N}$

$$\mathcal{L}_i = N_{\text{sh}}^{\text{atm}} \mathcal{P}_i^{\text{atm}} + N_{\text{sh}}^{\text{ast}} \mathcal{P}_i^{\text{ast}}$$

$$\mathcal{P}_i^{\text{atm}} = \left(\int_{E_{\text{dep}}^{\text{min}}}^{E_{\text{dep}}^{\text{max}}} dE_{\text{dep}} \int_{-1}^1 d \cos \theta_z \frac{d^2 N_{\text{sh}}^{\text{atm}}}{dE_{\text{dep}} d \cos \theta_z} \right)^{-1} \left(\frac{d^2 N_{\text{sh}}^{\text{atm}}}{dE_{\text{dep}} d \cos \theta_z} \Big|_{E_{\text{dep},i}, \cos \theta_{z,i}} \right)$$

PDF for this shower to be made by an atmospheric ν

$$\mathcal{P}_i^{\text{ast}} = \left(\int_{E_{\text{dep}}^{\text{min}}}^{E_{\text{dep}}^{\text{max}}} dE_{\text{dep}} \int_{-1}^1 d \cos \theta_z \frac{d^2 N_{\text{sh}}^{\text{ast}}}{dE_{\text{dep}} d \cos \theta_z} \right)^{-1} \left(\frac{d^2 N_{\text{sh}}^{\text{ast}}}{dE_{\text{dep}} d \cos \theta_z} \Big|_{E_{\text{dep},i}, \cos \theta_{z,i}} \right)$$

PDF for this shower to be made by an astrophysical ν

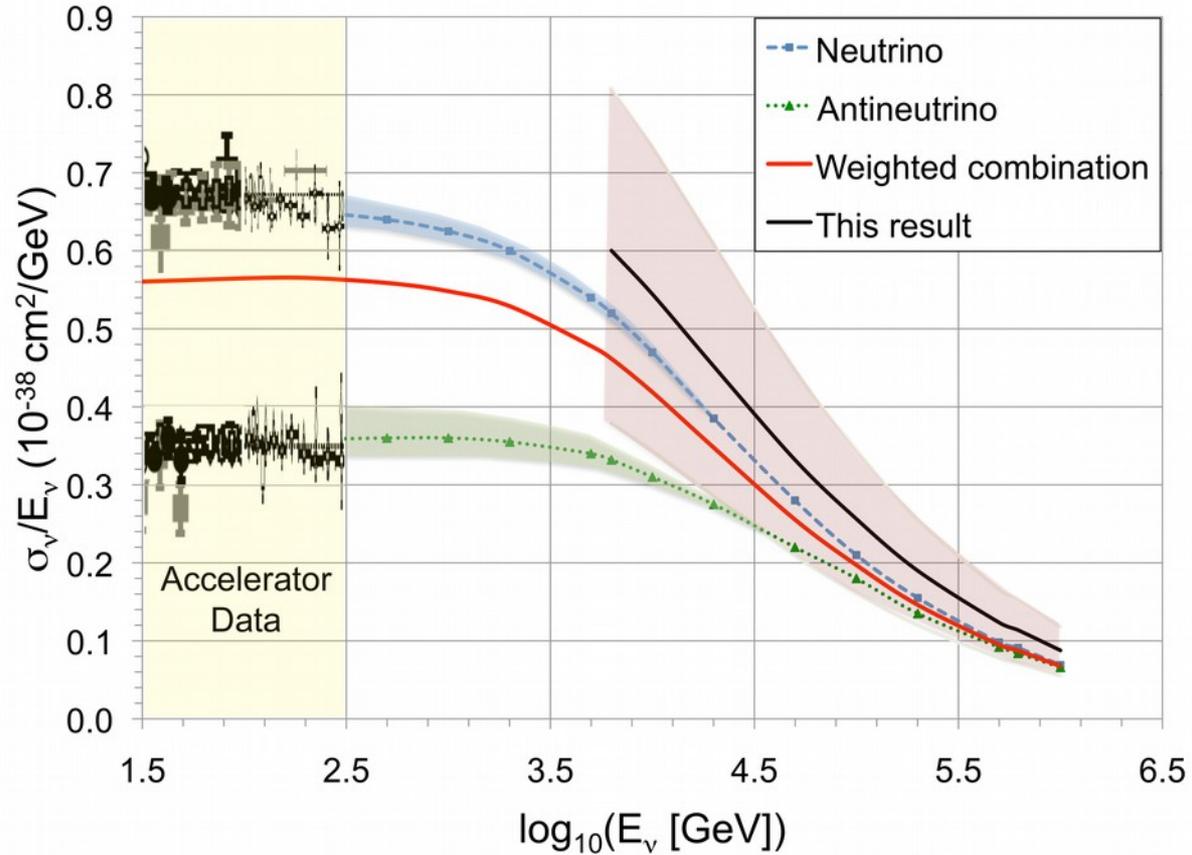
Depends on γ and $\sigma_{\nu N}$

The fine print

- ▶ High-energy ν 's: astrophysical (isotropic) + atmospheric (**anisotropic**)
→ We take into account the shape of the atmospheric contribution
- ▶ The shape of the astrophysical ν **energy spectrum** is still uncertain
→ We take a $E^{-\gamma}$ spectrum in *narrow* energy bins
- ▶ **NC showers** are sub-dominant to **CC showers**, but they are indistinguishable
→ Following Standard-Model predictions, we take $\sigma_{\text{NC}} = \sigma_{\text{CC}}/3$
- ▶ IceCube does not **distinguish ν from $\bar{\nu}$** , and their cross-sections are different
→ We assume equal fluxes, expected from production via pp collisions
→ We assume the avg. ratio $\langle \sigma_{\bar{\nu}N} / \sigma_{\nu N} \rangle$ in each bin known, from SM predictions
- ▶ The **flavor composition** of astrophysical neutrinos is still uncertain
→ We assume equal flux of each flavor, compatible with theory and observations

Using through-going muons instead

- ▶ Use $\sim 10^4$ through-going muons
- ▶ Measured: dE_μ/dx
- ▶ Inferred: $E_\mu \approx dE_\mu/dx$
- ▶ From simulations (uncertain):
most likely E_ν given E_μ
- ▶ Fit the ratio $\sigma_{\text{obs}}/\sigma_{\text{SM}}$
 $1.30^{+0.21}_{-0.19}(\text{stat.})^{+0.39}_{-0.43}(\text{syst.})$
- ▶ All events grouped in a single
energy bin 6–980 TeV



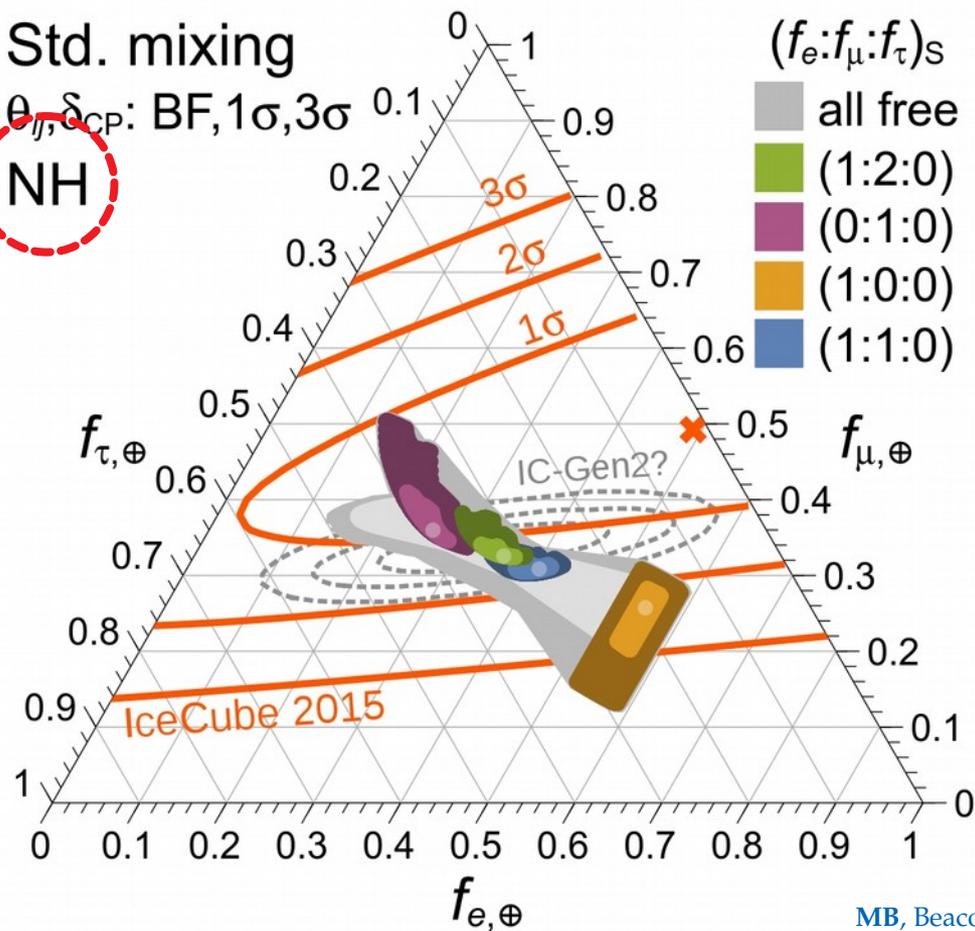
Flavor composition – a few source choices

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Std. mixing

θ_{ij}, δ_{CP} : BF, $1\sigma, 3\sigma$

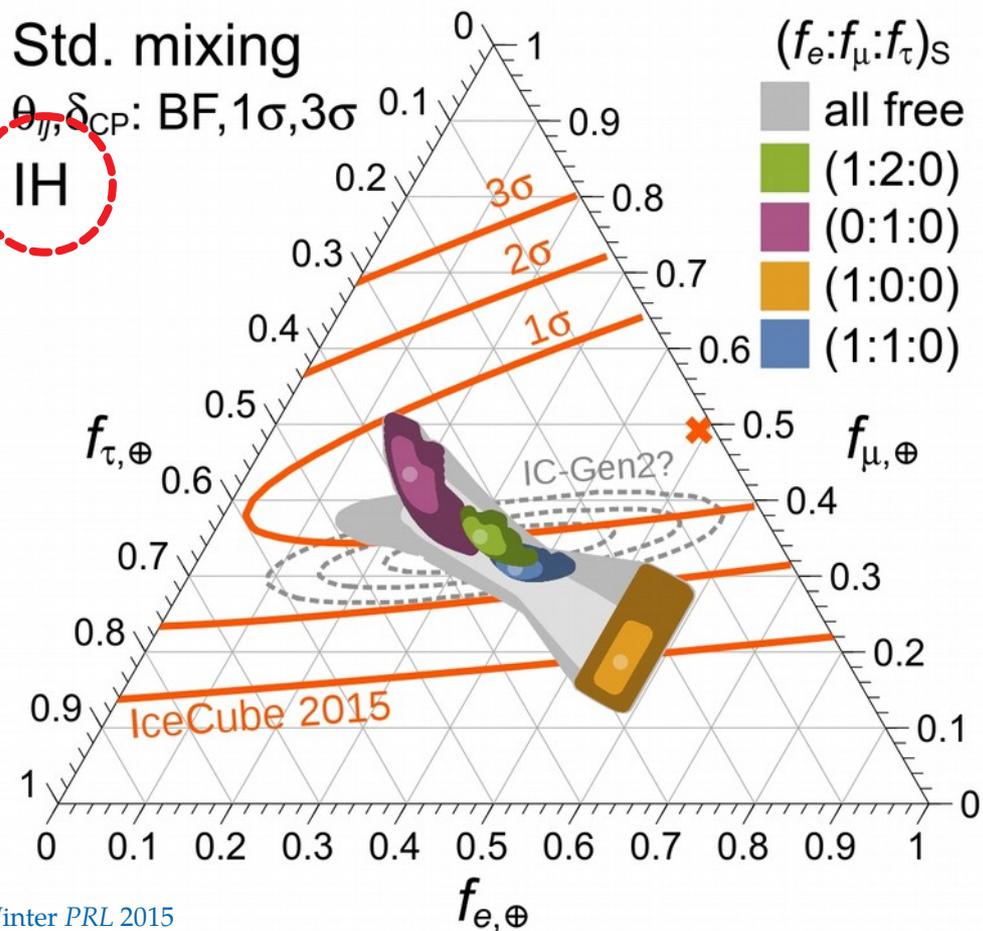
NH



Std. mixing

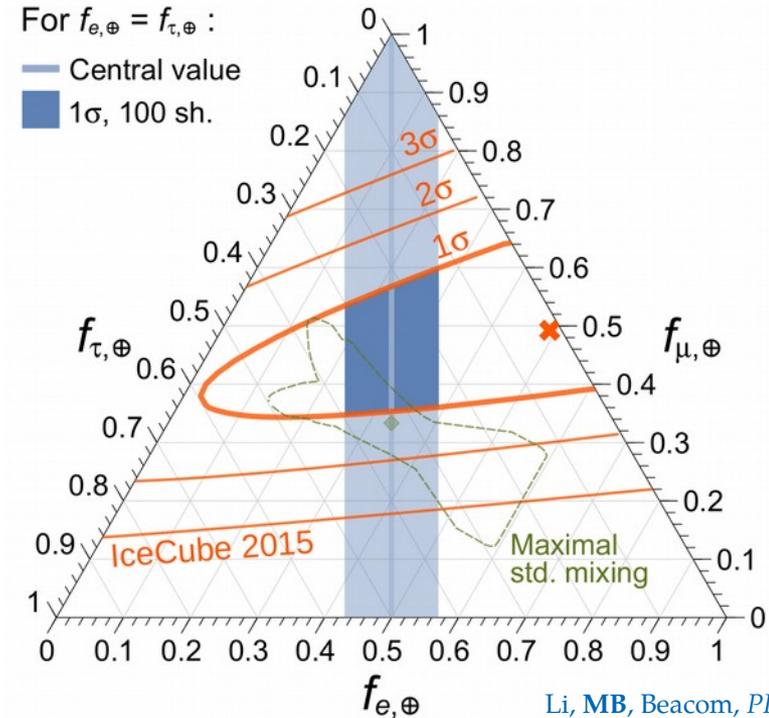
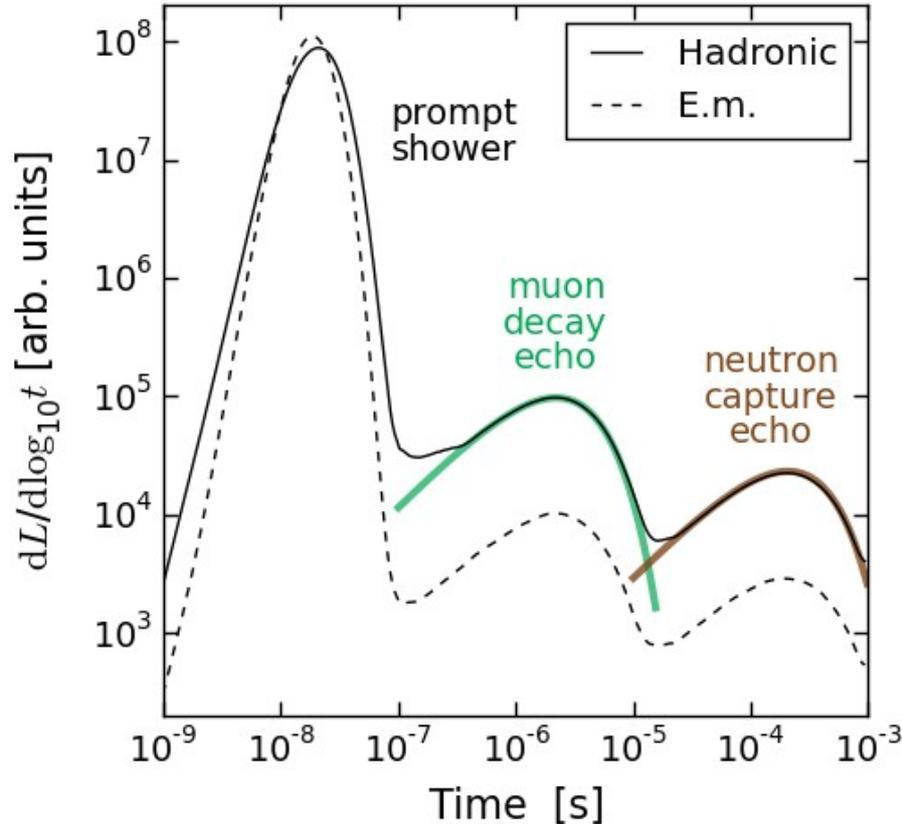
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IH



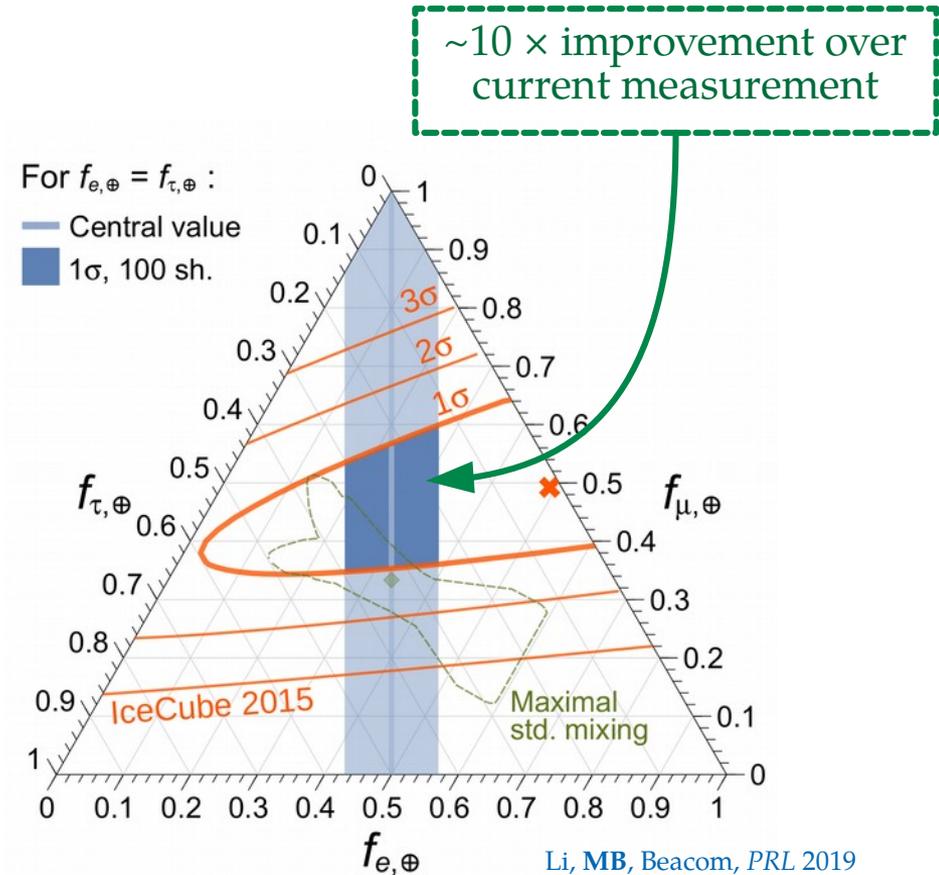
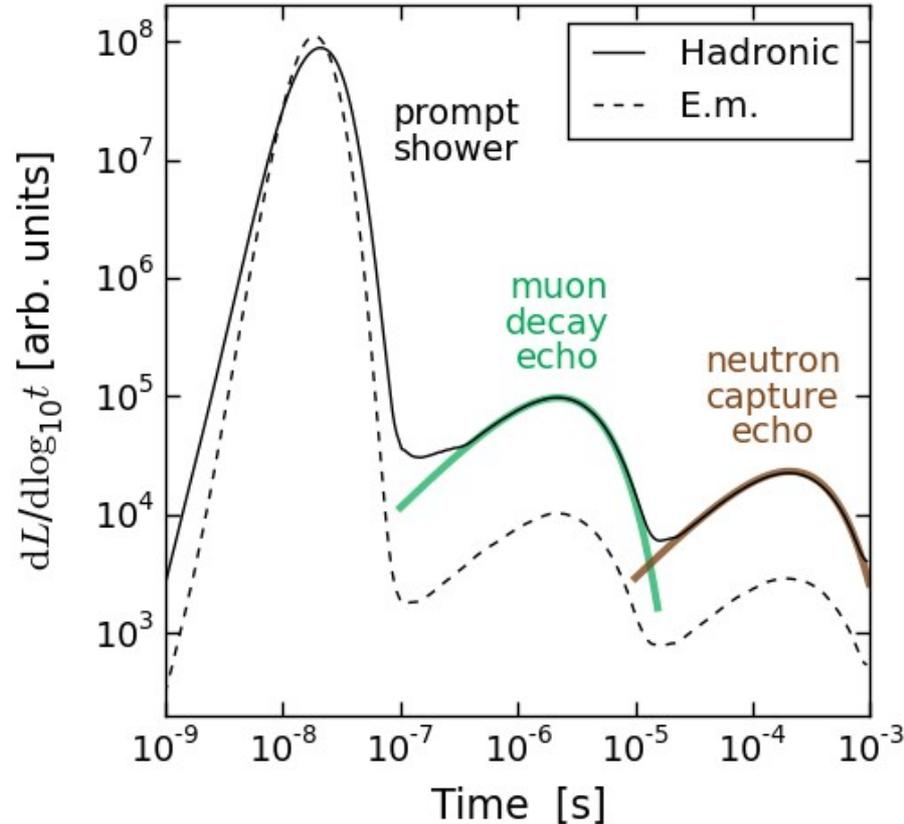
Side note: Improving flavor-tagging using *echoes*

Late-time light (*echoes*) from muon decays and neutron captures can separate showers made by ν_e and ν_τ –



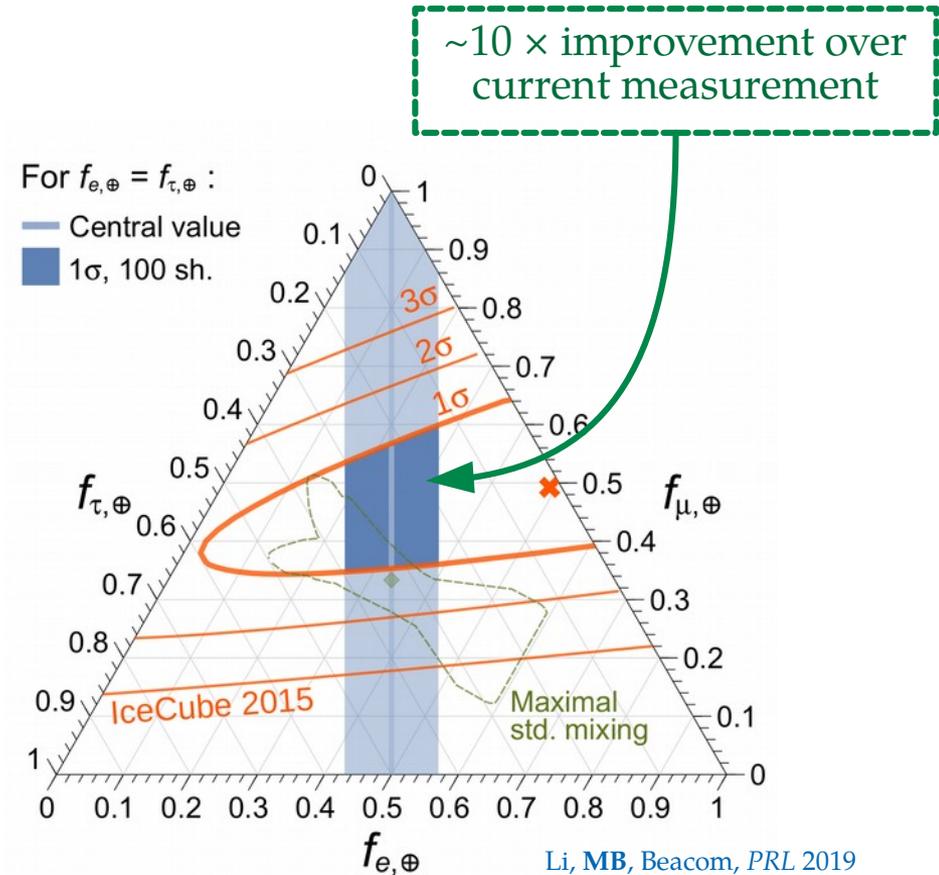
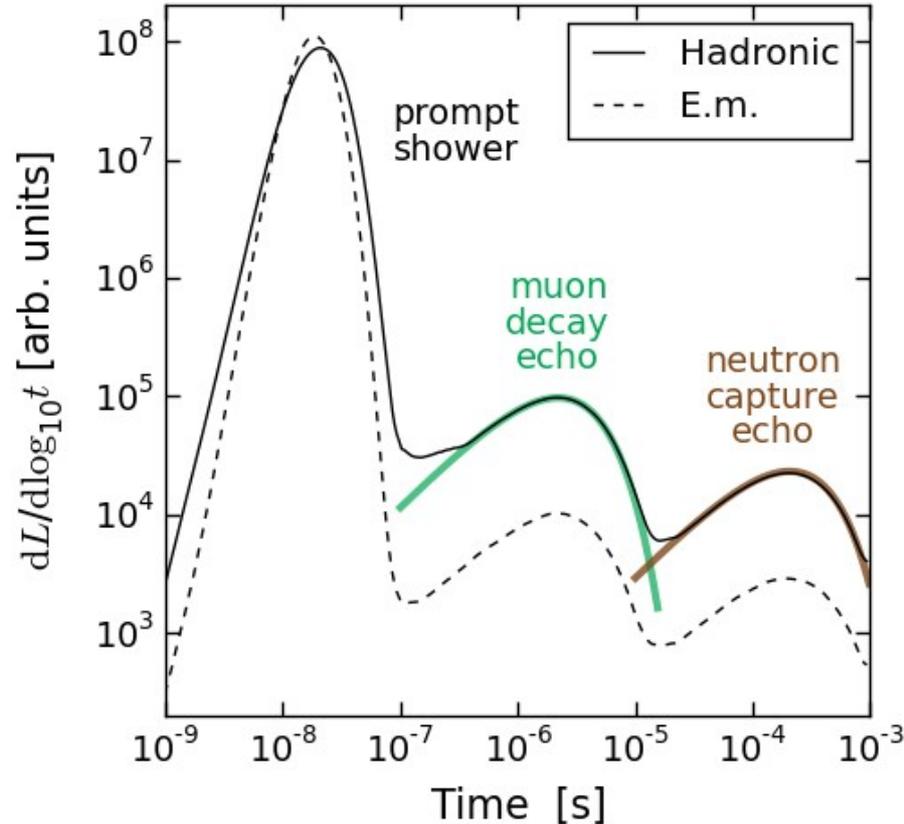
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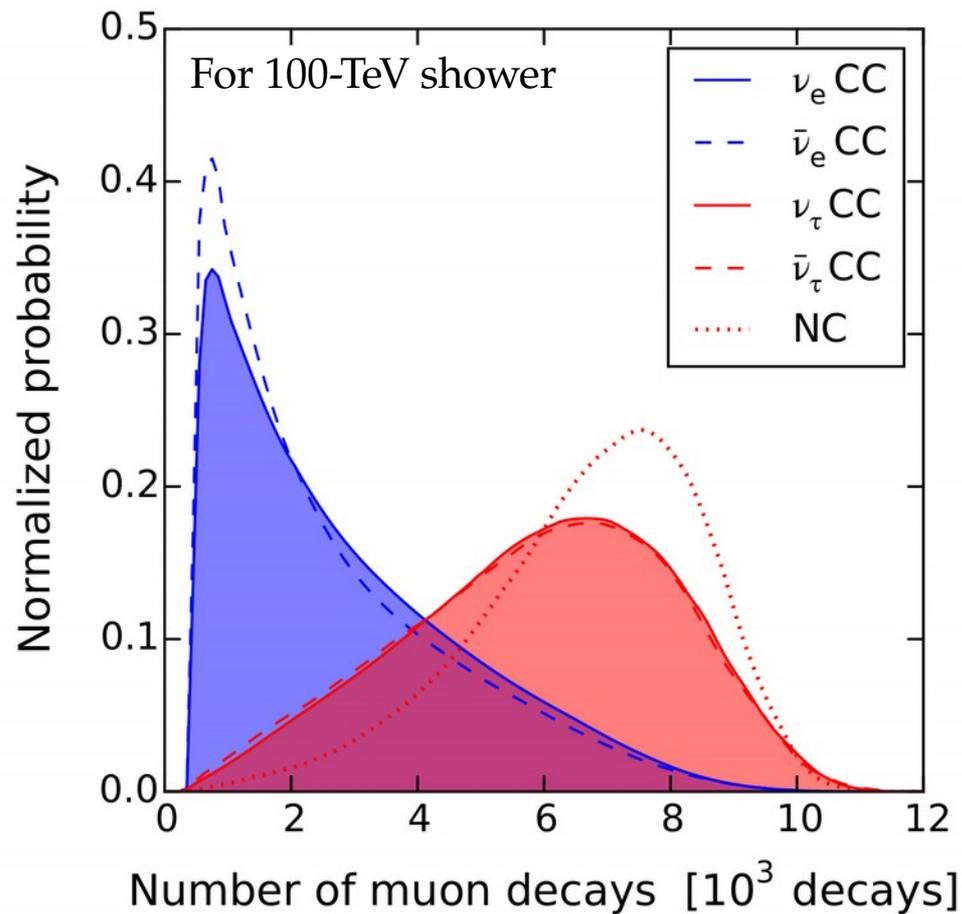
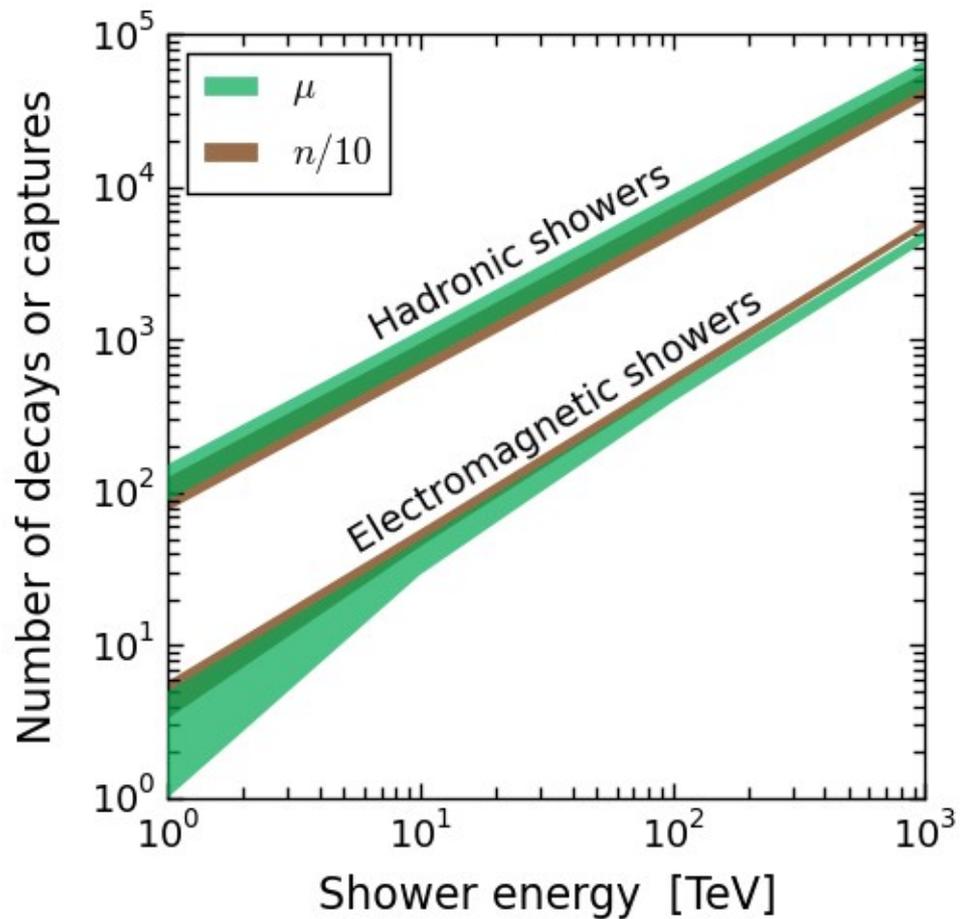


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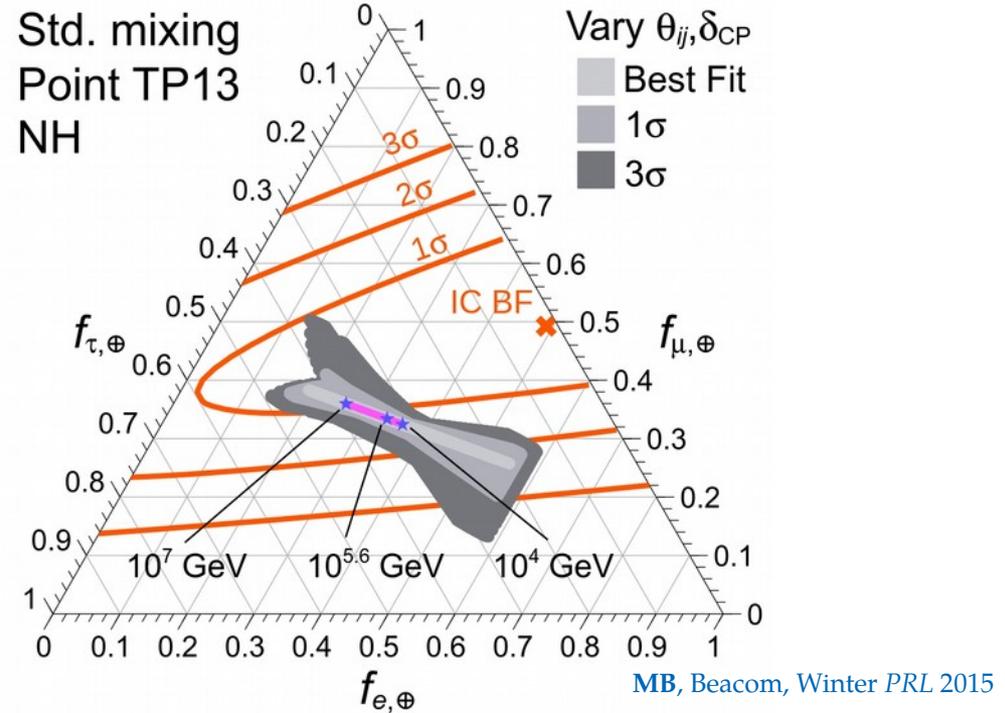
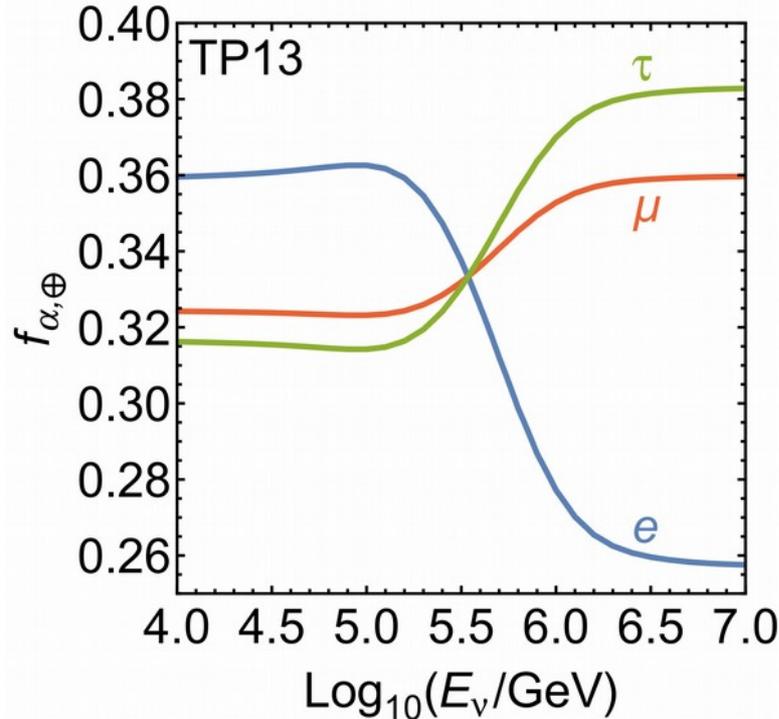


Hadronic *vs.* electromagnetic showers



Energy dependence of the flavor composition?

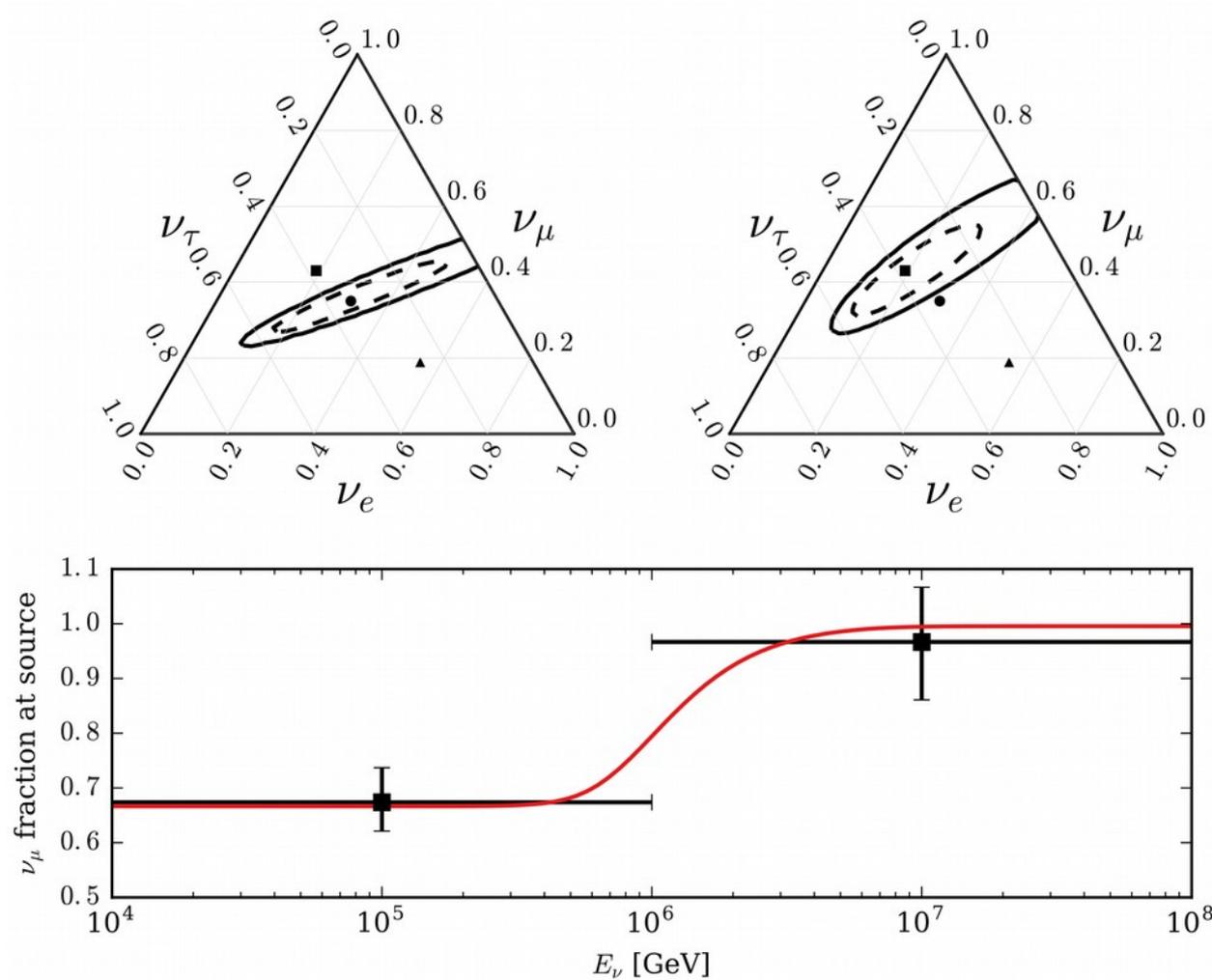
Different neutrino production channels accessible at different energies –



► TP13: $p\gamma$ model, target photons from electron-positron annihilation [Hümmer+, *Astropart. Phys.* 2010]

► Will be difficult to resolve [Kashti, Waxman, *PRL* 2005; Lipari, Lusignoli, Meloni, *PRD* 2007]

... Observable in IceCube-Gen2?



Borrowed from M. Kowalski

Flavor content of neutrino mass eigenstates

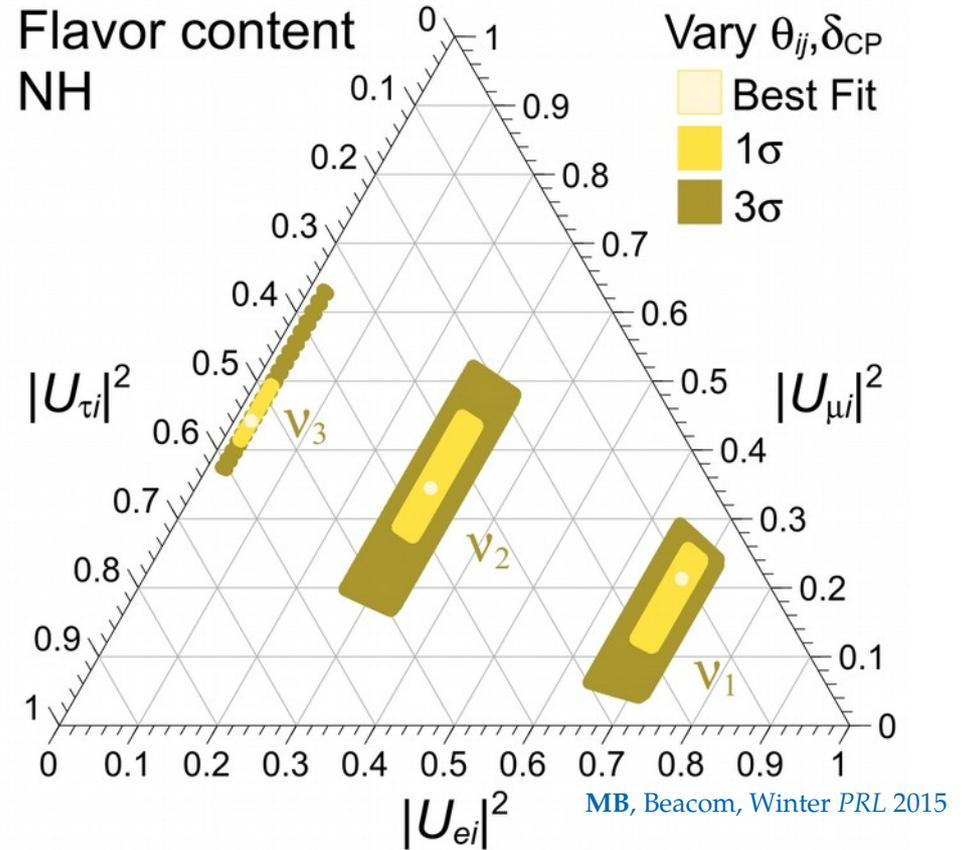
Flavor content for every allowed combination of mixing parameters –

$$|U_{\alpha i}|^2 = |U_{\alpha i}(\theta_{12}, \theta_{23}, \theta_{13}, \delta_{CP})|^2$$

Known to within 2%

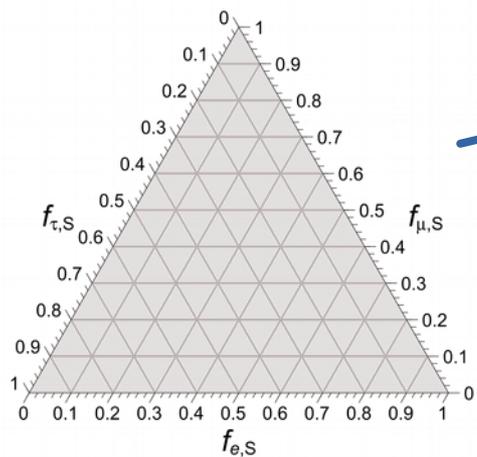
Known to within 8%

Known to within 20% (or worse)



Measuring the neutrino lifetime

Sources

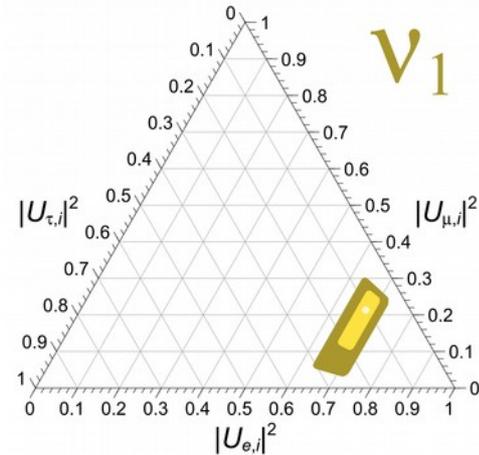


$\nu_2, \nu_3 \rightarrow \nu_1$
 ν_1 lightest and stable

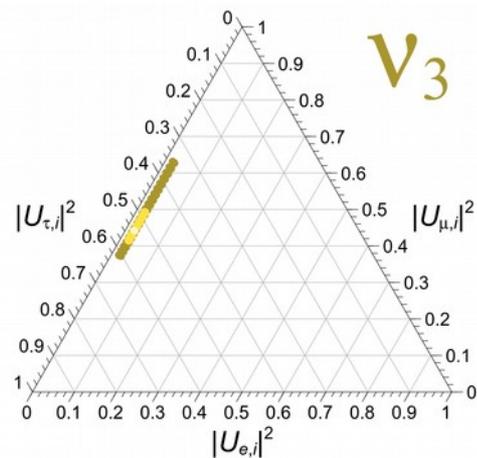
If all unstable
 neutrinos decay

$\nu_1, \nu_2 \rightarrow \nu_3$
 ν_3 lightest and stable

Earth



$$f_{\alpha,\oplus} = |U_{\alpha 1}|^2$$



$$f_{\alpha,\oplus} = |U_{\alpha 3}|^2$$

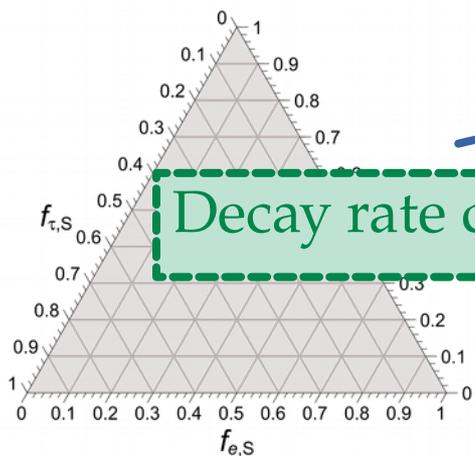
Measuring the neutrino lifetime

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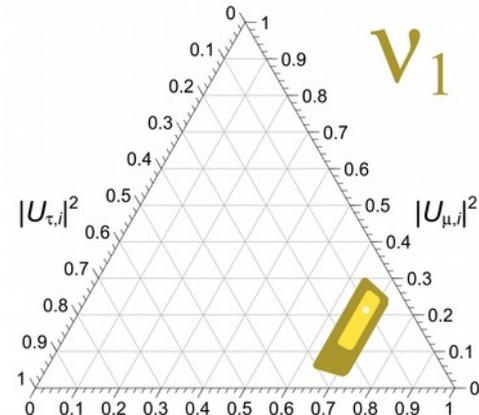
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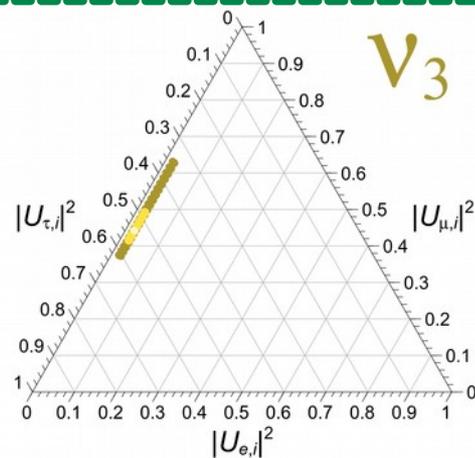
Decay rate depends on $\exp[-t / (\gamma\tau_i)] = \exp[-(L/E) \cdot (m_i/\tau_i)]$

$$\nu_1, \nu_2 \rightarrow \nu_3$$

ν_3 lightest and stable



$$f_{\alpha,\oplus} = |U_{\alpha 1}|^2$$



$$f_{\alpha,\oplus} = |U_{\alpha 3}|^2$$

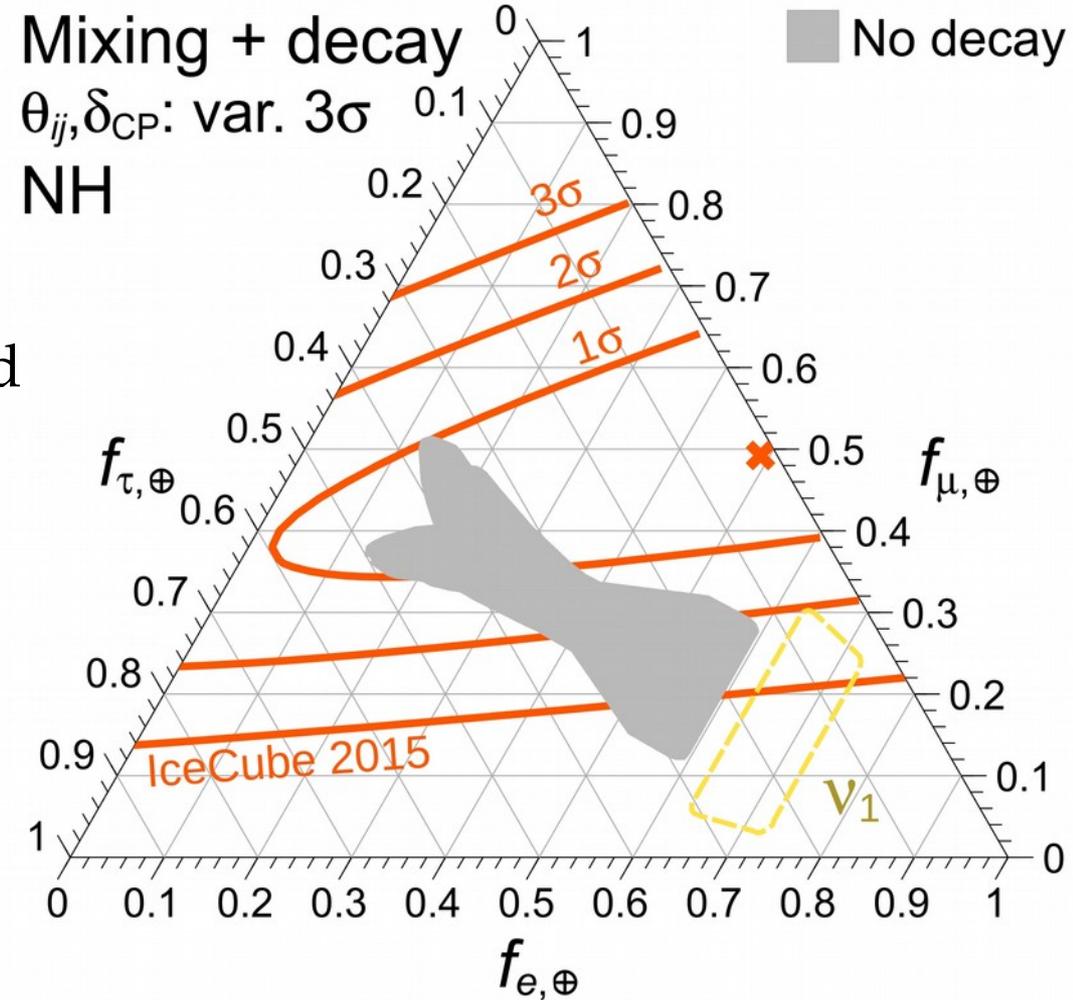
Measuring the neutrino lifetime

Find the value of D so that decay is complete, *i.e.*, $f_{\alpha,\oplus} = |U_{\alpha 1}|^2$, for

- ▶ Any value of mixing parameters; and
- ▶ Any flavor ratios at the sources

(Assume equal lifetimes of ν_2, ν_3)

MB, Beacom, Murase, *PRD* 2017
Baerwald, MB, Winter, *JCAP* 2012



Measuring the neutrino lifetime

Fraction of ν_2, ν_3 remaining at Earth

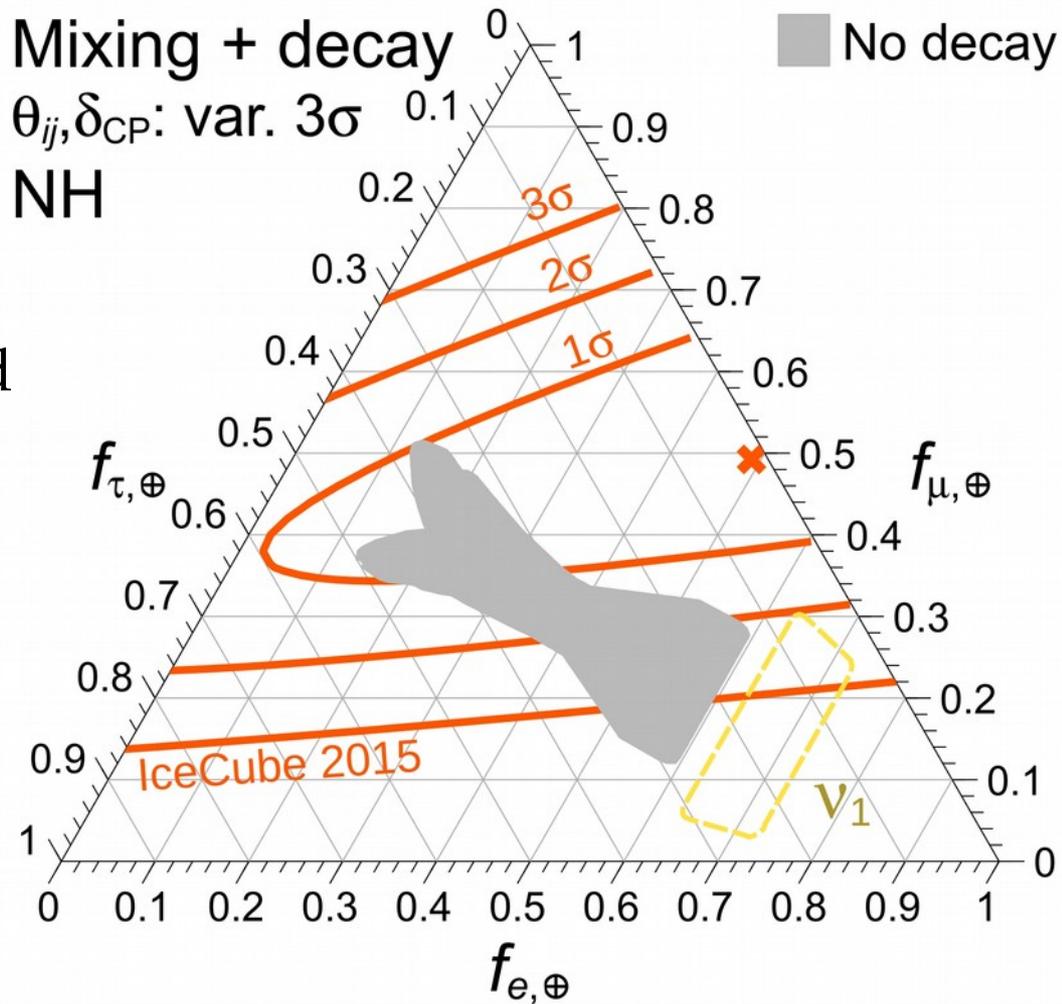


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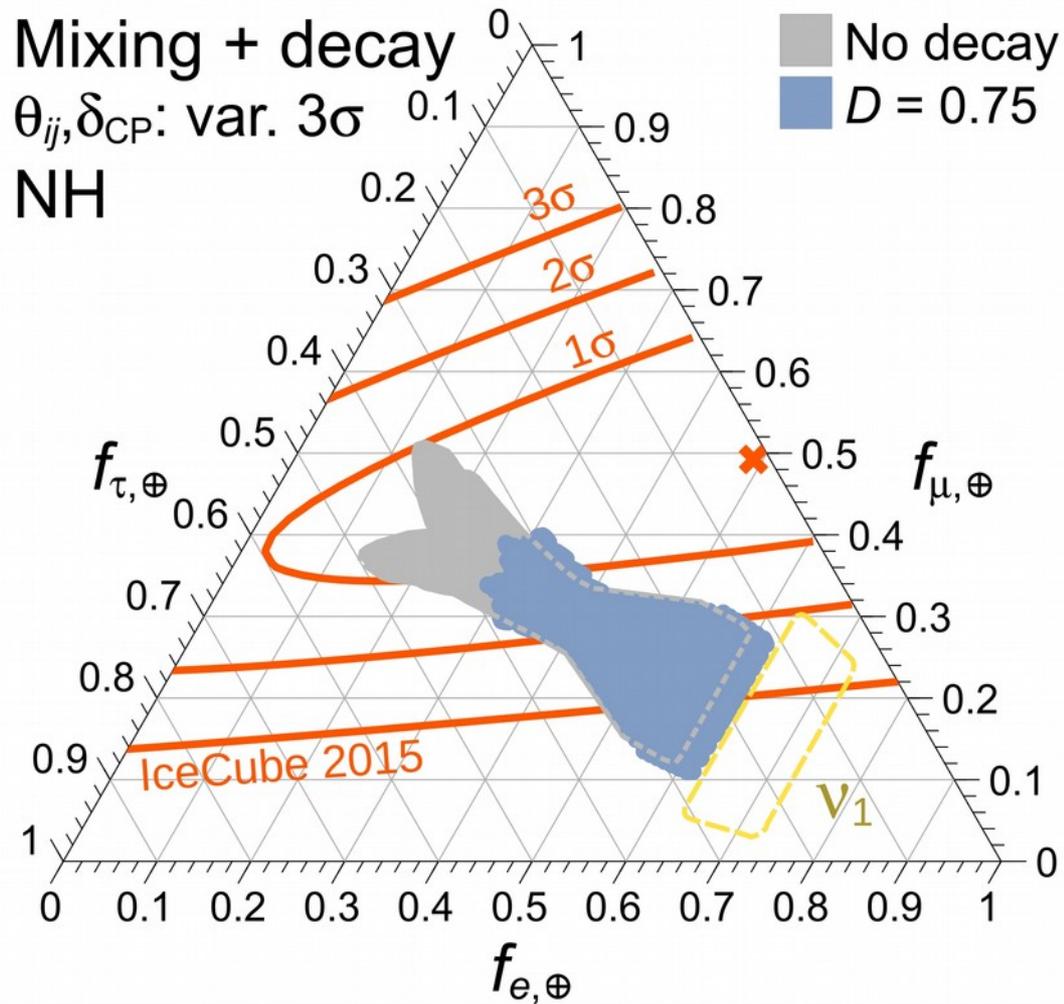


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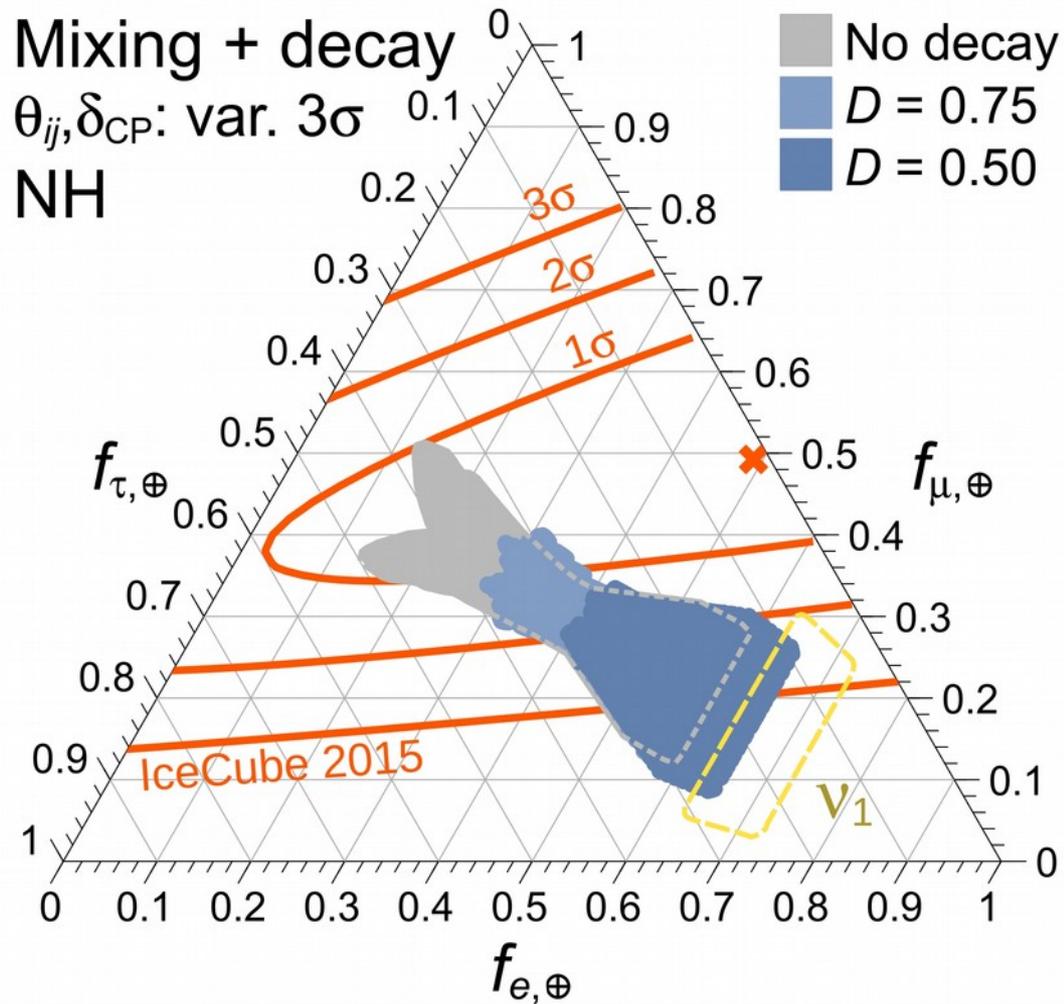


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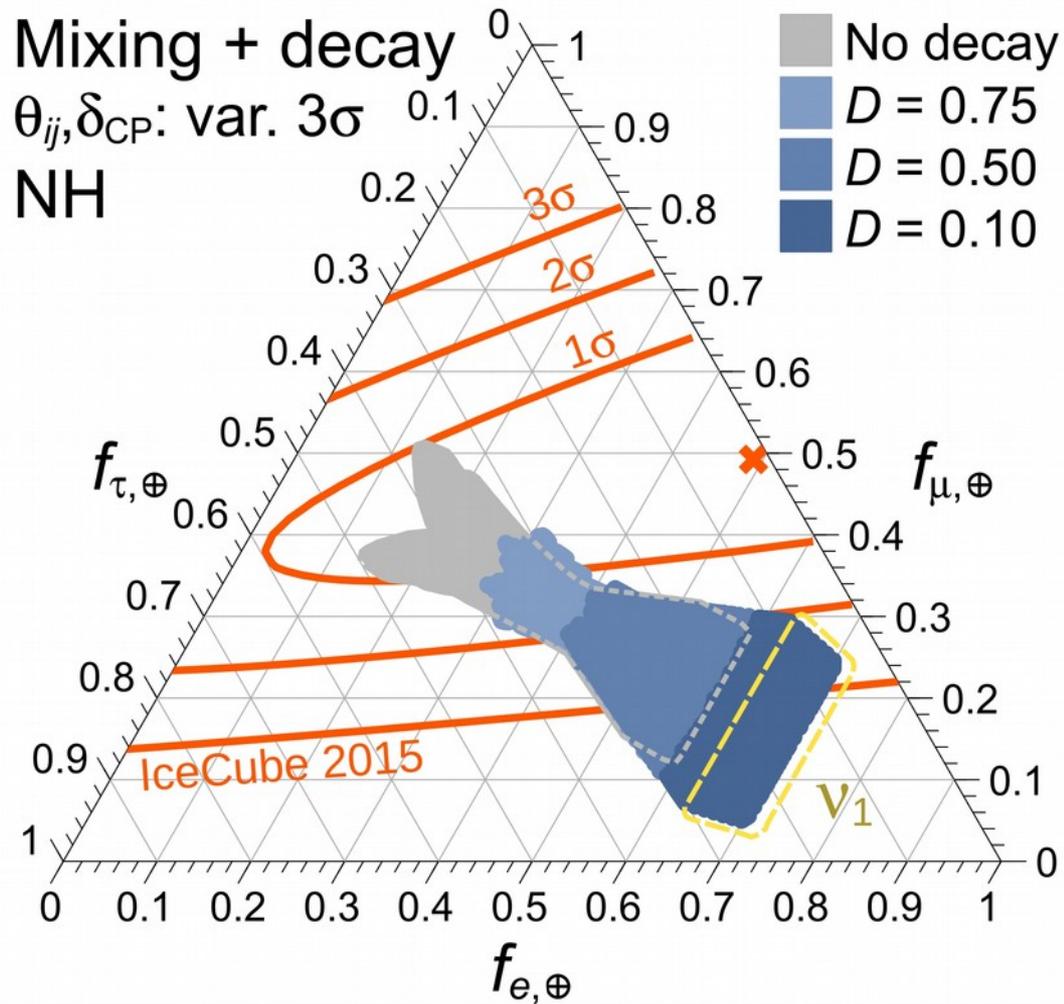


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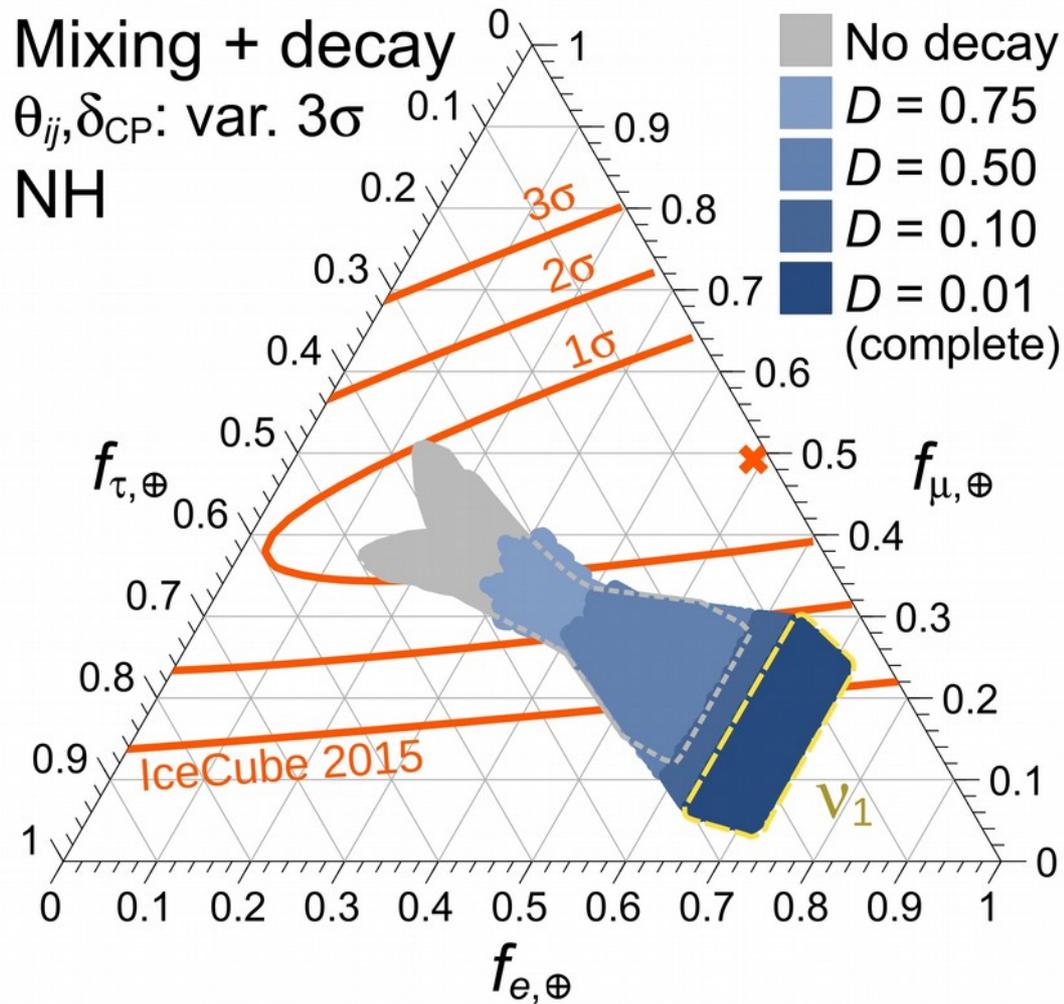


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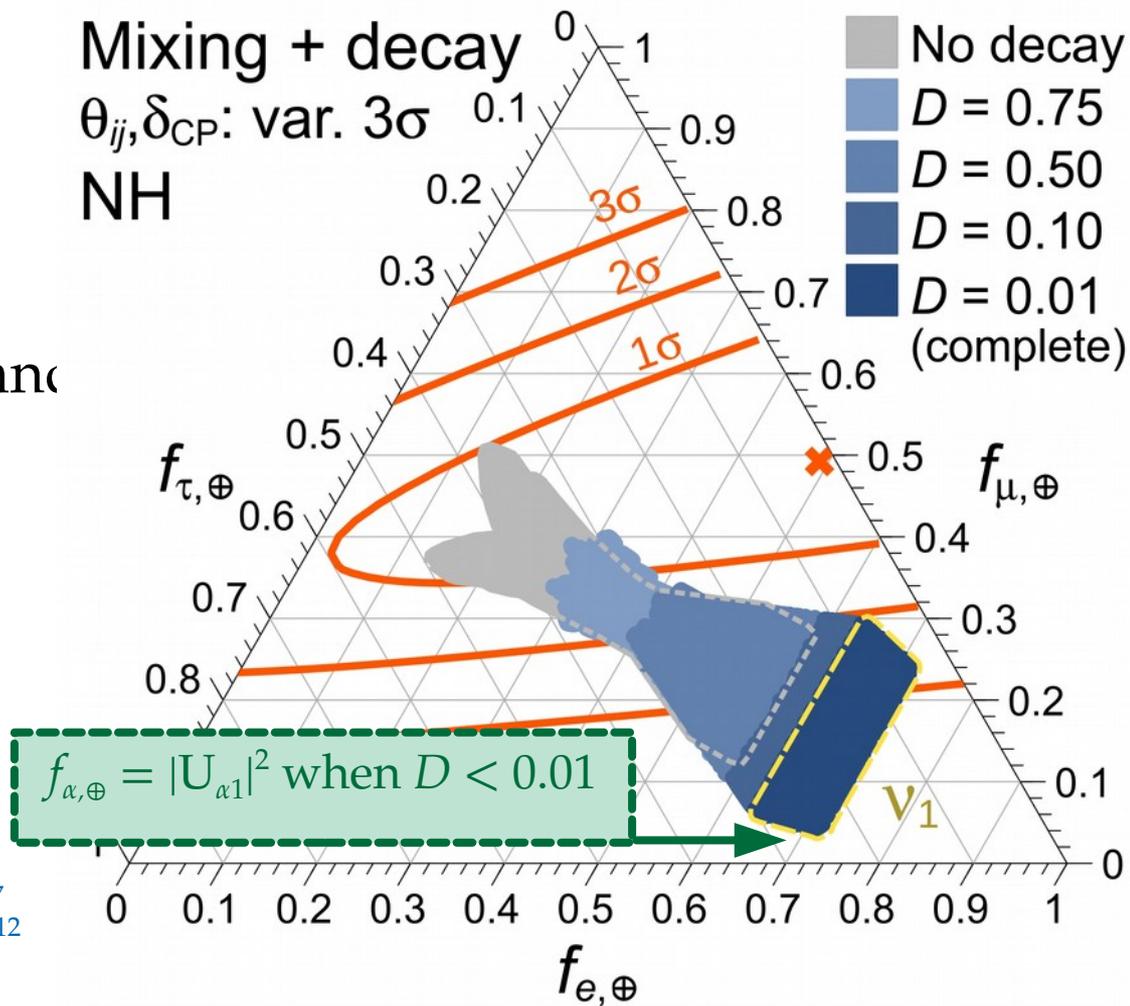
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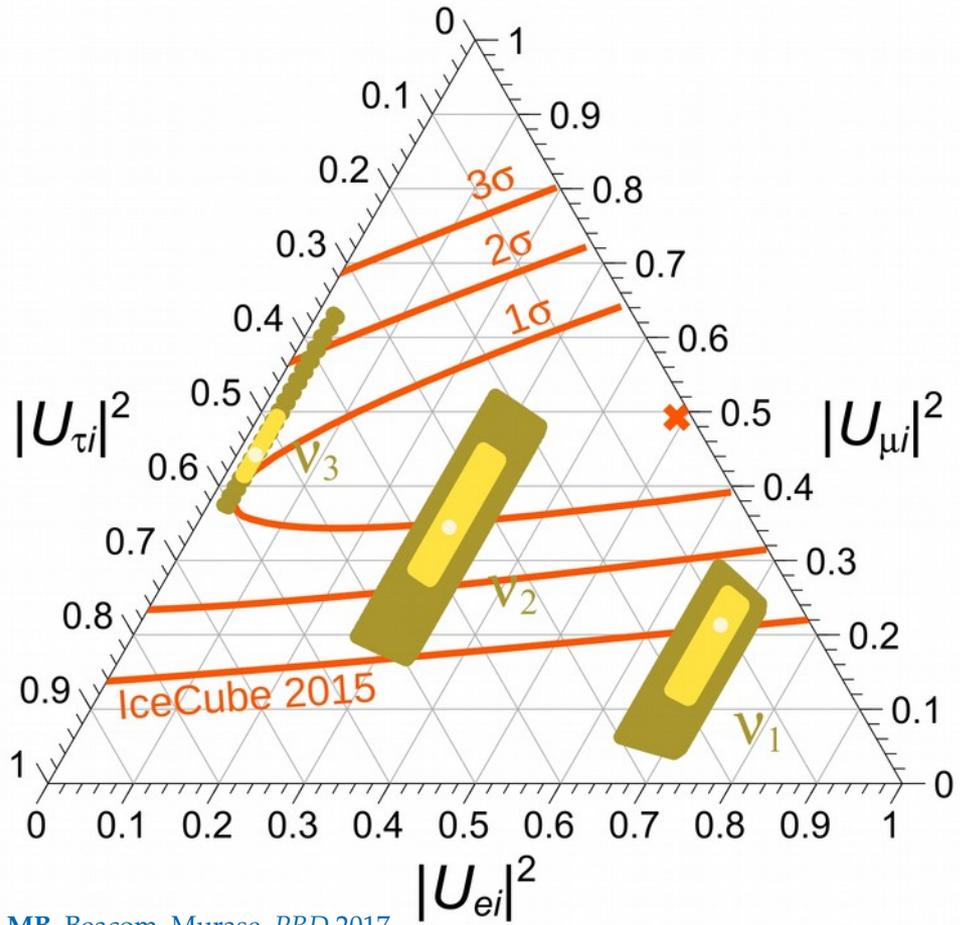


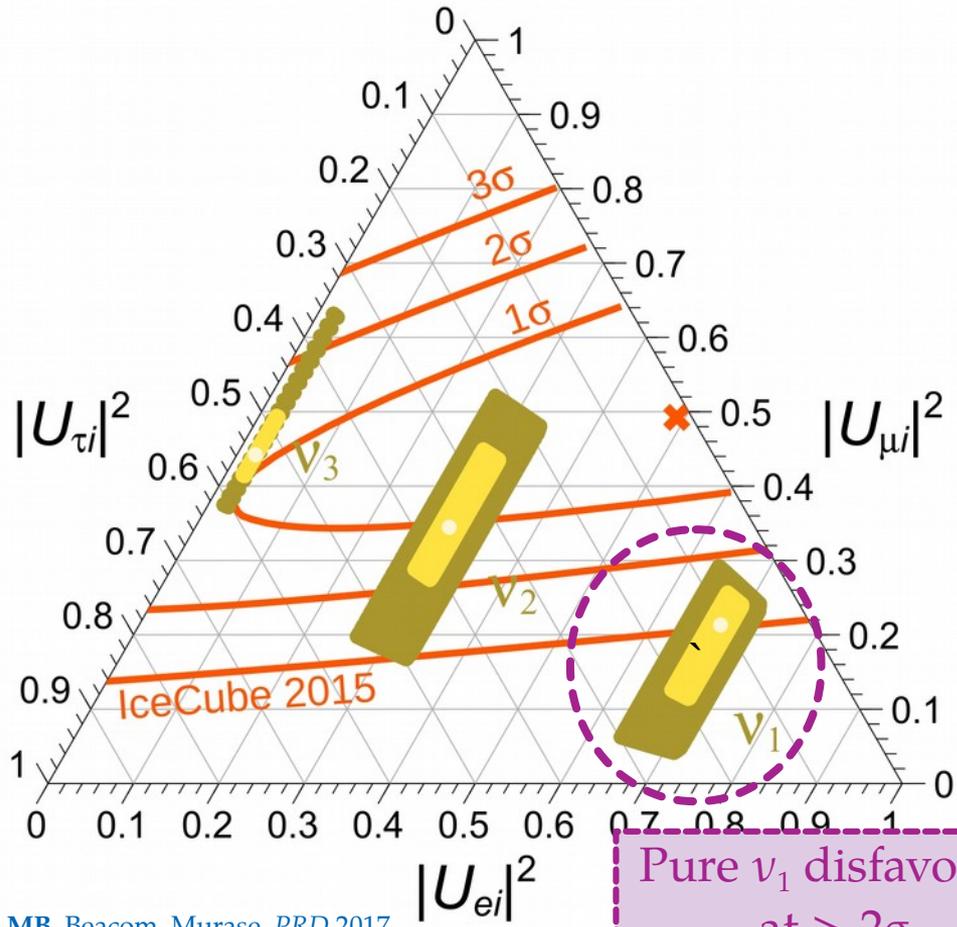
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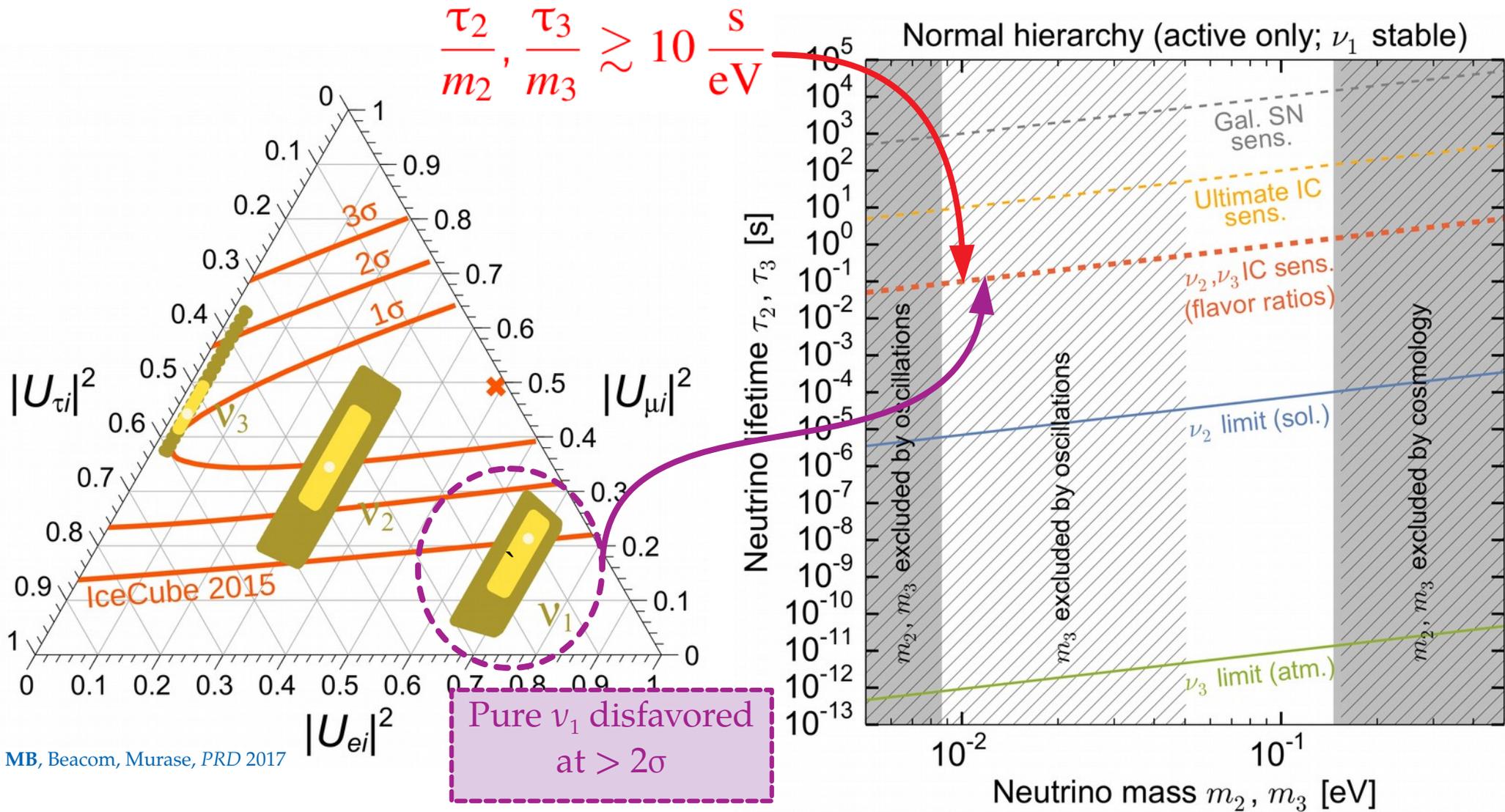
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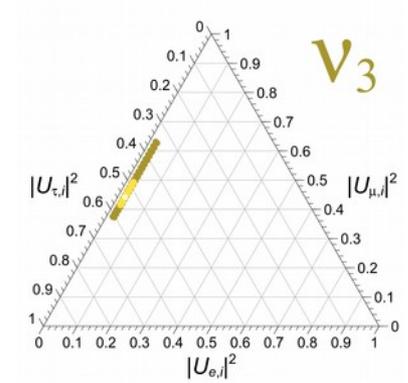
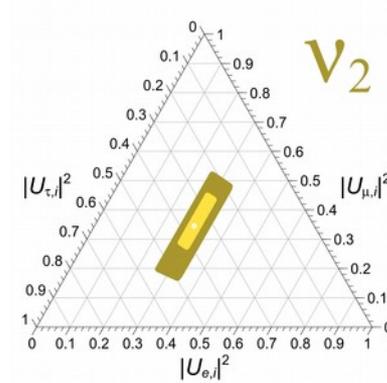
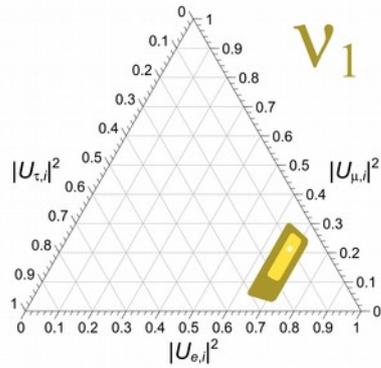


Pure ν_1 disfavored at $> 2\sigma$



Two classes of new physics

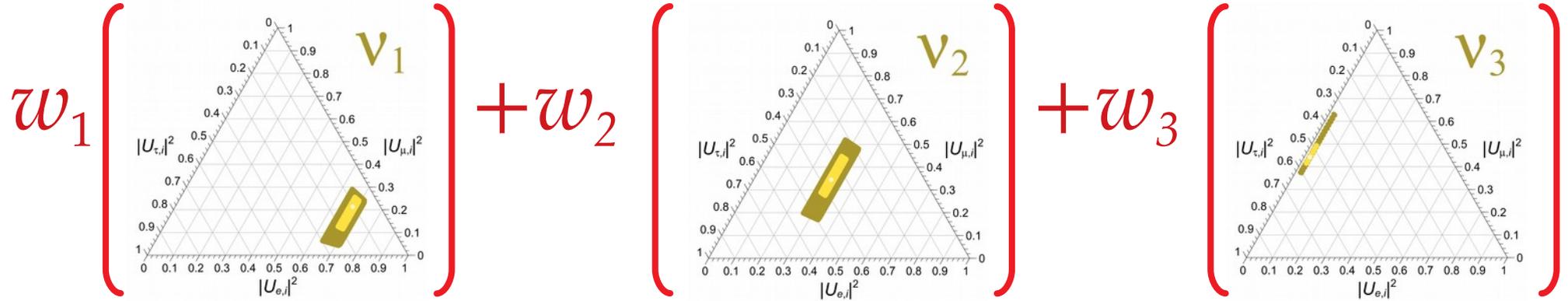
- ▶ Neutrinos propagate as an incoherent mix of ν_1, ν_2, ν_3
- ▶ Each one has a different flavor content:



- ▶ Flavor ratios at Earth are the result of their **combination**
- ▶ New physics may:
 - ▶ Only reweigh the proportion of each ν_i reaching Earth (*e.g.*, ν decay)
 - ▶ Redefine the propagation states (*e.g.*, Lorentz-invariance violation)

Two classes of new physics

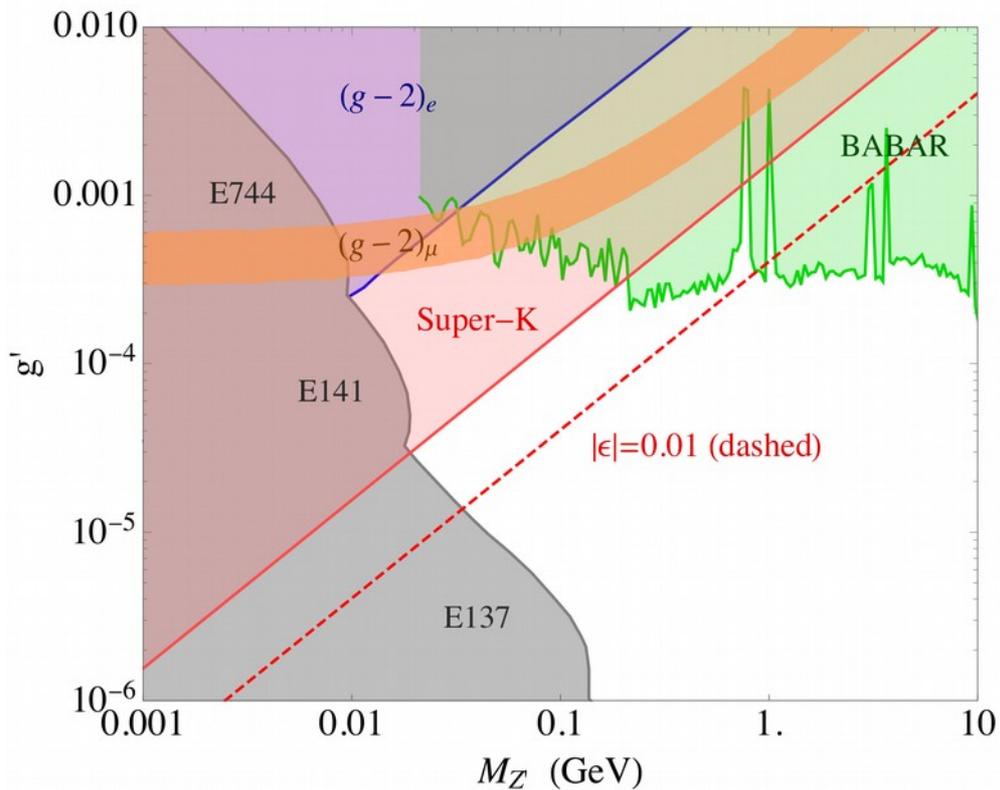
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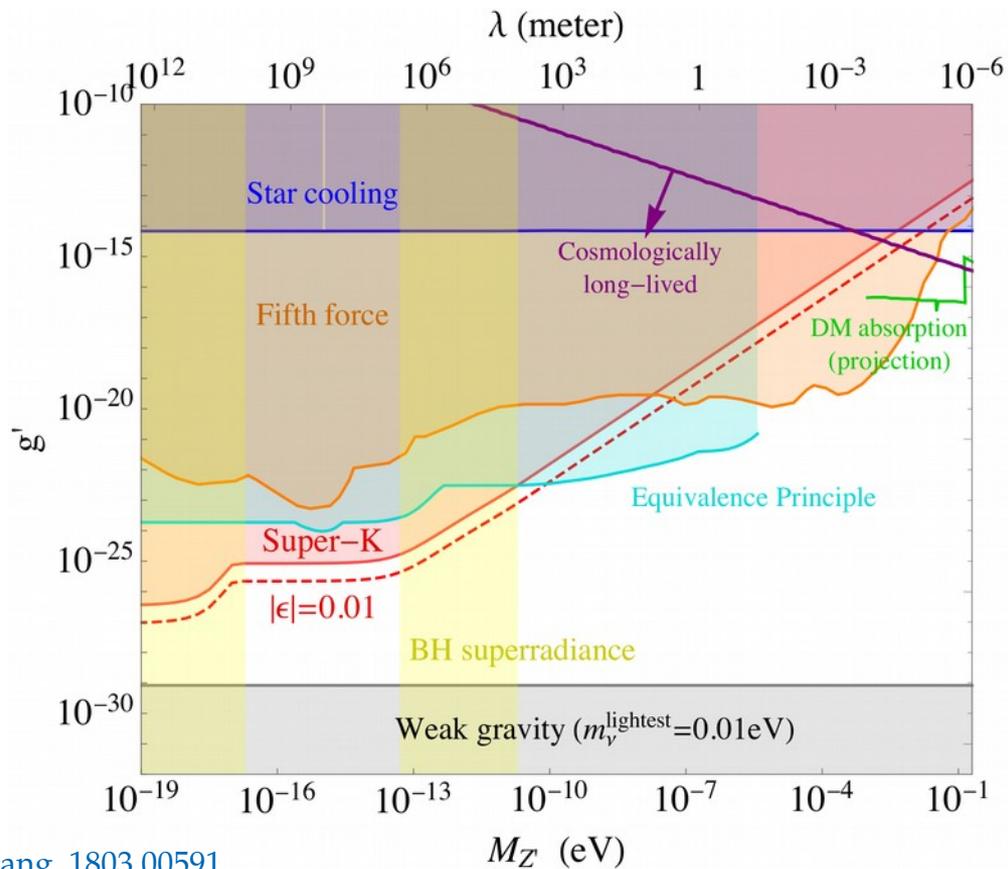
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 - ▶ Redefine the propagation states (*e.g.*, Lorentz-invariance violation)

Current limits on the Z'

MeV–GeV masses



Sub-eV masses



M. Wise & Y. Zhang, 1803.00591

Connecting flavor-ratio predictions to experiment

- 1 Integrate potential in redshift, weighed by source number density
→ Assume star formation rate

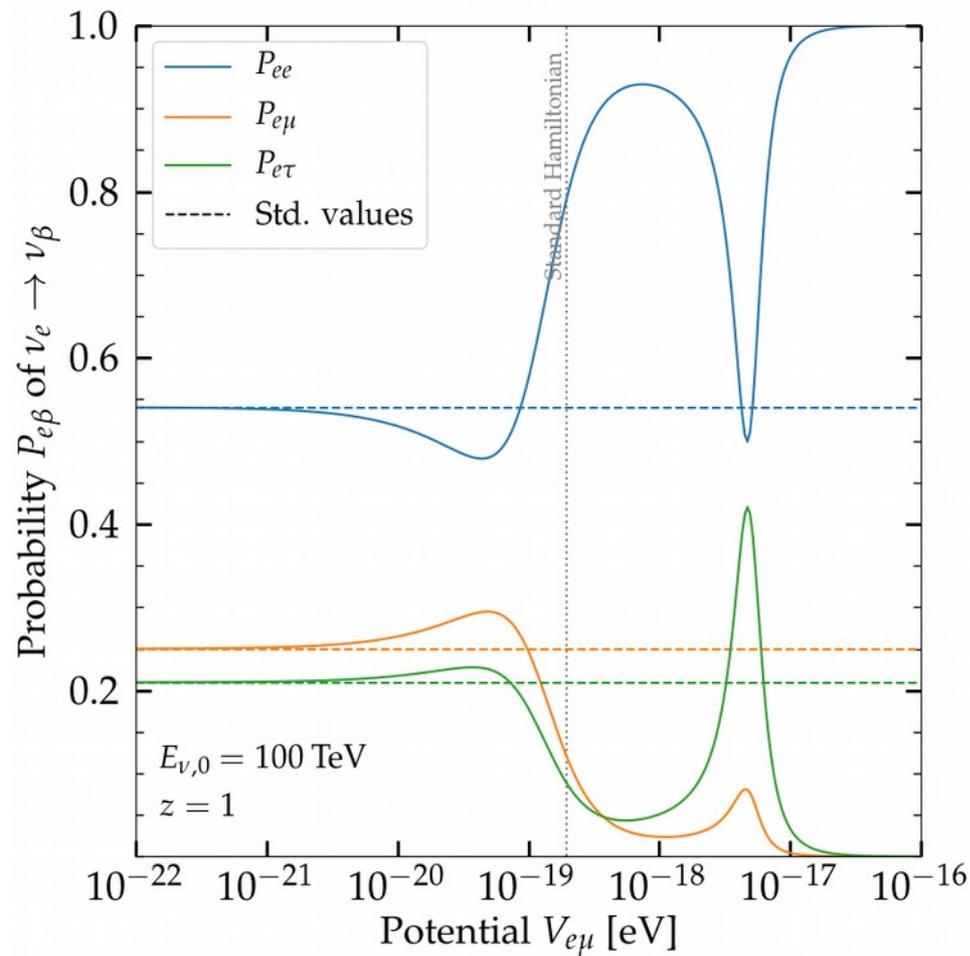
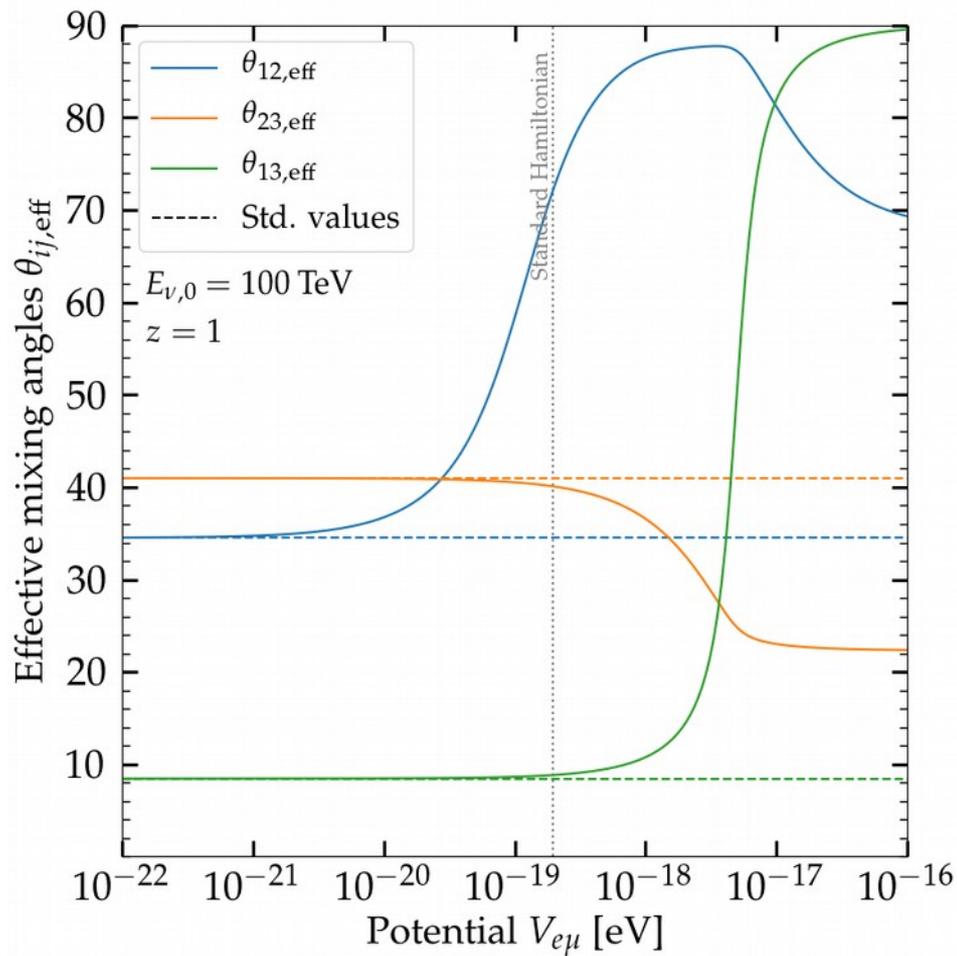
$$\langle V_{e\beta}^{\text{COS}} \rangle \propto \int dz \rho_{\text{SFR}}(z) \cdot \frac{dV_c}{dz} \cdot V_{e\beta}^{\text{COS}}(z)$$

Density of cosmological e grows with z

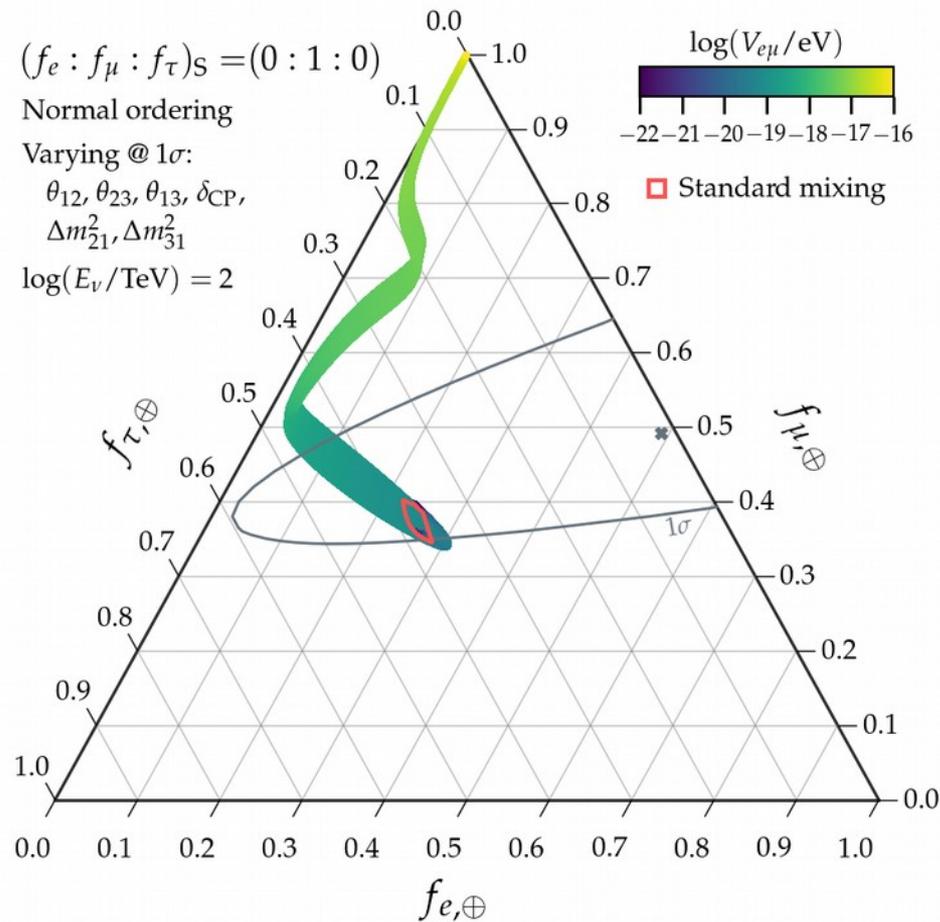
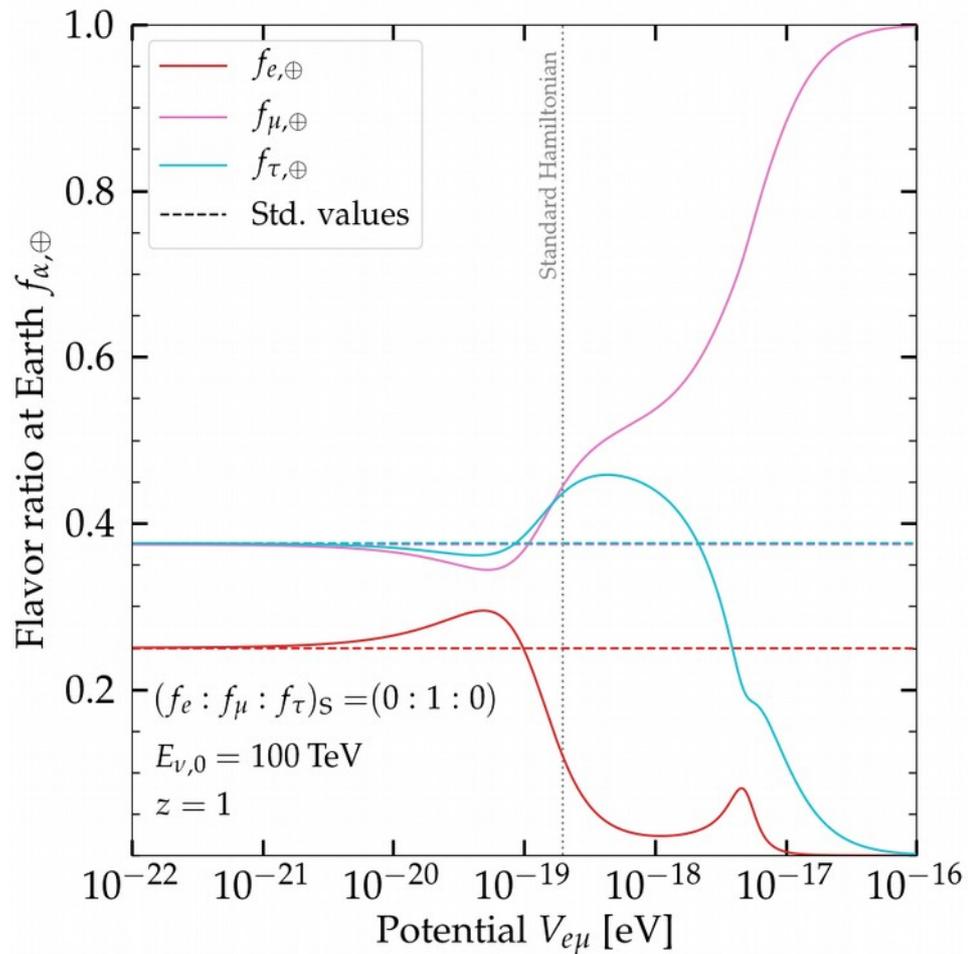
- 2 Convolve flavor ratios with observed neutrino energy spectrum
→ Either $E^{-2.50}$ (combined analysis) or $E^{-2.13}$ (through-going muons)

$$\underbrace{\langle \Phi_\alpha \rangle \propto \int dE_\nu f_{\alpha,\oplus}(E_\nu) E_\nu^{-\gamma}}_{\text{Energy-averaged flux}} \Rightarrow \underbrace{\langle f_{\alpha,\oplus} \rangle \equiv \frac{\langle \Phi_\alpha \rangle}{\sum_{\beta=e,\mu,\tau} \langle \Phi_\beta \rangle}}_{\text{Energy-averaged flavor ratios}}$$

Resonance due to the L_e-L_μ symmetry

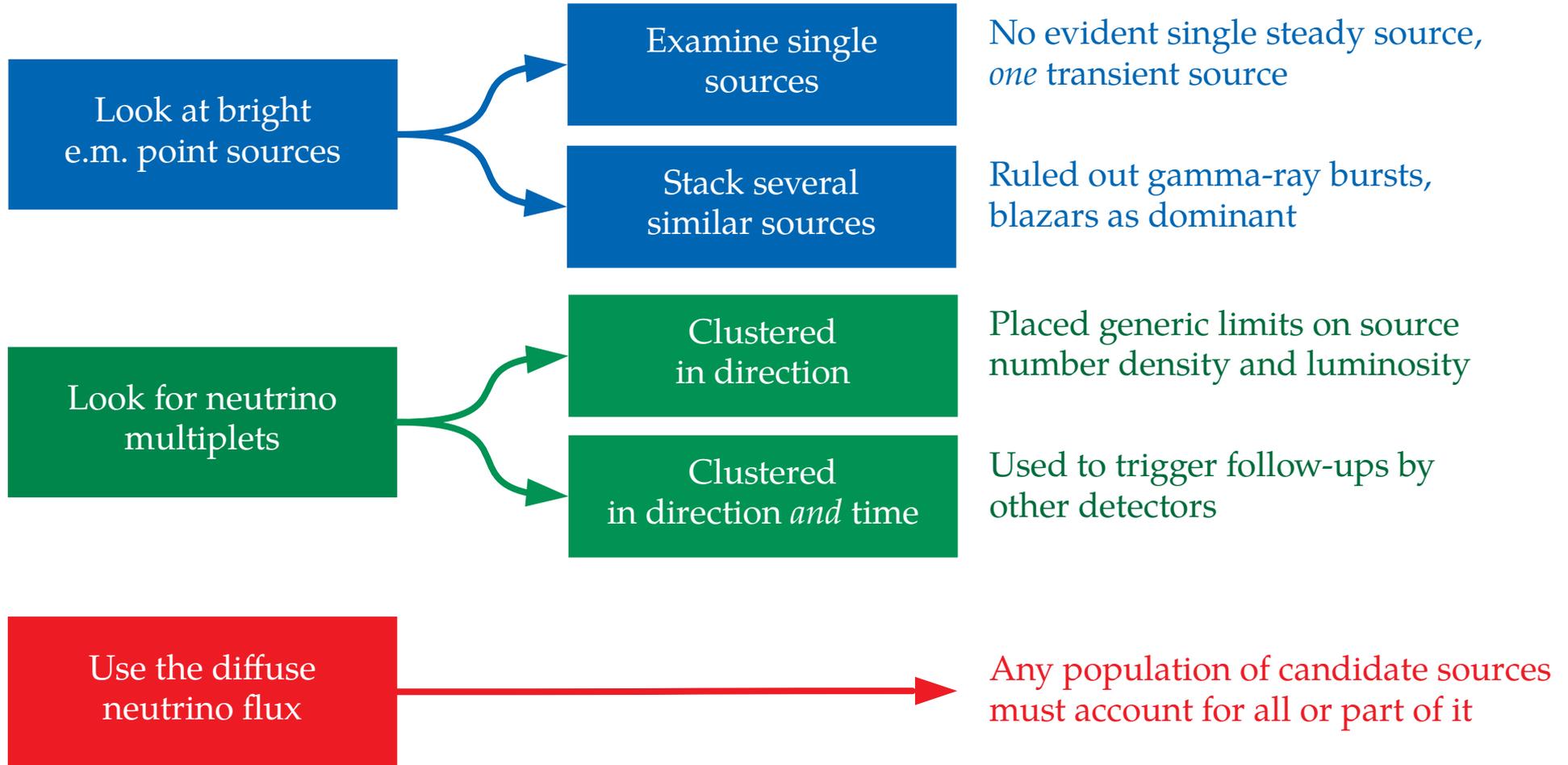


Resonance due to the L_e-L_μ symmetry (cont.)



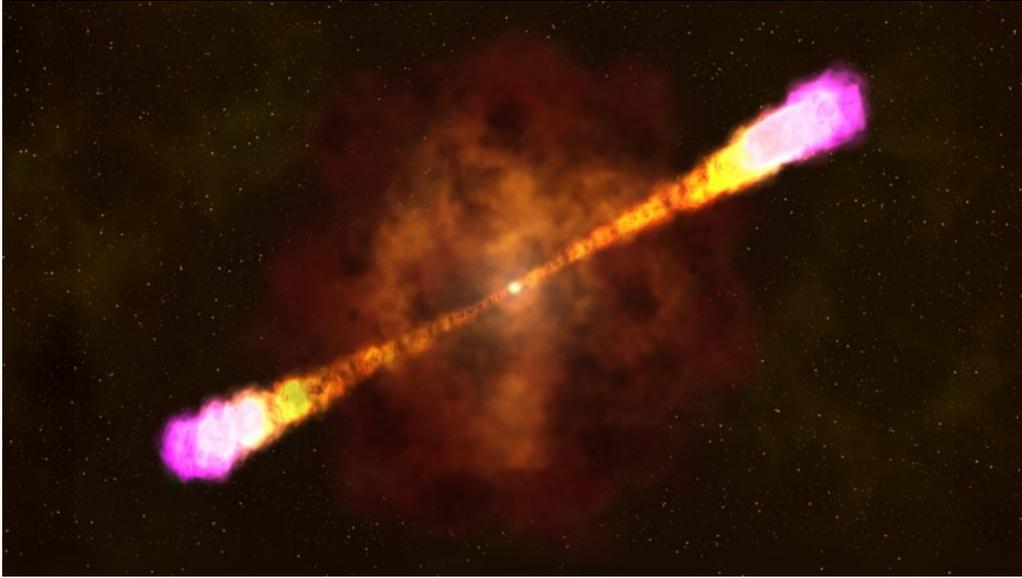
Looking for the sources

Three Strategies to Reveal Sources Using TeV–PeV ν

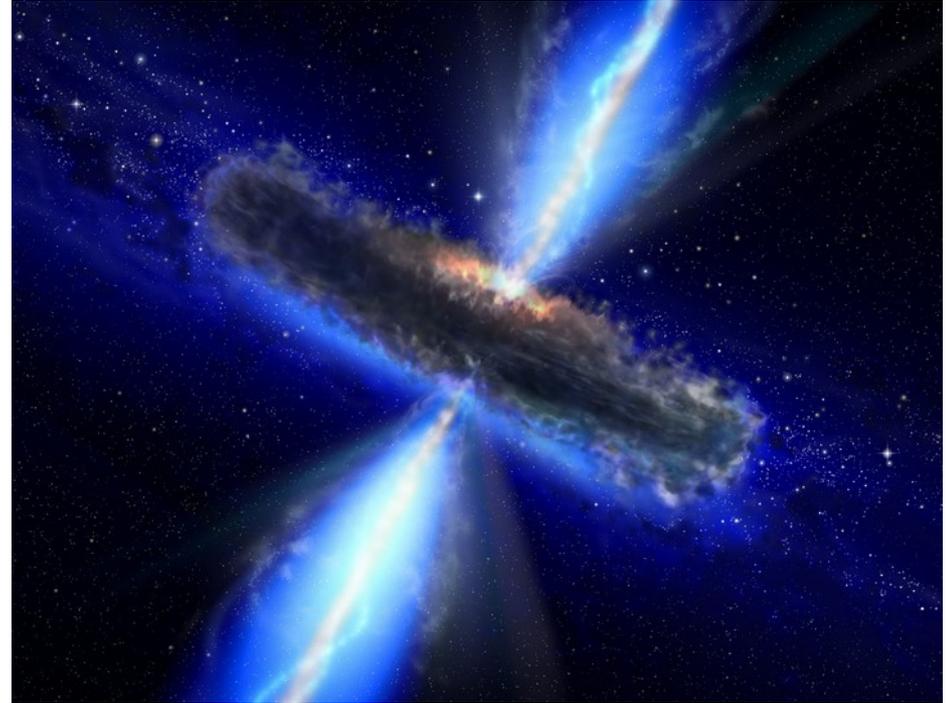


Gamma-ray bursts and blazars – *not* dominant

Gamma-ray bursts

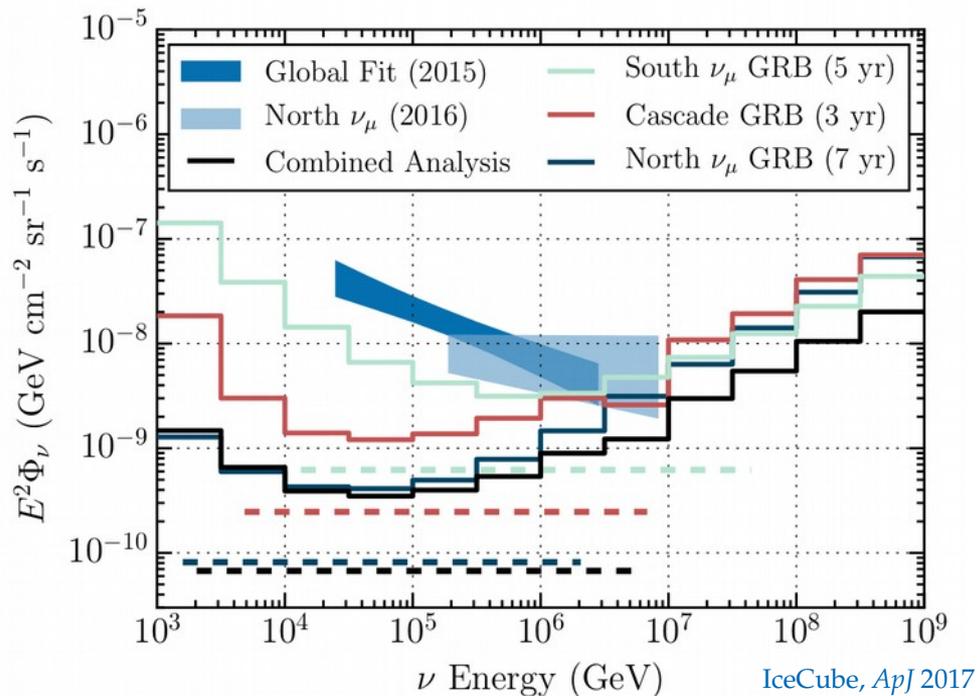


Blazars



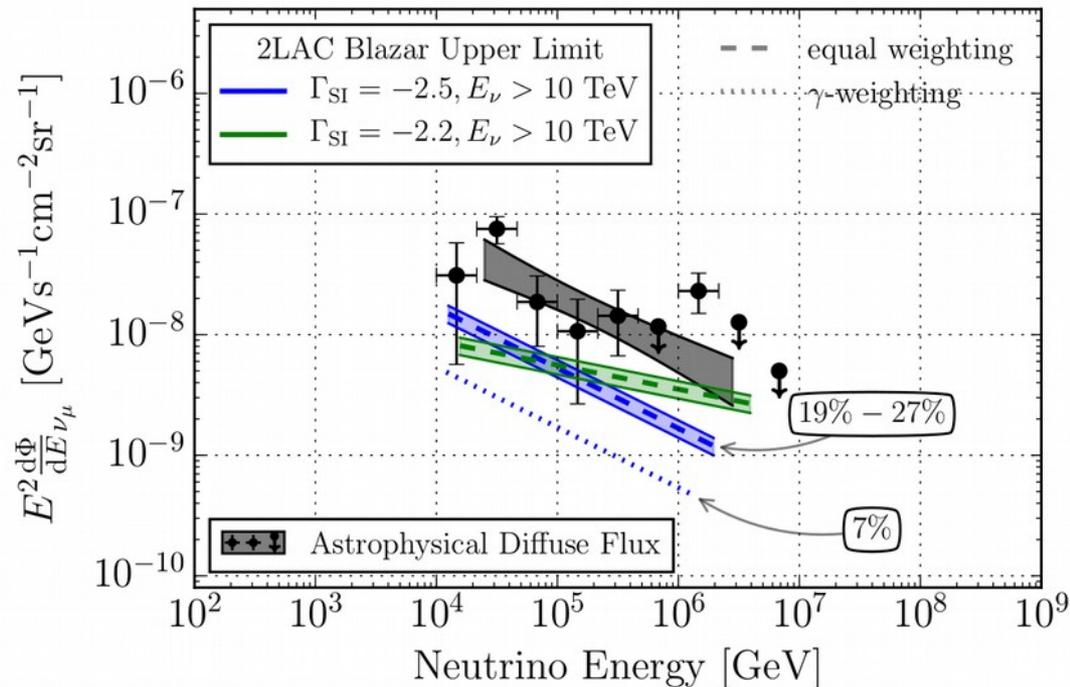
Gamma-ray bursts and blazars – *not* dominant

Gamma-ray bursts



1172 GRBs inspected, no correlation found
< 1% contribution to diffuse flux

Blazars

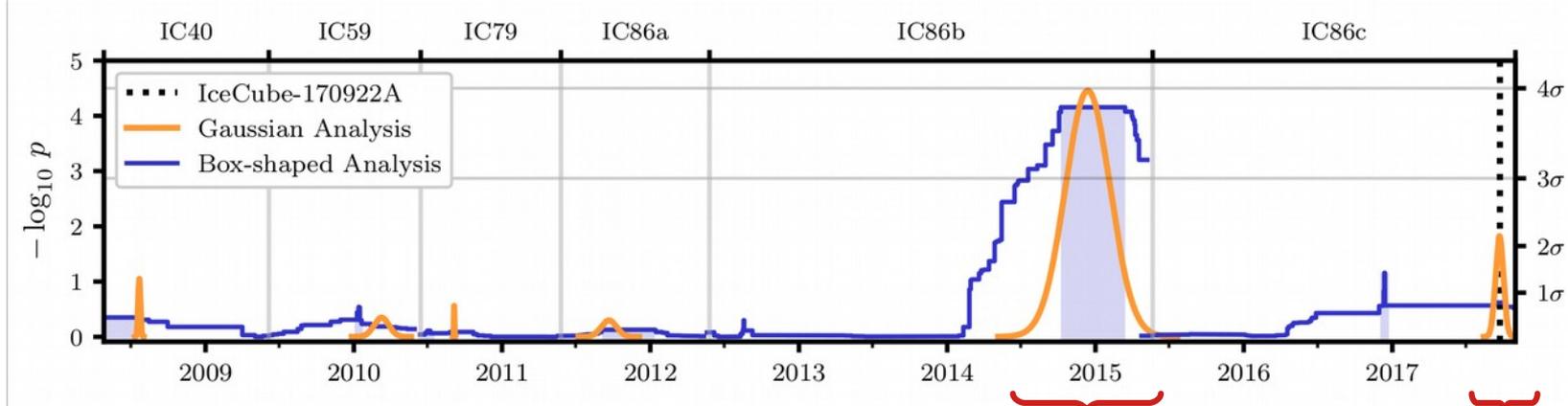


862 blazars inspected, no correlation found
< 27% contribution to diffuse flux

... but we have seen *one* blazar neutrino flare!

Recent news:
The starburst Seyfert galaxy NGC 1068 is also a potential neutrino source candidate (1908.05993)

Blazar TXS 0506+056:



Important:

If every blazar produced neutrinos as TXS 0506+056, the diffuse neutrino flux would be 20x higher than observed!

2014–2015: 13 ± 5 ν flare, no X-ray flare
3.5 σ significance of correlation (post-trial)

2017: one 290-TeV ν + X-ray flare
1.4 σ significance of correlation

Combined (pre-trial): 4.1 σ

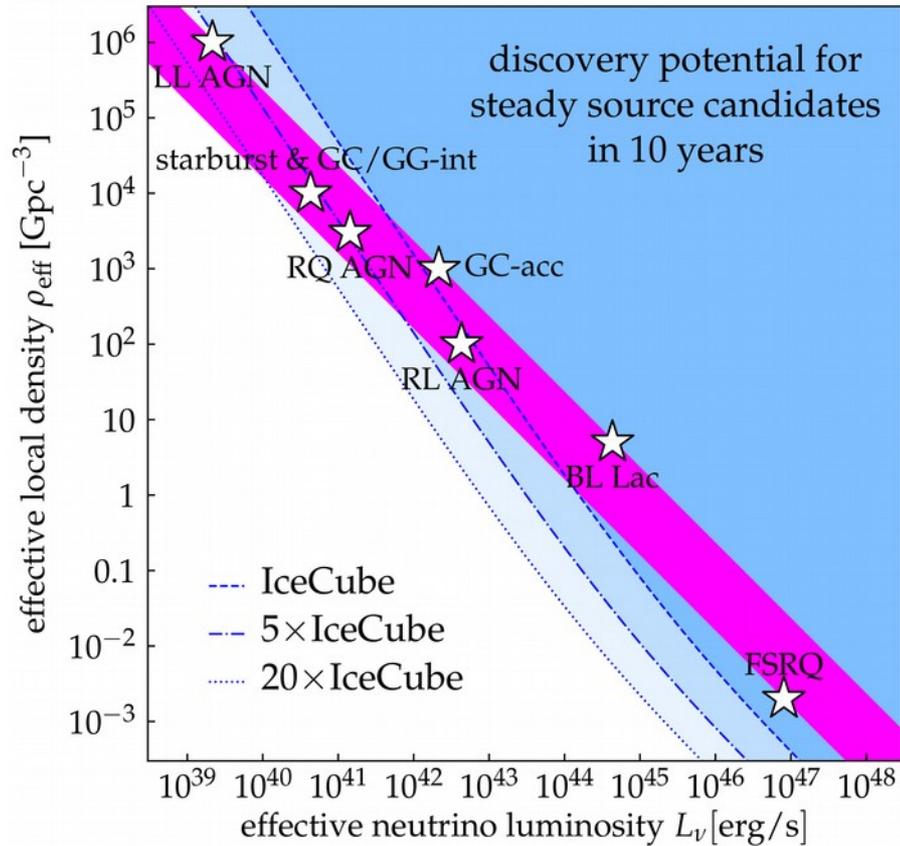
Hard fluence: $E^2 J_{100} = 2.1_{-0.7}^{+0.9} \left(\frac{E}{100 \text{ TeV}} \right)^{-2.1 \pm 0.2} \text{ TeV cm}^{-2}$

Joint modeling of the two periods is challenging; see ICRC 2019 talk by Walter Winter

Source discovery potential: today and in the future

■ Accounts for the observed diffuse ν flux (lower/upper edge: rapid/no redshift evolution)

Closest source with $E^2\Phi_{\nu_\mu+\bar{\nu}_\mu} = 10^{-12} \text{ TeV cm}^{-2} \text{ s}^{-1}$



Closest source with $E^2F_{\nu_\mu+\bar{\nu}_\mu} = 0.1 \text{ GeV cm}^{-2}$

