Probing dark matter with diffractive gravitational lensing of GRBs and FRBs

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w/Andrey Katz, Joachim Kopp and Wei Xue, Diego Blas Temino 1807.11495, 1912.xxxxx

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Multimessengers, Prague, 5/12/19

Which Dark Matter ?

MACHOs - MAssive Compact Halo Objects

 $10^{-17} M_{\odot} < M < 10^3 M_{\odot}$

Examples :

Carr & Hawking (1974)

- primordial black holes
- axion miniclusters / ultracompact minihalos





Vaquero, Redondo & Stadler (2018)



Eggemeier et al. (2019)



How to search?

Gravitational lensing



Many surveys inside Milky Way / Local Group: MACHO, EROS, OGLE, SUBARU. Use microlensing technique: look for variations of observed intensity of background stars

Lensing at cosmological scales

Advantages :

- optical depth $n_{dm}r_E^2D_S = 4G\rho_{dm}D_S^2 \propto (H_0D_S)^2\Omega_{dm}$
- DM mass function can be different in halos (e.g. miniclusters may be destroyed by tidal forces)

Challenges :

• the sources must be smaller than the Einstein ring = pointlike



- angular separation between the images $\sim \mu as~$ for lens of solar mass — too small to be resolved

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Gamma Ray Bursts and Fast Radio Bursts

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Diffractive lensing

Idea: interference between two unresloved images produces fringes in the frequency spectrum

$$Iensing time delay$$

$$F_{\rm obs}(t) = AF_{in}(t) + BF_{\rm in}(t - \Delta t)$$

$$f_{\rm obs}(\omega) = A f_{in}(\omega) + B f_{\rm in}(\omega) e^{i\omega\Delta t}$$

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longing time delay

$$|f_{\rm obs}(\omega)|^2 = |f_{in}(\omega)|^2 (A^2 + B^2 + 2AB\cos\omega\Delta t)$$

Diffractive lensing

lanaina tima dalau

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Lensing generalities



$$\psi(\theta) = \theta_E^2 \log \theta$$

Einstein angle $\theta_E = \left(4GM\frac{D_{LS}}{D_SD_L}\right)^{1/2}$
 $\theta_{\pm} = \frac{1}{2}\left(\beta \pm \sqrt{\beta^2 + 4\theta_E^2}\right)$
Magnifications: $\mu_{\pm} = \frac{y^2 + 2}{2y\sqrt{y^2 + 4}} \pm \frac{1}{2}$

Total magnification:

$$\mu = \frac{y^2 + 2}{y\sqrt{y^2 + 4}} + \frac{2}{y\sqrt{y^2 + 4}} \sin\left(\Omega\left[\frac{y\sqrt{y^2 + 4}}{2} + \log\left|\frac{y + \sqrt{y^2 + 4}}{y - \sqrt{y^2 + 4}}\right|\right]\right)$$

 $y = \beta/\theta_E$ $\Omega = 4GM(1+z_L)\omega$

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Lens equation: $\theta_{\pm} - \theta_0^{\delta - 1} \theta_{\pm}^{2 - \delta} = \beta$



two images



two images

three or one image



two images

three or one image

single image

Geometric vs. wave optics

Exact magnification:
$$\mu = \left| \frac{\Omega}{(2\pi i)\theta_E^2} \int d^2 \vec{\theta} \, e^{i\omega\Delta t(\vec{\theta})} \right|^2$$



Extended source



Intensity profile $W(ec{y};\sigma_y)$

Fringes average out:

$$\bar{\mu} = \frac{\int d^2 y \, W(\vec{y}; \sigma_y) \, \mu(\vec{y}; \Omega)}{\int d^2 y \, W(\vec{y}; \sigma_y)}$$



Application to GRBs (femtolensing)

Gould (1992) Stanek, Paczynski, Goodman (1993)

Duration Short GRBs from 0.1 to 2 s, Long GRBs from 2 to 200 s

Broad spectrum from 10 keV to 10 MeV



Are GRBs compact enough ?

Transverse size can be estimated from the variability time scale





- majority of GRBs have $a_S \sim 10^{11} {
 m cm} \gg r_E \sim 10^9 {
 m cm}$
- however, 10% have $a_S \sim 10^{10} {\rm cm}$. Is it possible that there are some with $a_S \sim 10^9 {\rm cm}$?

Constraints in future ?



We need to better understand GRB sizes...

Application to FRBs

Zheng et al. (2014) Eichler (2017)

Duration ~ ms

Typically at 400 MHz to 1.5 GHz, but also up to ~ 10 GHz

 $\nu\Delta t\sim 1$ potentially sensitive to $10^{-4}M_{\odot} < M < 0.1M_{\odot}$

No problems with the size of the source (Einstein radius 10^{-13} cm)

Farah et al. (2018)



This is **not** lensing

Scintillation



Scales of scintillation

• phase correlation length $r_{\rm diff}$: $\langle (\varphi(x + r_{\rm diff}) - \varphi(x))^2 \rangle = 1$

 $r_{
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- diffraction angle $\theta_{\rm diff} = \frac{1}{\omega r_{\rm diff}}$ $r_{\rm diff} < r_F \equiv \sqrt{D_{\rm ScO}/\omega}$ strong scintillation
- the light is collected from regions of size $r_{
 m ref}= heta_{
 m diff}D_{
 m ScO}$

includes many incorrelated patches

$$f_{\rm obs}(\omega) = \frac{\omega f_{\rm in}(\omega)}{2\pi i D_{\rm ScO}} \int d^2x \, e^{i\Phi(\omega,\vec{x})} \left(A e^{i\omega\delta t_A(\vec{x})} + B e^{i\omega[\Delta t + \delta t_B(\vec{x})]}\right)$$

variation of time-delay over the screen

 $|f_{\rm obs}(\omega)|^2 = |f_{\rm in}(\omega)|^2 \left(|\mathcal{A}(\omega)|^2 + |\mathcal{B}(\omega)|^2 + \mathcal{A}(\omega)\mathcal{B}^*(\omega) e^{i\omega\Delta t} + \mathcal{A}^*(\omega)\mathcal{B}(\omega) e^{-i\omega\Delta t} \right)$

peak in the FT of intensity spectrum if the amplitude has a slowly varying component

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(the random phases do not change, have smooth frequency dependence) contributes order one if

$$\omega \theta_E r_{\rm ref} < 1$$

(the images are distorted in the same way, "are not resolved by the screen")





Simulating lensing+scintillation



Scintillation outside MW

• in the host galaxy: symmetric, the same analysis as for MW



 in the intergalactic medium: unimportant for realistic assumptions, unless the line-of-sight crosses a galaxy or a cluster

Analysis of simulated data

Idea: Apply peak-finder algorithm to the FFT of the power spectrum



 $M = 0.001 M_{\odot}, \quad r_{diff} = 10^9 \,\mathrm{cm}, \quad y = 0.5, \quad D_L = 0.5 \,\mathrm{Gpc}$

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